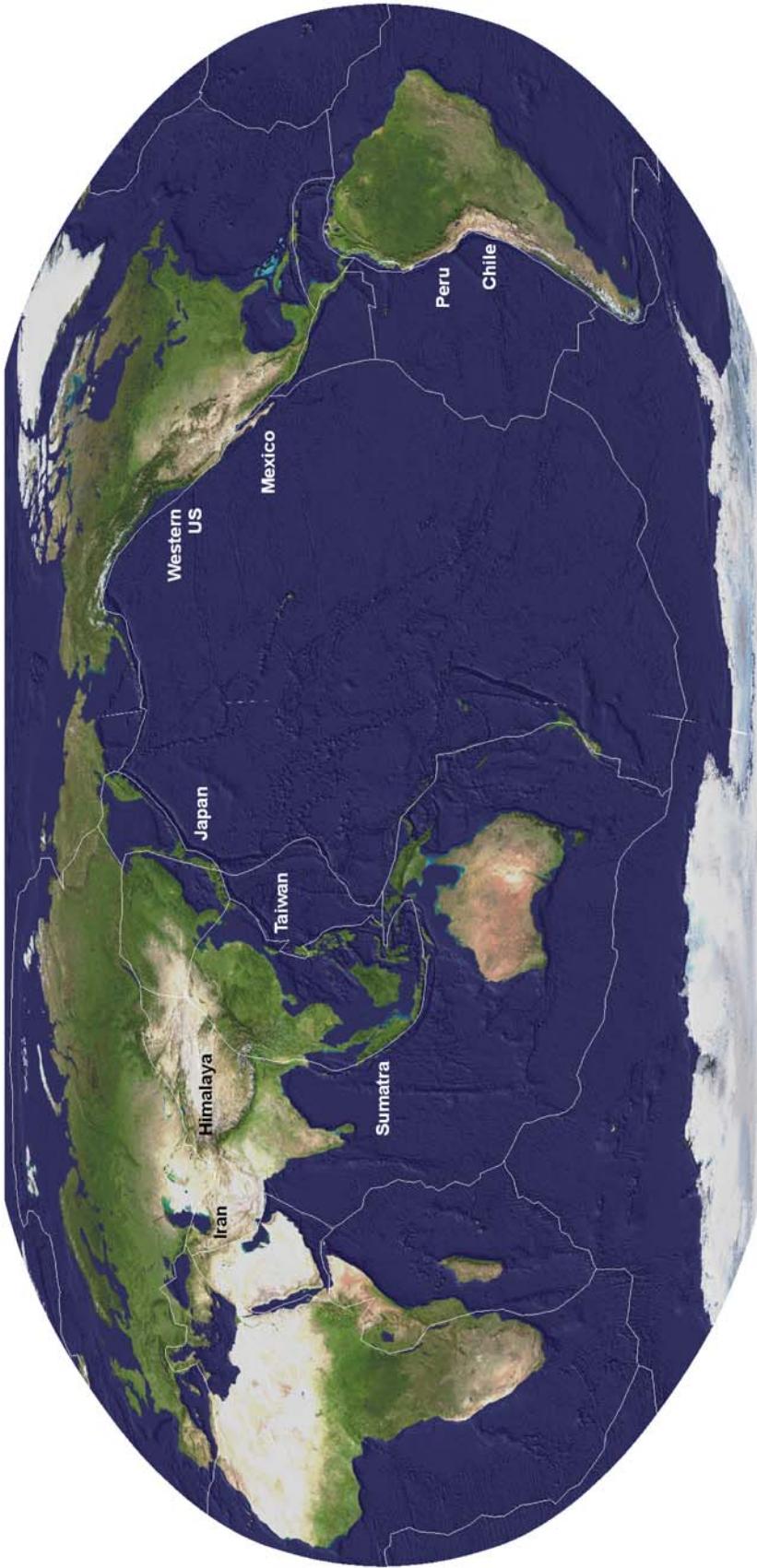


Tectonics Observatory



Annual Meeting 2006
Poster Session



^3He Cosmogenic Dating in Zircon, Apatite and Titanite

William Amidon¹, Kenneth Farley,¹ Douglas Burbank²

¹⁾California Institute of Technology

²⁾UC Santa Barbara

What is cosmogenic dating?

Cosmogenic isotopes are produced in rock by the interaction of energetic particles with "target" nuclei in the rock. By measuring the abundance of ^3He (or ^{10}Be , ^{21}Ne , ^{26}Al , ^{36}Cl), the exposure age of surfaces in the landscape can be calculated.

How does it apply to active tectonics?

Knowing the age of surfaces in the landscape allows the calculation of the rate at which the surface is deformed by tectonic or erosional processes. This includes; faulted and folded terraces, alluvial fans, offset moraines, bedrock and basin-scale erosion, the age and rate of soil development, and more!

Why develop ^3He dating in these phases?

Conventional cosmogenic dating using ^{10}Be , ^{26}Al or ^{36}Cl requires time consuming and expensive preparation, and measurements on an accelerator mass spectrometer. In contrast, ^3He is measured cheaply and easily on a mobile gas mass spectrometer. Previous use of ^3He has been primarily in olivine and pyroxene in volcanic rocks. ^3He dating in zircon, apatite and titanite will allow fast, cheap cosmogenic dating of more common rock types.

Our approach...

Our goal is to determine the production rate of ^3He in quartz and apatite by direct cross calibration to the known production rate of ^{10}Be in quartz. We separated minerals and measured He in glacial moraine boulders from the Nepal Himalaya, which were previously dated by ^{10}Be .

Methods

1) 5–25 mg of mineral is degassed under vacuum by heating of a platinum packet with a Nd-YAG laser.

2) Gas is purified in a vacuum line before ^4He and ^3He intensities are measured by peak switching between a faraday cup and an electron multiplier using an MAP-250 noble gas mass spectrometer.

3) ^3He Concentrations are calculated by comparison with a He standard.

Analytical concerns

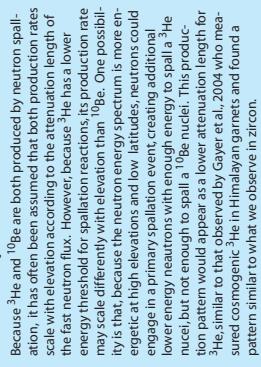
1) **Pressure broadening of the ^4He peak:** to address this concern samples of thoronite (ThO_2) were used to generate high pressures of pure ^4He . It was found that the ^4He peak was not "tailoring" onto the ^3He peak at high pressures.

2) **Space-charge effects:** because machine sensitivity is calibrated with standards at low He pressures, it may not reflect the true sensitivity for higher pressures induced by the sample gas. This problem is solved by introducing a "spike" of ^3He rich standard during the analysis to evaluate the sensitivity during each analysis.

Previous Work:

Farley et al., 2006 calibrated the production rate of ^3He against ^{21}Ne in a Bolivian tuff. A production rate of 80 at/g/yr was found for zircon, 105 at/g/yr in apatite and 90 at/g/yr in titanite. These depth profiles show exponential decay of ^3He with depth in the rock, confirming a cosmogenic source for the ^3He .

Is overproduction increasing with elevation?

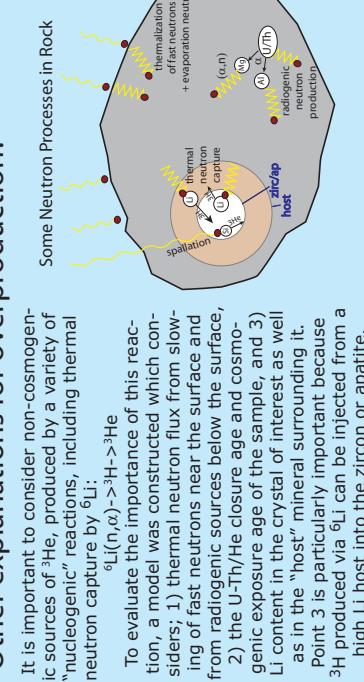


Plot A shows the importance of grain size in determining the production of ^3He from ^6Li . Production in small grains is dominated by implanted ^3He , and is linearly related to Li content in the host. Larger grains have most of their ^3He produced internally, and thus the apparent production rate from Li flattens when the host Li content becomes very low. Plot B shows the expected ^3He produced from ^6Li as a function of exposure age and $\text{U}-\text{Th}/\text{He}$ closure age. This result is for a 1 ppm zircon/apatite embedded in a quartzite at 4500 m, and shows that production from ^6Li is <5% of the total observed ^3He .

Conclusions and future work:

- Based on the work of Farley et al., 2006 it appears that this technique is ready to be applied to low-Li samples, in settings outside of Nepal.
- It is likely that a problem exists with the currently accepted scaling models when applied in the Nepal Himalaya. To confirm this we plan to:
 - measure Li in more samples of zircon and apatite
 - explore hand sample petrography to determine "host" minerals
 - measure 3He and closure age in Li-rich phases to estimate thermal neutron flux in the rock

Other explanations for overproduction?



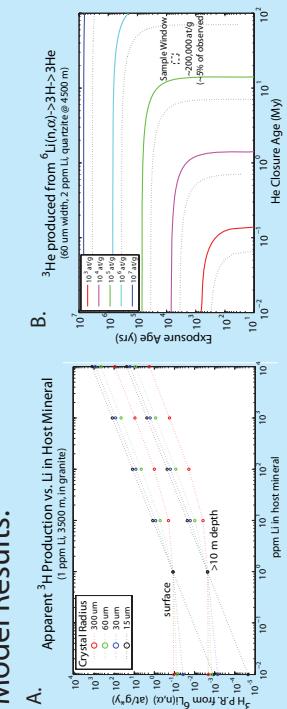
It is important to consider non-cosmogenic sources of ^3He , produced by a variety of "nucleogenic" reactions, including thermal neutron capture by ^6Li :

To evaluate the importance of this reaction, a model was constructed which considers: 1) thermal neutron flux from slow-moving fast neutrons near the surface and from radiogenic sources below the surface, 2) the $\text{U}-\text{Th}/\text{He}$ closure age and cosmogenic exposure age of the sample, and 3) Li content in the crystal of interest as well as in the "host" mineral surrounding it. Point 3 is particularly important because ^3He produced via ^6Li can be injected from a high Li host into the zircon or apatite.

Measuring Li:

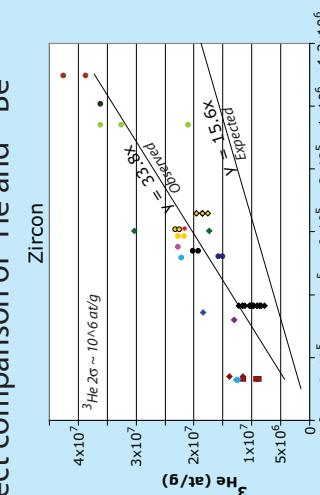
Li contents have been measured on apatites using the ICPMS at Caltech. Results show good reproducibility, suggesting Li can be accurately measured by this technique. Observed values range from ~1–3 ppm, with similar or lesser values expected for zircon.

Model Results:

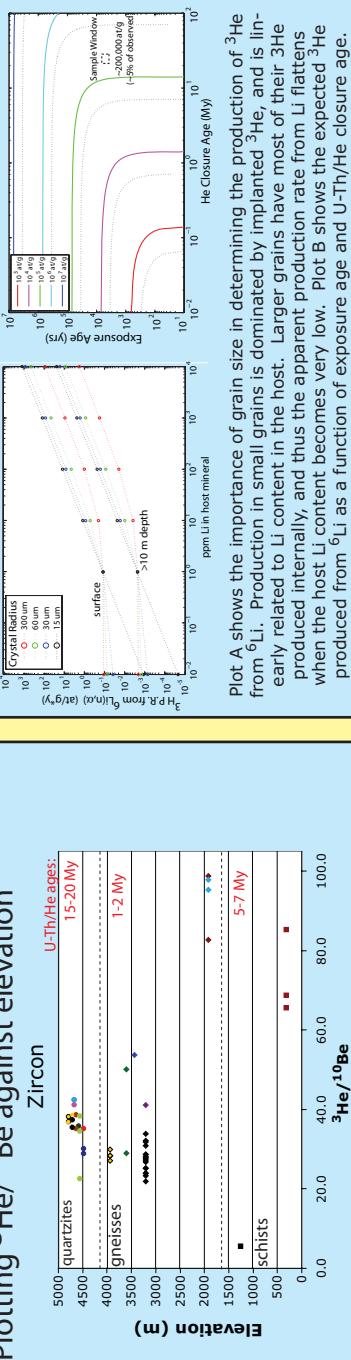


Plot B shows the importance of grain size in determining the production of ^3He from ^6Li . Production in small grains is dominated by implanted ^3He , and is linearly related to Li content in the host. Larger grains have most of their ^3He produced internally, and thus the apparent production rate from Li flattens when the host Li content becomes very low. Plot B shows the expected ^3He produced from ^6Li as a function of exposure age and $\text{U}-\text{Th}/\text{He}$ closure age. This result is for a 1 ppm zircon/apatite embedded in a quartzite at 4500 m, and shows that production from ^6Li is <5% of the total observed ^3He .

Direct comparison of ^3He and ^{10}Be



Plotting $^3\text{He}/^{10}\text{Be}$ against elevation



Because ^3He and ^{10}Be are both produced by neutron spallation, it has often been assumed that both production rates scale with elevation according to the attenuation length of the fast neutron flux. However, because ^3He has a lower energy threshold for spallation reactions, its production rate may scale differently with elevation than ^{10}Be . One reason for this is that the neutron energy spectrum is more energetic at high elevations, and lower latitudes, neutrons could engage in primary spallation events, creating additional lower energy neutrons with enough energy to split a ^{10}Be nucleus but not enough to split a ^3He nucleus. This production pattern would appear as a lower attenuation length for ^3He similar to that observed by Gayer et al., 2004 who measured cosmogenic ^3He in Himalayan garnets and found a pattern similar to what we observe in zircon.

Gayer et al., 2004



Sedimentary record of erosion and deformation rates in the northern Tianshan as constrained from magnetostratigraphic sections in the Junggar basin, western China

¹ Institut des Sciences de la Terre d'Orléans, France; ² California Institute of Technology, USA; ³ Université de Montpellier II, France; ⁴ Institut des

Introduction

INTRODUCTION The growth of a mountain range results primarily from crustal thickening driven by horizontal shortening. As mountain ranges grow, the processes that become more active and operate more frequently from the uplifted zone of the mountain range transfer material to the interior. Mountain building is therefore a complex process resulting from coupling between tectonic deformation, surface processes and climate. The redistribution of mass at the surface is key factors determining tectonic deformation. In Central Asia, the high Tian Shan mountains are blanketed by the Tarim basin, which is a large clastic system derived from the Tian Shan. This study aims to understand the formation and through-going of the southern boundary of the Tarim basin and the development of the basin. This study is particularly appropriate to both sedimentary and southern pre- Andean geological settings of the Andes.

(1) Kinematic framework

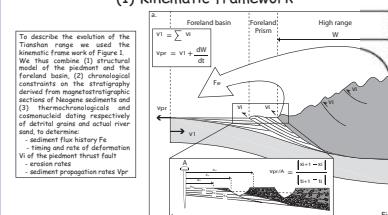
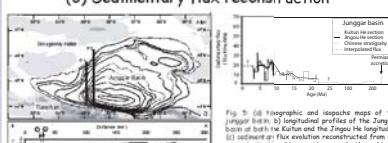


Fig. 1: schematic sketch defining the kinematic framework used in this study. V_I is the underthrusting velocity of the foreland; V_t are the shortening rates of the contractional thrust system and the extensional thrust system.

(3) Sedimentary flux reconstruction



Assuming that the relative thickness of U is uniform throughout, the total value V above a z in the whole basin may be expressed as follows: $V = \rho_z / \rho_{\text{B}}$, where ρ_z is the true density of a z layer at a given vertical z , ρ_{B} is the depth of basin bottom at the same profile. The volume V is proportional to the area of the basin. Yet, horizontal action along z in the midpoint may be considered as vertical profile z through the basin. We thus reconstruct the sediment flux accumulation in the Junggar Basin based on the stratigraphy constrained by the seismic reflection data.

(6) Preliminary results of U-Th/He thermochronology

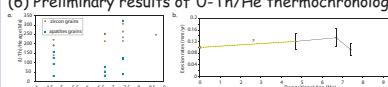


Fig. 8: (a) plot of U-Th/He of detrital grains ages versus depositional ages ; (b) Evolution of the erosion rates deduced from the Lag-time concept with respect to time, yellow circle represent erosion rates calculated from cosmogenic dating of actual river sand of the Kuitun River

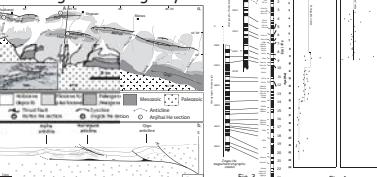
In order to decipher the late exhumation history of the Tienhsin mountains, we carried out U-Th/He dating on both apatite and zircon detrital grains collected in the Kehan section. Results are summarized on Figure X. From the young ages obtained, we conclude that the Tienhsin mountains were exhumed during the late stage of the log-time concept. We found an average of ~ 1.1 m.y./stage.

In order to decipher the late exhumation history of the Tienhsin mountains, we carried out U-Th/He dating on both apatite and zircon detrital grains collected in the Kehan section. Results are summarized on Figure X. From the young ages obtained, we conclude that the Tienhsin mountains were exhumed during the late stage of the log-time concept. We found an average of ~ 1.1 m.y./stage.

assign it to a global effect of climate change on erosion rates (Reznichenko et al., 2001). Our interpretation is that this gravelsheet was prograded over the foreland as it was underthrusting beneath the orogenic wedge. We reject the common assumption that these conglomerates would reflect climate change or a particular episode of thrusting and uplift of the Tien

assign it to a global effect of climate change on erosion rates (Reznichenko et al., 2001). Our interpretation is that this gravelsheet was prograded over the foreland as it was underthrusting beneath the orogenic wedge. We reject the common assumption that these conglomerates would reflect climate change or a particular episode of thrusting and uplift of the Tien

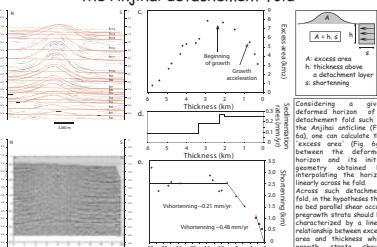
(2) Geological setting and magnetostratigraphy



2: (a) geological map of the northern piedmont with localisation of the Kuitun He and the Jingou He section; (b) magnetic cross section of the northern piedmont at the Jingou He longitude.
 3: Compacted stratigraphic column of the Jingou He and the Kuitun He section and their correlation to the reference scale.
 4: Magnetostriatigraphic column of the Jingou He and the Kuitun He section deduced from the magnetostriatigraphic

Tianshan is a 2500-km-long range with an average altitude of 2500 m with summit up to 7000 m high. The range was probably built during Devonian to Carboniferous. During Cenozoic, it has been living tectonically. The Tianshan plateau is situated in the Urumqi prairie capital to the Dzungaria basin. It is constituted by three rows of thrust fault and fold fields.

(4) Deformation rates across the Anjihai detachment fold



6. (a) line drawing of a seismic line ran across the Anjihai anticline and localisation of the horizons used to calculate the Excess area; (b) original seismic line (from Dengfa et al., 2005); (c) Excess area against the thickness of the 15 horizon; (d) sedimentation rates against space, deduced from the microfacies/marker-concretion; (e) Shalestone content of the

(5) The conglomerate Xiyu formation and its propagation notes

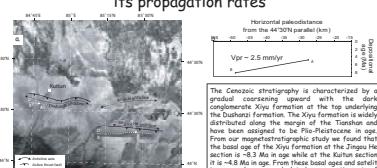


Fig. 7. (a) Lateral image of the Dushanze area with localities of the Kuhun He and the Jiusu He section and the line between the Dushanze and the Xian formation; (b) east of the headwall of the Xian formation across the Dushanze area, we measured a progradation rate of sedimentation Vpr = -2.5 cm yr⁻¹ before shortening, the initial position of the section before shortening.

The Izu-Bonin-Mariana and Costa Rican Subduction Factories Modeled by GyPSM-S: Evidence for a Low-Viscosity Channel

Laura Baker*, Paula Smith, Michael Gurnis, and Paul Asimow



Tectonics Observatory
California Institute of Technology, Pasadena, CA 91125
*labbaker@gps.caltech.edu

Abstract:

We use GyPSM-S, a model that fully couples the petrological model PHMELTS with a 2D thermal and variable viscosity flow model, to describe and compare fundamental processes occurring within the Izu-Bonin-Mariana and Costa Rican subduction zones. By prescribing forcing functions within GyPSM-S we are able to establish the thermal state and paleo-kinematics of the subducting oceanic slab and adjacent mantle wedge and constrain fluid flux specific to different subduction zones. This allows us to describe the process of hydration of the mantle wedge adjacent to the slab, which leads to the development of a low-viscosity channel. We discuss the impact of the low-viscosity channel on petrophysical observables such as dynamic topography and aridity. Also, we define regions of melting and discuss melt and residue chemistry. This leads to predictions of major and trace element chemistry of arc volcanoes.

Parallel computing within the context of GyPSM-S uses a Lagrangian particle distribution to perform thousands of thermodynamically enzatillated calculations during each iteration allowing for a continuous updating chemical dataset and a strong feedback mechanism with the fluid dynamics. Advection of melt, dynamics and compositionally-inherently independent porosity, a condition linking between the effect of water addition to the mantle wedge and the resulting velocity field leading to large-scale changes in the flow field of the high- and wedge relative to the anhydrous wedge adjacent to the site of hydrated material at subduction zones. Adult analysis is available for chemical feedback. Additionally, we account for a transient feature by introducing latent heat to account for a seismic topic: capillary in a phreatic zone.

The flexibility of PHMELTS allows for a slow fractionation in melting to be included and for the calculation of a realistic composition of arc volcanoes.

We present four cases, encompassing the Costa Rican and Izu-Bonin-Mariana subduction zones (Southeastern Costa Rica, Central Costa Rica, Northern Izu-Bonin, and Northern Marianas), and compare our model results with geophysical and geochemical observations from these localities. Among these examples, model-specific changes in potential energy, slab dip, thermal age, subduction velocity, and overriding lithosphere thickness allow us to compare a wide variety of behavior previously unresolved by petrodynamic models of subduction. A continuous, slab-adjacent low-viscosity channel defined by hydrous mineral stability and higher concentrations of nominally anhydrous minerals develops.

Parallel-e-computing structure allows thousands of PHMELTS calculations to occur at each iteration:

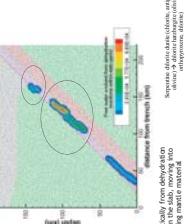
Method:

GyPSM-S is a Geodynamic and Petrological Synthesis Model for Subduction zones. It couples two modules: ConMan, a 2D thermal and variable viscosity flow model (Kring et al., 2009) and PHMELTS, a petrological model that couples reaction kinetics and phase equilibria of melts, solids, and water solubility in nominally-anhydrous minerals (via a next-generation interface robust; Jephcott & Asimow, 2004) with governing Perseus scripts. GyPSM-S is built to run in a parallel environment on Clusters.

GyPSM-S is a fully-coupled formulation, whereby iterative transfer of PHMELTS to finite particle data to the finite element mesh resolves the dynamic feedbacks between the element dynamics and the chemistry that are unobserved in other more simple models.

The evolution of the pressure-temperature state of the downgoing slab produces a flow field and a hydrodynamic reaction rate field in the overlying wedge. The water is vertically advected and reacted with surrounding partially-hydrated peridotite and olivine, producing hydrous phases or an increase in olivine content on a nominally-anhydrous minerals.

Modifications to ConMan include the use of an entropic-avective scheme to allow for latent heat effects as well as a compositional component to the viscosity formulation.



Central Costa Rica:
18 Ma subducting slab age
15 mm/yr convergence rate
35 degree slab dip
20 km overriding lithosphere

SE Izu-Bonin:
90.0 Ma subducting slab age
30 degree slab dip
30 km overriding lithosphere

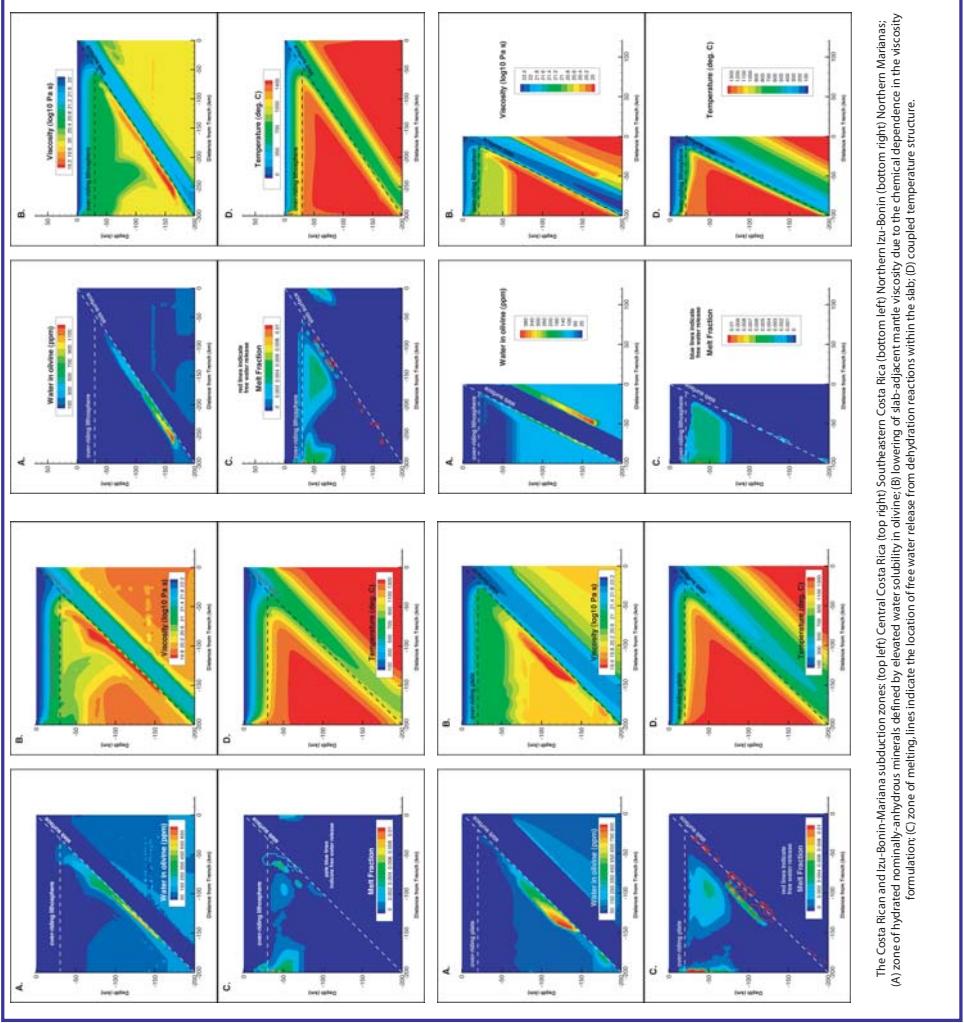
Northern Marianas:
165 Ma subducting slab age
60 degree slab dip
20 km overriding lithosphere

Conclusions:

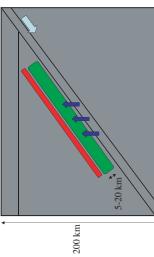
GyPSM-S is a powerful, flexible technique that allows for the discovery of new phenomena that can only be observed within the setting of a fully-coupled model integrating both chemistry and dynamics. The presence of the low-viscosity channel is not a transient feature within subduction zones, being defined by the water release from dehydration reactions and by the partially-hydrated peridotite solidus. It has dynamic implications for evolution of slab dip as the slab is allowed to decouple from the mantle wedge.

References:

- Smith & Asimow (2003); Hacker et al. (2003); Billen & Gurnis (2001); Hirth & Kohlstedt (1996); Kring et al. (1990); Mosenfelder et al. (2006); Werckman & Hart (2003); Ghiorso & Sack (1995)



Low-Viscosity Channel:



Dehydration reaction within the downgoing slab (at the altered basalt and serpentinite layers) release free water (Hacker et al., 2003), which flows into the overlying mantle wedge and reacts with the peridotite to form hydrous phases, and to hydrate nominally-anhydrous minerals, such as olivine.

The influence of elevated hydrolytic dehydrations in lithology on the mantle wedge to produce accretionary prism is accounted for in the coupled formulation (Herr & Kohlstedt, 1996). The relatively thin layer of hydrated nominally-anhydrous minerals is bounded by the slab and the partially-hydrated solidus.

Water rise from the slab through the hydrated olivine layer (green), and into serpentinite layers release free water (red). The relatively thin layer of hydrated nominally-anhydrous minerals is bounded by the slab and the partially-hydrated solidus.

Water rise from the slab through the hydrated olivine layer (green), and into serpentinite layers release free water (red). The relatively thin layer of hydrated nominally-anhydrous minerals is bounded by the slab and the partially-hydrated solidus.

Water rise from the slab through the hydrated olivine layer (green), and into serpentinite layers release free water (red). The relatively thin layer of hydrated nominally-anhydrous minerals is bounded by the slab and the partially-hydrated solidus.

Elastic versus permanent deformation above the Sunda megathrust

Rich Briggs¹, Willy Amidon¹, Kerry Sleaf¹, John Galvez¹, Bambang Suwargadi², Danny Natawardja², Nugraha Sastra², Dudi Prayud²

¹ Scripps Observatory, Division of Geodesy and Remote Sensing, California Institute of Technology

² Research Center for Geotechnology, Indonesian Institute of Sciences (LIPI), Bandung, Indonesia

Informed by the National Science Foundation

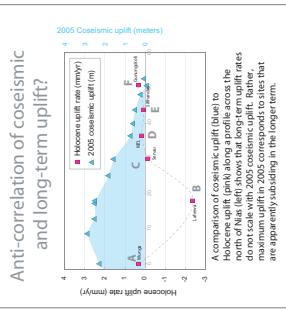
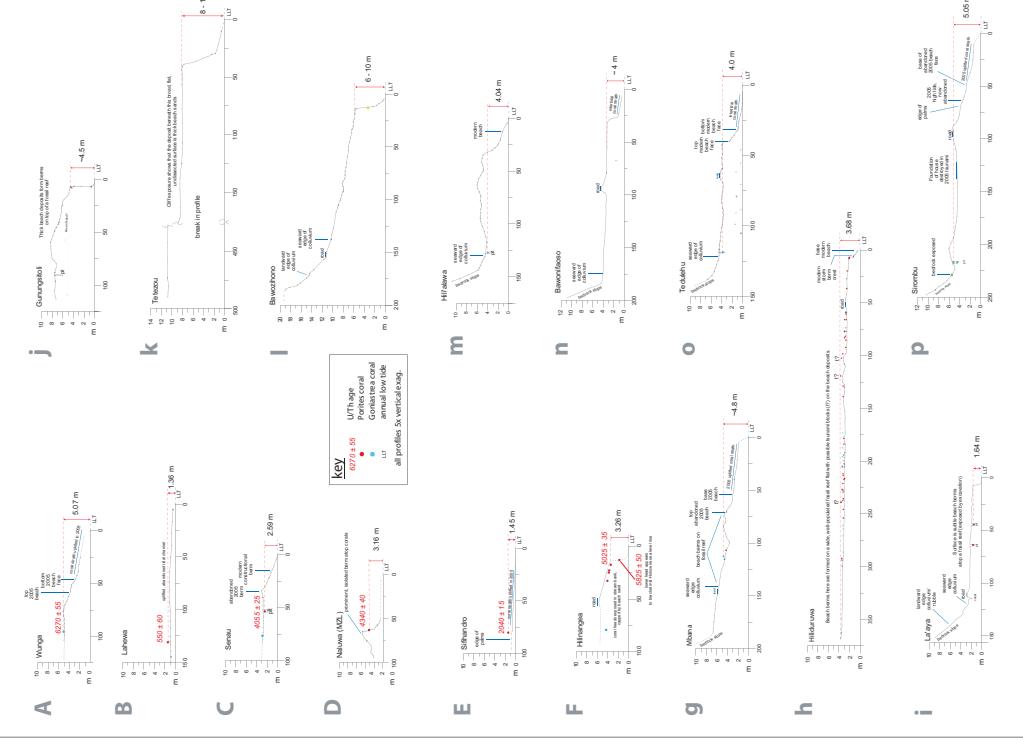
and grant from The Gordon and Betty Moore

Foundation to the Tectonic Observatory, California

Institute of Technology

Topographic profiles of uplifted Holocene surfaces

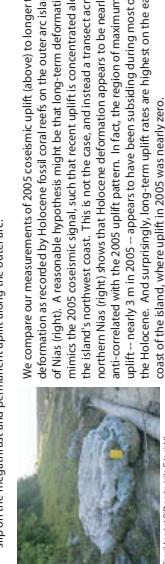
We compare uplifted fossil coalheads, which are reliable recorders of paleo-sea-level, to present-day low tide to determine total apparent uplift. From this raw number we subtract 2005 coseismic uplift, and also the appropriate mid-Holocene sea level transgression to calculate raw uplift rates. All profiles below show raw uplift values. Locations on map at left.



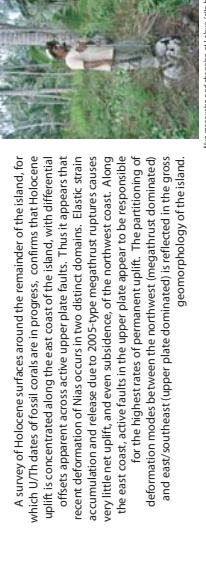
Location of study area and contour map of 2005 coseismic uplift. Most of Nias island is the west of the island – 2.2 to 3.5 m – runs along the NW coast, while the line of zero uplift is nearly aligned along the east coast.

Summary

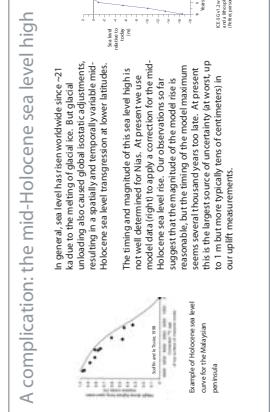
The 2005 Sunda megathrust rupture provides a rare opportunity to explore the relationship between megathrust slip and permanent uplift along the outer arc.



We compare our measurements of 2005 coseismic uplift above to longer term deformation as recorded by Holocene fossil coal records on the outer arc island of Nias (Fig.). A reasonable hypothesis might be that long-term deformation mimics the 2005 coseismic signal, such that recent uplift is concentrated along the island's northwest coast. This is not the case, and instead a transect across the northern part of Nias shows that Holocene deformation appears to be nearly anti-correlated with the 2005 uplift pattern. In fact, the region of maximum uplift – nearly 3 m in 2005 – appears to have been subsiding during most of the Holocene. And surprisingly, long-term uplift rates are highest on the east coast of the island, where uplift in 2005 was nearly zero.



Nias measuring and choosing a line of site to go!

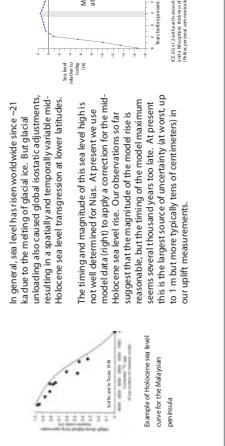


A correlation between megathrust slip and basin has long been noted (e.g. Logg, 1999) and more recently Wei et al., 2002, and Logg and Simons, 2003 find that during the long interseismic period of the megathrust, the basin subsides above the future megathrust rupture, and an accumulation of sediments occurs. This is a significant portion of plastic, non-conservative deformation takes place during the megathrust cycle. However, the model does not account for the megathrust cycle. Alternatively, some type of regional sea level solution. Alternatively, some type of basin erosion may take place coseismically (e.g. von Heune and Scholl, 1991). At right is schematically shown the possible interplay of shortening a long brittle upper, but structures and surface controlled by the earthquake. See the next paragraph.

Do paleo-coal measures explain the pattern of subsidence?

A correlation between megathrust slip and basin has long been noted (e.g. Logg, 1999) and more recently Wei et al., 2002, and Logg and Simons, 2003 find that during the long interseismic period of the megathrust, the basin subsides above the future megathrust rupture, and an accumulation of sediments occurs. This is a significant portion of plastic, non-conservative deformation takes place during the megathrust cycle. However, the model does not account for the megathrust cycle. Alternatively, some type of regional sea level solution. Alternatively, some type of basin erosion may take place coseismically (e.g. von Heune and Scholl, 1991). At right is schematically shown the possible interplay of shortening a long brittle upper, but structures and surface controlled by the earthquake. See the next paragraph.

A complication: the mid-Holocene sea level high



In general, the megathrust is responsible for a 21 cm of relative sea level rise per year due to a combination of tectonic subsidence and eustatic sea level rise.

The timing and magnitude of this sea level high is not well determined for Nias. A transient sea level model data fit right to apply a correction for the mid-Holocene sea level rise. Our observations so far suggest that the amplitude of the model rise is consistent with the one occurring during the megathrust cycle. Alternatively, some type of regional sea level solution. Alternatively, some type of basin erosion may take place coseismically (e.g. von

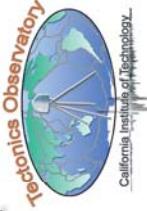
Heune and Scholl, 1991). At right is schematically shown the possible interplay of shortening a long brittle upper, but structures and surface controlled by the earthquake. See the next paragraph.



The origin of high-Ca annuli in garnets from the Rand schist of the San Emigdio Mountains, southern California

Alan D. Chapman, Peter Luffi, and Jason B. Saleeby

Division of Geological and Planetary Sciences, California Institute of Technology, Pasadena, CA



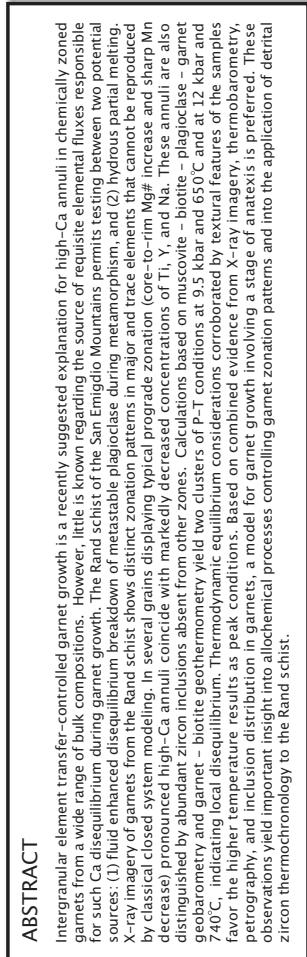
ABSTRACT

Intergranular element transfer-controlled garnet growth is a recently suggested explanation for high-Ca annuli in chemically zoned garnets from a wide range of bulk compositions. However, little is known regarding the source of reusable elemental fluxes responsible for such Ca disequilibrium during garnet growth. The Rand schist of the San Emigdio Mountains permits testing between two potential sources: (1) fluid enhanced breakdown of metastable plagioclase during metamorphism, and (2) hydrous partial melting. X-ray imagery of garnets from the Rand schist shows distinct zonation patterns in major and trace elements that cannot be reproduced by classical closed-system modeling. In several grains displaying typical prograde zonation (core-to-rim $Mg\#$ increase and sharp Mn decrease) pronounced high-Ca annuli coincide with markedly decreased concentrations of Ti, Y, and Na. These annuli are also distinguished by abundant zircon inclusions absent from other zones. Calculations based on muscovite – biotite – plagioclase – garnet geobarometry and garnet – biotite geothermometry yield two clusters of P-T conditions at 9.5 kbar and 650°C and at 12 kbar and 740°C, indicating local disequilibrium. Thermodynamic equilibrium considerations corroborated by textural features of the samples favor the higher temperature results as peak conditions. Based on combined evidence from X-ray imagery, thermobarometry, petrography, and inclusion distribution in garnets, a model for garnet growth involving a stage of anatexis is preferred. These observations yield important insight into abiochemical processes controlling garnet zonation patterns and into the application of deritified zircon thermochronology in the Rand schist.

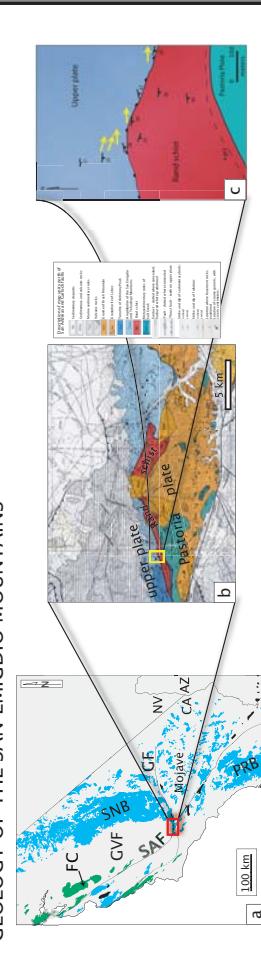
Intergranular element transfer-controlled garnet growth is a recently suggested explanation for high-Ca garnets from a wide range of bulk compositions. However, little is known regarding the source of requisite elements for such Ca disequilibrium during garnet growth. The Rand schist of the San Emigdio Mountains permits test sources: (1) fluid enhanced disequilibrium breakdown of metastable plagioclase during metamorphism, and (2) X-ray imagery of garnets from the Rand schist shows distinct zonation patterns in major and trace elements by classical closed system modeling. In several grains displaying typical prograde zonation (core-to-rim Mn decrease) pronounced high-Ca annuli coincide with markedly decreased concentrations of Ti, Y and Zr, distinguished by abundant zircon inclusions absent from other zones. Calculations based on muscovite - biotite geobarometry and garnet - biotite geothermometry yield two clusters of P-T conditions at 9.5 kbar and 740 °C, indicating local disequilibrium. Thermodynamic equilibrium considerations corroborated by texture favor the higher temperature results as peak conditions. Based on combined evidence from X-ray imaging, petrography, and inclusion distribution in garnets, a model for garnet growth involving a stage of anatexis observations yield important insight into allochemical processes controlling garnet zonation patterns and interzircon thermochronology to the Rand schist.

BACKBOUND: DISCUSSION ELEMENT INCORPORATION MODELS

thermodynamic models of rocks with MORB composition (Konrad-Schmolle et al., 2005) and garnet zoning patterns from micaeous quartzites (Chemoff and Carlson, 1997) show evidence for disequilibrium incorporation of Ca during garnet growth. These studies show that high-Ca annuli in garnets can be produced during sporadic porphyroblast growth in localized chemical inhomogeneities (figure 1). Zoning profiles from garnets found in the Rand schist of the San Emigdio Mountains, southern California, display similar zonation patterns. However, the observed Ca zoning patterns are ubiquitous at the outcrop scale and coincide with pronounced fluctuations in trace element concentrations. Open system processes such as fluid-fluxed partial melting and/or catastrophic breakdown of Ca-silicate minerals are the common mechanism for Ca zoning in garnets.



SOCIETY FOR THE STUDY OF LITERATURE



Sample	Unit
04SE14	upper pl.
04SE16a	upper pl.
04SE10	Rand sc.
04SE9	Rand sc.
04SE23	Rand sc.

卷之三

If the Rand schist of the San Emigdio Mountains underwent anatexis, melts may have lubricated the margins of the schist and facilitated its exhumation from depth.

CONCLUSIONS

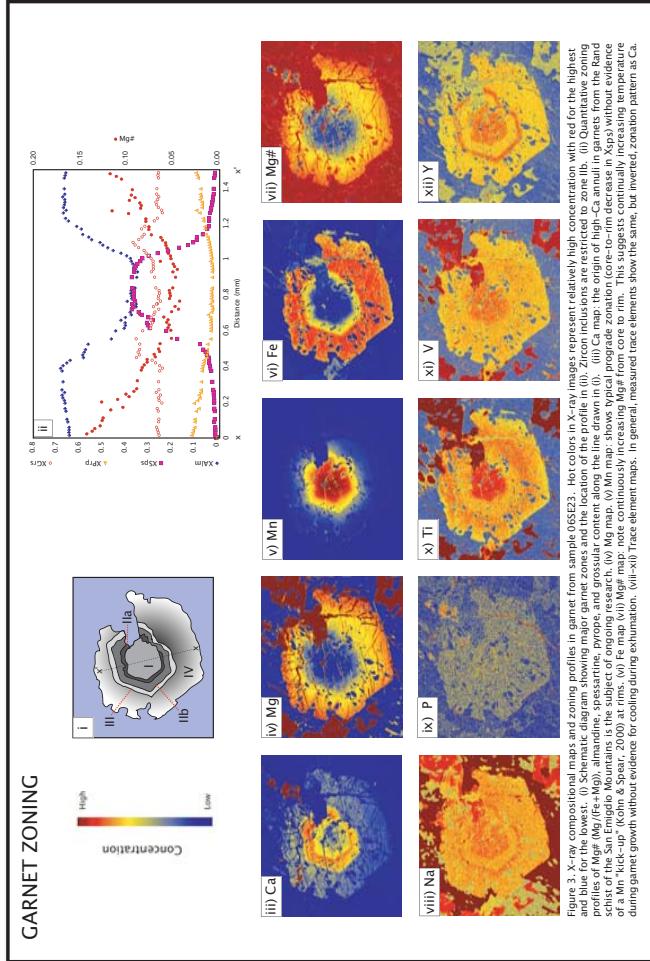
- Garnet growth history
X-ray imagery, thermobarometry, petrography, and zircon inclusion distribution in garnets suggest a model for garnet growth involving a stage of hydrous partial melting.

ELITI BE WORK

- Zr-in-utile thermometry** to correctly interpret the thermal evolution of the schist, it is critical to have an independent check on the peak temperatures experienced by the schist (figure 4).



same as in figure 3. Temperature has a strong positive correlation with the Zr content in rutile. Note a narrow zone of high Zr near the rim of each rutile grain. This zone may have been the result of a late stage of heating.





Waveform Modeling of the Slab underneath Japan

Min Chen, Jeroen Tromp, Don Helmberger & Hiroo Kanamori
Seismological Laboratory, California Institute of Technology, mchen@gps.caltech.edu

Data selection and model testing

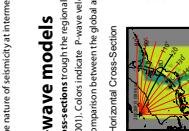
Figure 4. 3-component seismograms from more than 600 Japan-Hinet stations produced by two earthquakes in the Japan subduction zone, one which occurred in the downgoing Pacific plate at depths greater than 400 km. We simulate subducting peridotite in the 3-D P-wave model (Zhao et al. 1994) using 2-D finite-difference (FDM) and 3-D spectral-element (SEM) methods.

As illustrated by cross-correlation between synthetics and data (Zhao et al. 1994), typically explained by the waveform anomaly of the interface and some of the waveform complexity, but fail to predict the extended Shallow event. In this study we take advantage of the densely distributed Hinet stations and use 2-D FDM modeling to simulate the S-wave and SH-wavefronts. Our 2-D model suggests a thin elongated low-velocity layer (LVL) exists above the slab, extending down to a depth of 300 km with an S-wave velocity reduction of 4% if the thickness of 20 km is assumed. Furthermore, the water-reddish phase transition of ultra-mafic rocks, and the nature of seismicity at intermediate depths (>70–300 km).

3-D P-wave models

Figure 1 shows cross-sections through the regional and global tomographic P-wave models derived by Zhao et al. (1994) and Zhao (2001). Colors indicate P-wave velocity anomalies relative to a 1-D Earth model (ASPEP91). (a) The regional model (b) Global model.

Horizontal Cross-Section



3-D Spectral-Element mesh

Figure 2 shows velocity anomalies are superimposed on the one-chunk SEM mesh. For parallel computing purposes, the one-chunk SEM simulations are subdivided in terms of 20-slices. The center of the chunk is at 33.8°N, 137.5°E, and the lateral dimensions are 30' × 30'. (a) Full view of two neighboring slices. (b) Close-up view of the upper mantle mesh.

Events and Stations

Figure 3 is a map view of the study area. Hinet stations are indicated by red triangles and contours of the upper plate boundary of the Pacific are indicated by black lines. The location of event 20020915 is marked by a red star and event 20030331 is indicated by the green star.

FDM simulations

Figure 11. Base model with a slab inside the transition zone (a) 2D snapshots of waveforms and vertical cross-correlation. (b) Vertical S-waves. (c) Radial S-waves. (d) Vertical S-waves. Preferred scaling factor $f = 1.5 - 2.0$.

Figure 12. Slabmodels with a low-velocity layer (LVL) (a) LVL (80%; DL: HJ). Depth (km)

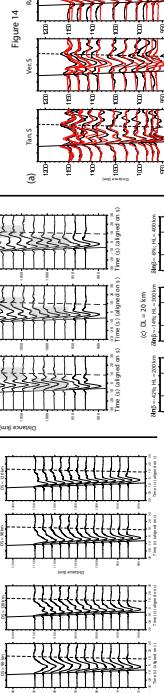


Figure 13. Three models with different types of mantle wedges

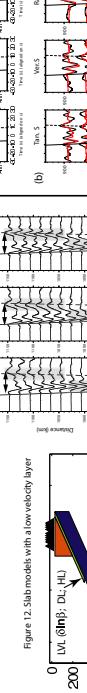


Figure 14. Three-component S-waveforms and vertical S-waveforms comparison between data (black lines) and 3-D SEM synthetic (red lines). SEM synthetics are calculated for Model 2 in Figure 13. (a) Event data and synthetics for 5 waves as it is filtered between 0–29 s, and for P-waves between 3–25 s. (b) Event 20020915 depth 589 km.

(c) Event 20030331 (depth 492 km). Model 2 is our preferred model and fits the data for both events on all three components adequately.

Figure 15. 2-D FD snapshots of SH-wave propagation.

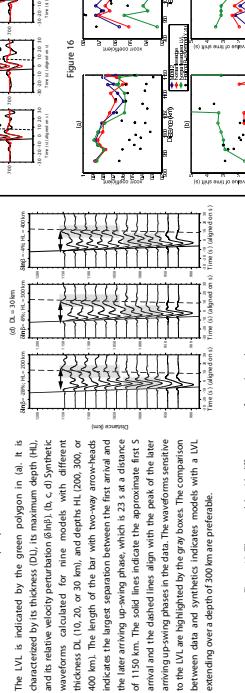


Figure 16. Comparison of cross-correlation coefficients (upper panel) and traveltimes anomalies (bottom panel) between data and SEM synthetics for four different models.

Black star: Reginald model.

Red circle: Model 2.

Blue circle: Model 2 without a slow mantle wedge but with a LVL.

Green circles: Model 2 without a slow mantle wedge and a LVL.

Figure 17. Comparison of cross-correlation coefficients (upper panel) and traveltimes anomalies (bottom panel) between data and SEM synthetics for four different models.

Black star: Reginald model.

Red circle: Model 2.

Blue circle: Model 2 without a slow mantle wedge but with a LVL.

Green circles: Model 2 without a slow mantle wedge and a LVL.

Conclusion The 2-D slab model indicates there is an elongated low-velocity layer above the slab extending down to a depth about 300 km, around 20 km thick. However, the thickness of the LVL trades off to some extent with the low-wave velocity in the LVL. We compare the normal mode traveltimes of the LVL beneath the Japan to be composed of hydrated mafic and/or titanite rocks – above a depth of 150 km the LVL could be composed of hydrous olivine and/or serpentinized peridotite above and/or below the descending crust; below a depth of 150 km this hydrous layer is more likely composed of serpentized peridotite (or at the greatest depth, phengite) above or below the fully eclogitized oceanic crust. Water released from the dehydration reactions in this hydrous zone could cause the abundant arc volcanism, the intermediate intra-slab seismicity (70–300 km), and possible silent epis events, which have been observed in other subduction zones.



Body Wave Attenuation Structure in Southern Mexico

Ting Chen¹, Robbert W. Clayton¹
¹ California Institute of Technology, Seismological Laboratory 252-21, Caltech, Pasadena, CA 91125

Abstract

The 2-D model of attenuation structure is determined along a 250 km trench-perpendicular profile in southern Mexico. Velocity spectra from 14 moderate earthquakes recorded by the Meso American Subduction Experiment (MASE) array which consists of 100 broadband sensors from Acapulco to Tampico are used. By assuming a Brune-type source, a path-averaged frequency-independent Q is obtained for each seismogram in the frequency band 0.5 Hz to 7.30 Hz depending on the signal quality. These measurements are then inverted for spatial variations in Q . The 1-D tomography result shows a pattern of Q qualitatively similar to other subduction zones, with low attenuation crust ($Q \sim 1000$), and high attenuation in the mantle wedge beneath the Trans-Mexico-Volcanic-Belt ($Q < 250$). The location of the low- Q region and the variation of the Q value also provides some constraints on the geometry of the subducting slab, or with the structure provided by other methods such as receiver functions, the Q estimates can be used to estimate variations in viscosity.

Data and Analysis

We used moderate earthquakes with magnitudes ranging from 4.5 to 5.0. The dataset is shown in Figure 1, and consists of 14 local earthquakes recorded by the instruments of MASE from February 2005 to May 2006.



The smoothed spectra of the vertical velocity-component for P wave is calculated from a 3.5 s time window beginning 0.5 s before the arrival pick, after 5% cosine taper is applied. A smoothed spectra of noise is also calculated from a 3.5 s time window immediately preceding each signal window in the same way. Tests show that changing the window length from 2–5 s does not produce significantly different spectra in the frequency band used, so a constant 3.5 s window was used for all the events. The signal is kept for further analysis if the signal-to-noise ratio is greater than 2, in a frequency band 0.5 to 7.30 Hz.

Method

Spectral Analysis

The Fourier velocity spectral amplitude of a body wave from event j , recorded at station i , can be written as [e.g. Garcia, 2004]

$$A_i(f_j) = S_i(f_j) / (N_i(t_{\text{wave}} - t_j)) \quad (1)$$

where $S_i(f)$ is the source spectra, $N_i(t)$ is the instrument response, C is the frequency-independent amplitude term associated with geometric spreading, seismic moment radiation pattern, and other static effects. The exponential term describes the attenuation effect. The term t^* can be expressed as $t^* = t/Q$, where t is the travel time, and Q is the quality factor.

Assuming a Brune-type source [Brune, 1970], the source velocity spectrum of event j , can be written as

$$S_i(f_j) = \frac{M_0}{1 + (f/f_c)^2} \quad (2)$$

where M_0 is the signal moment, and f_c is the corner frequency.

Since only data in the flat portion of the pass band of the recording system is used in this study, t^* can be neglected. The corner frequency is estimated based on the result given by Daniel Garcia (2004) for central Mexico, which takes the form of

$$f_c = 350 \times 10^{-12} M_0^{1/3} \quad (3)$$

where M_0 is the seismic moment in dyne-centimeters. This result is actually for shear waves, so (3) needs to be modified in this study by taking the P-S wave corner frequency ratio to be 1.5 [Boatwright, 1985], which means the corner frequency for P wave has the form of

$$f_c = 2.39 \times 10^{-12} M_0^{1/3} \quad (4)$$

Now we can rewrite (1) as

$$\log A_i = \log C + \log f - \log t^* - \log \left(\frac{M_0}{1 + (f/f_c)^2} \right) \quad (5)$$

where C is a frequency-independent term. Taking the logarithm of equation (5), we obtain

$$\log A_i = \log C + \log f - \log t^* - \log \left(\frac{M_0}{1 + (f/f_c)^2} \right) \quad (6)$$

We then solve this equation by a least-squares method to obtain the average frequency-independent Q for each ray path.

Tomographic Inversion

One-dimensional tomographic inversion is used to determine the attenuation structure. Based on the data coverage, the study region is divided into eight blocks parallel to the trench, and each block is assumed to have constant Q . The observed t^* for the i th ray path is

$$t_i^* = \sum_j T_{ij} Q_j \quad (7)$$

where T_{ij} is the travel time in block i for the i th ray, and Q_j

is the quality factor of block j . To compute the travel time, we currently assume that the ray path is a straight line connecting the source and station.

The inversion problem can be written in matrix form as

$$[\beta_{ijkl}] Q_j = [\tau_{ikl}] \quad (8)$$

where M is the number of blocks and N is the number of data.

Summary

We have studied the attenuation structure in southern Mexico using the spectral decay method. The results show a low- Q zone in the last block (Figure 5). The low- Q blocks correspond to the Trans-Mexico-Volcanic-Belt on the map. We interpret the low- Q region as the mantle wedge as has been shown in attenuation studies of other subduction zones.

Our future work is to obtain a more detailed 3-D tomographic attenuation model, which takes Q variation in depth and varying subducted slab geometry into consideration. This work may involve earthquake relocation and ray tracing. Then, we can convert the attenuation model into viscosity model by using the approximation $\eta = \pi Q / g^2$ [Billett and Gurnis, 2001]. The viscosity value obtained will help us to conduct a more realistic geodynamic modeling in southern Mexico.

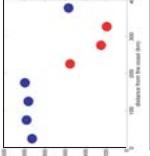


Figure 4

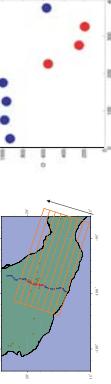


Figure 5



Figure 6

Inversions show high Q in the first few blocks, low Q in the next three blocks, and a slightly higher Q again in the last block (Figure 5). The low- Q blocks correspond to the Trans-Mexico-Volcanic-Belt on the map. We interpret the low- Q region as the mantle wedge as has been shown in attenuation studies of other subduction zones.

Our results indicate that the mantle wedge begins about 200 km, and ends about 350 km from the coast, and gradually transitions to normal upper mantle wedge beyond that (Figure 6).

Figure 2



In Figure 2, two velocity spectra are calculated for two paths with the same distance from an event. The blue one is mainly through the crust, and the red one is mainly through the mantle wedge. From figure 2, we can see that the wave passing through the mantle wedge attenuates more than the wave passing through the crust.

A complete set of path-average Q determined for one event is provided in Figure 3 for example.

We can see that the average Q for the wave path through the mantle wedge is smaller than that for the path through the crust. This indicates that the mantle wedge is characterized by higher attenuation property than the crust.

Our future work is to obtain a more detailed 3-D tomographic attenuation model, which takes Q variation in depth and varying subducted slab geometry into consideration. This work may involve earthquake relocation and ray tracing. Then, we can convert the attenuation model into viscosity model by using the approximation $\eta = \pi Q / g^2$ [Billett and Gurnis, 2001]. The viscosity value obtained will help us to conduct a more realistic geodynamic modeling in southern Mexico.

Figure 3

Coseismic Slip and Afterslip of the Great (Mw9.15) Sumatra-Andaman Earthquake of 2004.

Mohamed Chlieh¹, Jean-Philippe Avouac¹, Vala Hjorleifsdottir¹, Teh-Ru Alex Song¹, Chen Ji^{1,2}, Kerry Sieh¹, Anthony Sladen³, Helene Hebert³, Linette Prawirodirdjo⁴, Yehuda Bock⁴, and John Galetzka¹

¹Tectonics Observatory, Division of Geological and Planetary Sciences, California Institute of Technology, Pasadena, CA 91125, USA.

²Laboratoire de Détection et de Géophysique, CEA, Bruyères-le-Châtel, France

³Cecil H. and Ida M. Green Institute of Geophysics and Planetary Physics, Scripps Institution of Oceanography,

University of California San Diego, La Jolla, CA 92093, USA.

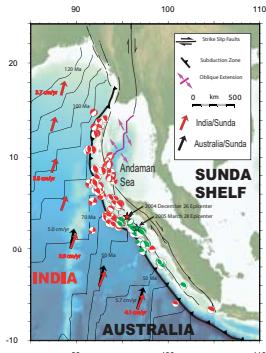
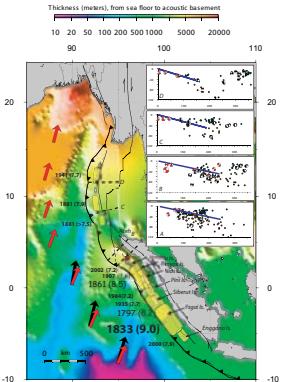
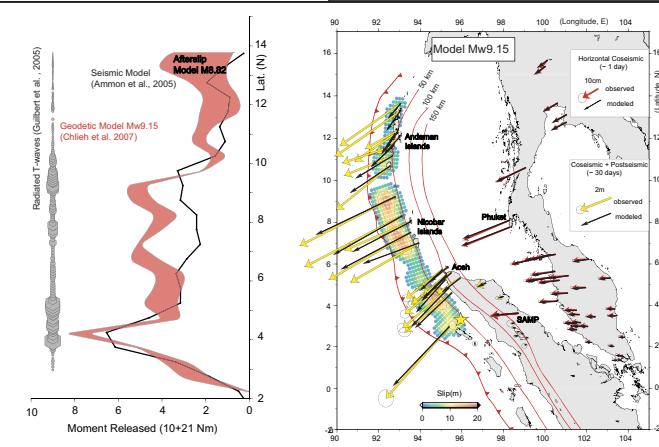


Figure 1: Neotectonic setting of the great Sumatra-Andaman earthquake. The map shows the Sunda Shelf, Andaman Sea, India, and Australia. The age of the sea floor [Cande and Kent 1995] and the location of the megathrust [Bock et al., 2003] are shown. The red star indicates the epicentral area to 90 Ma near Andaman Islands. The red star indicates the location of the 2004 Sumatra-Andaman megathrust and its aftershocks. The green star indicates the epicenter of the 28 March, 2005, Nias Mw6.7 earthquake. CMT associated to the aftershocks of the 2004 Sumatra-Andaman megathrust is in red and to the aftershocks of the 2005 Nias earthquake is in green.



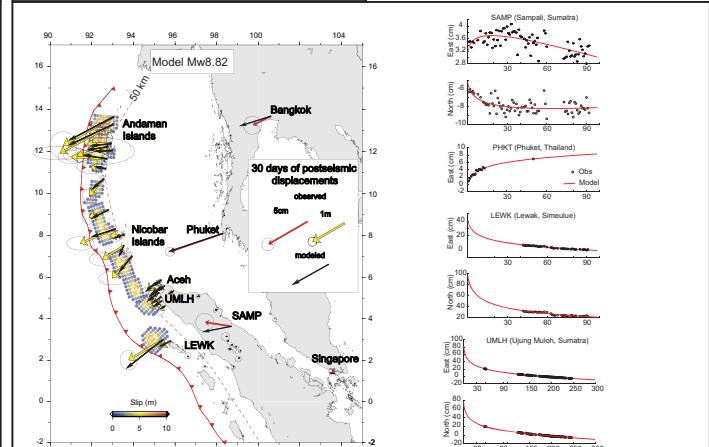
Estimated region area of the major historical earthquakes along the Sunda subduction zone between 1833 and 2004 [Vernon and McCann, 1987; Zachariasen et al., 1999; Natuwidja et al., 2004]. The background shows the sediment thickness from sea floor to acoustic basement [Held et al., 2000]. Cross sections of coseismic, relocated seismicity and CMT solutions of aftershocks.

Coseismic Slip



Latitudinal variations of scalar moment as determined from seismic model Mw9.15 (Kagan et al., 2003) and geodetic model Mw9.15 coseismic slip model Mw9.15. Both our model and model derived from the latitudinal variation of radiated energy by T-waves (Guillet et al., 2005) show three distinct peaks. The third peak corresponds to the relative amplitude released since the earthquake by hydrophone sensors at Diego Garcia island.

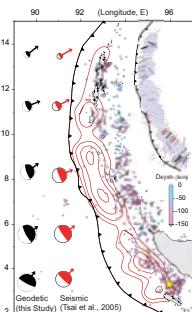
Afterslip



Continuous GPS time series and best-fitting analytical function corresponding to instantaneous slip (red line). Average displacement per day was determined from the best fit to 5 minutes (SAMP) and Phuket (PHKT) time series and applied to fit the Ujung Muluh (UMLH) and Lewak (LE WK) time series.

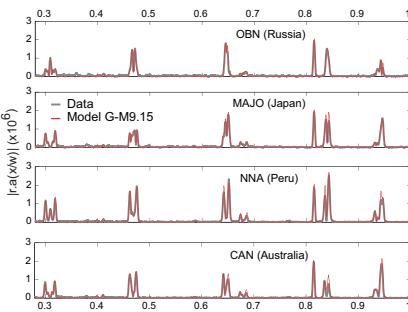
Consistencies of the Coseismic Slip model with:

CMT



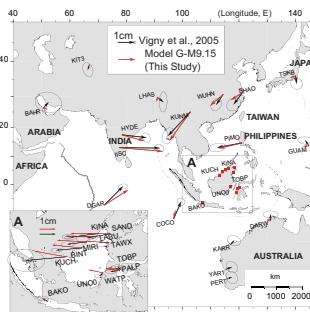
Summary of slip distribution on the megathrust during the 2004 earthquake. The coseismic Mw9.15 slip contours are each 5 cm increments. Most of the 2004 coseismic slip occurred trenchward of the megathrust. The upper plot shows the trenchward component of the coseismic double couple mechanism for five regions of the rupture for 2004. The lower plot shows the total coseismic double couple mechanism for 2004, based upon centroid-moment-tensor (CMT) analysis of the earthquake. The upper inset shows the azimuth of slip on the megathrust (0° to 180°) and the lower inset shows the azimuth of slip on the far-field north. Bars are the slip vector azimuths of shallow dipping aftershocks and aftershocks from the Harvard CMT catalog.

Normals Modes



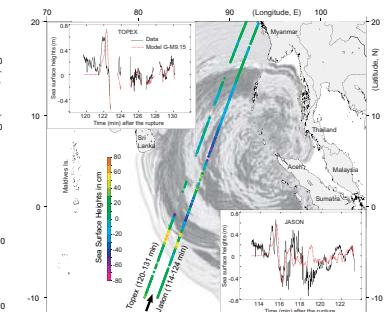
Comparison between normal modes data (thick line) and geodetic model G-M9.15 synthetics at four stations: OBN (Obninsk, Russia), MAJO (Matsushiro, Japan), NNA (Nana, Peru), CAN (Canberra, Australia).

Very Far-Field



Comparison of the very far-field GPS (Vigny et al., 2005) and predicted displacements of model M9.15 computed by SEM between 2000 and 4000 seconds. This model accounts for 3D structure (Model Crust 2.0, [Bassin et al., 2000]), ellipticity, gravity, rotation, topography and ocean load.

Tsunami



Sea surface heights observed by the JASON and TOPEX-Poseidon satellites compared to numerical simulations of the tsunami based on the coseismic model M9.15 (dotted lines).

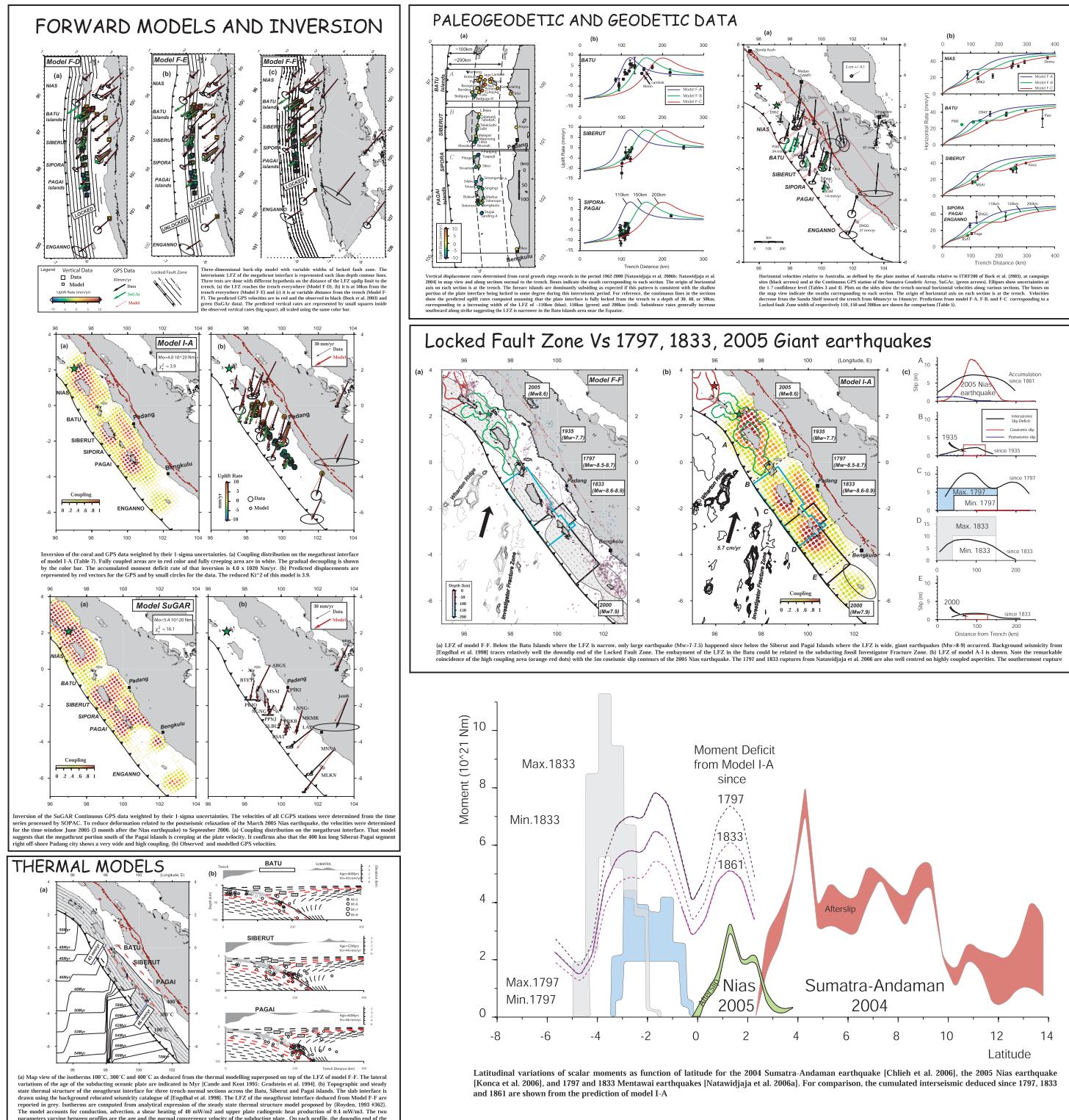
The model that fits best the geodetic measurements recorded within the first day of the 2004 earthquake is M9.15. This model is consistent with seismological, tsunami and T-waves observations. We deduce that the seismic rupture must have propagated as far as 15°N. The latitudinal distribution of moment in the model has three distinct peaks. This pattern is consistent with latitudinal variations in energy released by T-waves and high-frequency diffracted seismic waves. The general pattern in the model is a gradual northward decrease in slip. The fact that this mimics the northward decrease of the convergence rate across the plate boundary suggests that this pattern might be a characteristic feature of the large ruptures along this stretch of the megathrust.

Although our data place only low constraints on slip near the trench it seems that the coseismic rupture didn't reach to the trench everywhere. This inference is based on the slip distribution obtained from the inversion of the geodetic data and the consistency of that model with the amplitude of the deep-sea tsunami wave. Possibly that would reflect the effect of the poorly lithified sediments at the toe of the accretionary prism on the rheology of the plate interface, which would have inhibited the propagation of the seismic rupture due to a rate-strengthening friction mechanism [Byrne et al. 1988; Scholz 1998]. If this is so, one would expect afterslip on the megathrust proximal to the trench in response to stresses induced by the coseismic rupture [Marone et al. 1991]. A model of frictional afterslip explains to first order the evolution of postseismic deformation. Within 60 days of the earthquake, post-seismic moment release equaled about 35% of the coseismic moment, the equivalent of an Mw 8.82 earthquake. The ratio of coseismic to postseismic slip is higher than this average north of 11°N. In fact, afterslip in this portion of the megathrust in the month following the earthquake was larger than the coseismic slip. Perhaps this is evidence that the rheology of the megathrust there is strongly influenced by subduction of the exceptionally thick sedimentary sequence of the Bengal fan. Although the spatial distribution is poorly resolved, afterslip seems to have occurred over about the same width of the megathrust as coseismic slip.

Geodetic and paleogeodetic resolution of locked patches on the Sunda megathrust, offshore Sumatra

Mohamed Chlieh, Jean-Philippe Avouac, Kerry Sieh,
Danny H. Natawidjaja, John Galetzka,
Tectonic Observatory, Caltech, Pasadena, CA
Research Center for Geotechnology, Bandung, Indonesia

Geodetic and paleogeodetic measurements of strain above the Sumatran portion of the Sunda subduction zone reveal a heterogeneous pattern of coupling along the subduction megathrust. Annual banding in coral heads provides vertical rates of deformation spanning the last half of the 20th century and repeated GPS surveys between 1991 and 2004 and continuous measurements at GPS stations operated since 2002 provide horizontal velocities. The area of the plate interface within which the coupling is high is only ~75 km wide near the Equator but increases to ~175 km farther south. Major sections of this Locked Fault Zone (LFZ) coincide with the rupture areas of major Mw>8.5 interplate earthquakes. The section that ruptured during the Mw 8.7 Nias earthquake of 2005 released about 2/3 of the slip deficit that had accumulated since its previous rupture in 1861. Farther south, beneath the Mentawai islands, overlapping ruptures of the LFZ produced giant earthquakes in 1797 and 1833. The accumulated slip deficit since these events is slowly reaching the amount of slip that occurred during the 1833 earthquake but already exceeds the slip that occurred during the 1797 earthquake. Thus, re-rupture of the Mentawai patch in the near future seems quite likely. In contrast, coupling is low in the Batu islands near the Equator and around Enganno island at about 5S, where only moderate earthquakes have occurred in the past two centuries. Temperature might influence the mode of slip along the plate interface, through its effect on the rheology of sediments at the plate interface. Other influences, such as structures on the subducting plate, could also play a role. In particular, subduction of the Investigator Fracture Zone near the Equator coincides with the relatively low coupling there.



MASE: Shallow Subduction in Central Mexico

Robert W. Clayton, Xyoli Pérez-Campos, Paul Davis, Arturo Iglesias, YoungHee Kim, Ting Chen, Alan Husker, Fernando Greene, Lizbeth Espejo, Luca Ferrari, Dante Moran, John Eiler, Mike Gurnis, Vlad Manea, Carlos Valdez, Joann Stock, Vladimir Kostroglodov

ABSTRACT

The objective of the MASE (Middle America Subduction Experiment) is to construct a geodynamical model of the subduction process. The Middle America Trench was chosen as the first example because of the relatively simple plate geometry (a linear margin with near normal subduction) and a significant along-strike slab-dip variation. The initial deployment along the Acapulco to Tampico transect in central Mexico is designed to investigate the case of shallow subduction.

The main results to date are:

- The discovery that the slab underplates the continental crust to a distance of 200 km from the trench. This result is interesting because there is no geologic or geodetic indication of coupling in this zone – the coupling that is measured geodetically is confined to the initial 75 km near the coast where the slab is dipping down. There is also no fluid signature in the magnetotelluric (MT) data of the flat-slab portion of the line.
- The relative attenuation in the mantle under the Mexican Volcanic Belt (presumed location of the mantle wedge) is a factor of 2 higher than the surrounding mantle.
- Modeling studies indicates that a shrinking low-viscosity mantle wedge can lead to flat-slab subduction as observed.
- A slow earthquake appears to be in progress on the southern 200 km of the MASE line. The last slow event occurred in 2002. This one started in March, 2006.

El objetivo del proyecto MASE (*Middle America Subduction Experiment*) es construir un modelo geodinámico que represente un proceso de subducción. Se ha escogido como primer ejemplo la Trinchera de América Central debido a que presenta una geometría relativamente simple (una maraña con geometría lineal y con una dirección de convergencia cercana a la normal) y una variación significante del manto de la placa a lo largo del rumbo. La fase inicial de este proyecto, que consiste en un experimento a lo largo de un transecto localizado en la región central de México entre Acapulco y Tampico, se ha diseñado para investigar este tipo de subducción subhorizontal. Los principales resultados obtenidos hasta el momento son:

- El descubrimiento de que la placa se encuentra en contacto con la corteza continental hasta una distancia de 200 km de la trinchera. Este resultado es interesante dado que no se ha encontrado indicación geológica o geodésica de que exista acoplamiento en esta zona – el acoplamiento medida en forma geodésica está confinado a los primeros 75 km cercanos a la costa, donde la placa se encuentra buceando. En la porción del transecto que pasa por la zona de subducción subhorizontal, los datos magnetotelúricos (MT) no indican la presencia de fluidos.
- Bajo el cinturón volcánico Mexicano (donde se piensa que está ubicada la curva del manto) el manto tiene una atenuación relativa con el doble del valor de la del manto circundante.
- Estudios a través de modelos numéricos, indican que la existencia de una cuña de baja viscosidad en el manto puede ser la causa de la presencia de zonas de subducción subhorizontal como la que se observa en este experimento.
- A partir de Marzo de 2006, se observa que un terremoto lento se está desarrollando en los 200 km ubicados en la zona sur del transecto de MASE. El último evento de este tipo ocurrió el año 2002.

CURRENT MASE SEISMIC ARRAY



RECEIVER FUNCTION STUDY

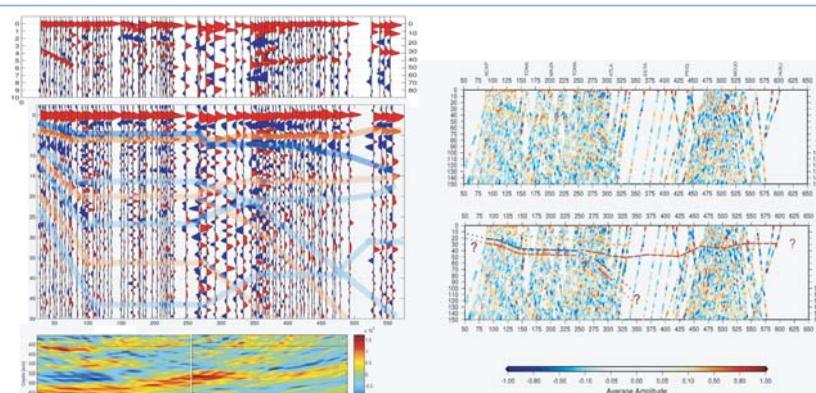
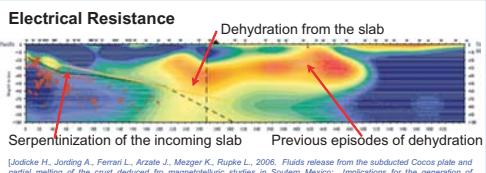


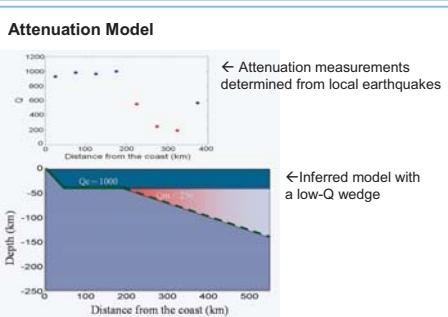
Image by back-projection and migration of receiver function. Top panel is the image and the bottom panel is current interpretation.

MAGNETOTELLURICS STUDY

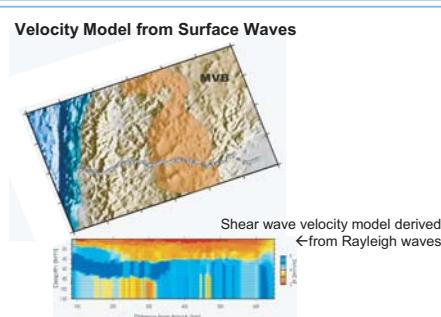


[Jodicke H., Jonck A., Ferrari L., Kruse J., Meier K., Rupke L., 2006. Fluid release from the subducted Cocos plate and partial melting of the crust deduced from magnetotelluric studies in Staudem Mexico. Implications for the generation of volcanism and subduction dynamics. *Journal of Geophysical Research*, v. 111, B08102, doi:10.1029/2005JB003739.]

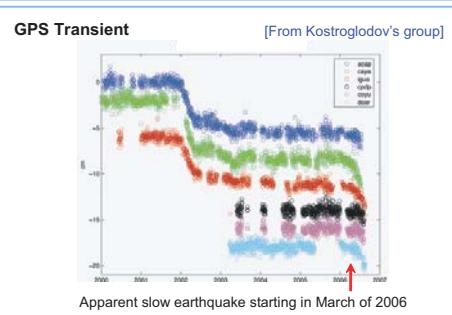
ATTENUATION STUDY



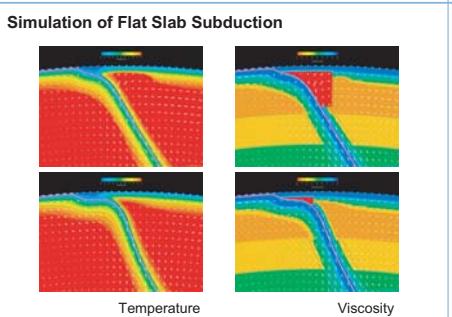
SURFACE WAVE STUDY



GPS STUDY



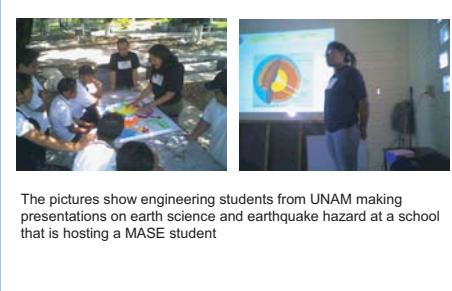
GEODYNAMICAL MODELING



FUTURE MASE SEISMIC ARRAY



OUTREACH



The pictures show engineering students from UNAM making presentations on earth science and earthquake hazard at a school that is hosting a MASE student

Presentation

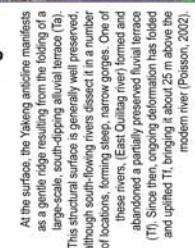
Modeling the shortening history of a fault-tip fold using structural and geomorphic records of deformation

M. Daëron*, J.-P. Avouac*, J. Charreau**

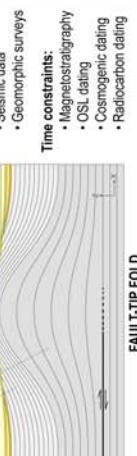
(* California Institute of Technology; ** Institut des Sciences de la Terre d'Orléans)

[Daëron et al., 2006, JGR, in press, <http://www.whooshingsounds.net/ecto/public/daeron-jgr-2006.pdf>]

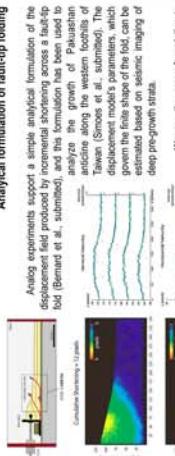
Yakeng fold



At the surface, the Yakeng anticline manifests as a gentle ridge resulting from the folding of a large-scale surface-cutting alluvial terrace (Ta). This structural surface is generally well preserved, although such surface terraces are dissected in a number of locations, forming steps, narrow gorges. One of these rivers (east Qilong river) and abandoned a partially preserved fluvial terrace (T1). Since then, ongoing deformation has folded and uplifted T1, bringing it about 25 m above the modern river (Poisson, 2002).



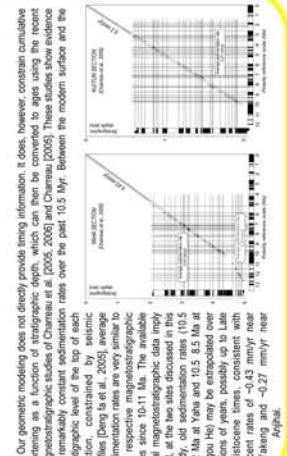
Analytical formulation of fault-tip folding



Regional setting

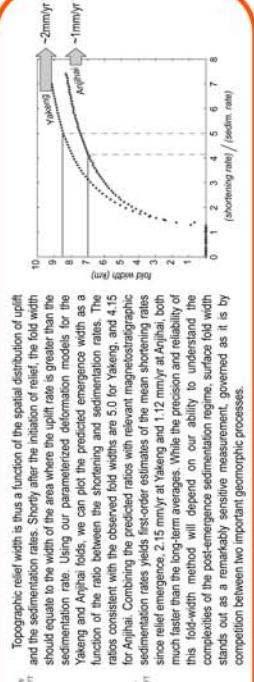


Sedimentation rates



Discussion

The study of these case examples highlights some simple interactions between folding, sedimentation and erosion. Topographic relief can only accrue where and when tectonic uplift is faster than sedimentation (see also discussion in Simone et al. 2006). Thus, in the early phases of the histories of both folds, synthetic sedimentary units extend continuously across the fold, and no topographic relief builds up (A). As shortening rate increases, maximum uplift rates overcome the sedimentation rate. The width is a function of the spatial distribution of the uplift rates. As long as the hydrographic system has enough erosional power to sweep laterally back, and form and erode ridges as they are uplifted, relief remains negligible, and an antecedent surface is eroded, unconformably overlying older ridges (B), as observed on the northern flanks of Alphai. If the river is forced to entrenched in a narrow gorge because it does not have enough power to abrade laterally all the uplifted ridges, relief starts building up above the core of the anticline (C), producing something similar to the current situation of Yaqing. Eventually, the old ridge is expected to undergo secondary erosion driven by its own relief, as observed in the exposed core of Apilhai.



The seismic profile reveals a smooth, rather symmetric subsurface structure strongly suggestive of detachment folding, although the lower part of the section might be suggestive of small-scale ramping near the fold's core. Our fine-markers interpretation of the seismic data allows mapping 7 distinct markers across the fold (L1 to L7). For all seven markers, the structural relief areas are well-correlated with depth, constant, with <-1 km of finite shortening over a basal detachment located ~ 5 km below the surface. We conclude that the sediments below L7 are pre-tectonic strata.

Simple scenarios of shortening (constant shortening & sedimentation rates) provide a good fit to the seismic data alone. However, all of these models predict that the fold should have no or negligible topographic relief, because the sedimentation rate exceeds the uplift rate. This is clear indication that the ratio of shortening to sedimentation rates has recently increased.

Two-stage scenarios involving a recent acceleration of shortening yield a better fit to the combined subsurface + surface data set. Additional data is still needed to quantify this acceleration.

Topographic relief width is thus a function of the spatial distribution of uplift and the sedimentation rates. Shortly after the initiation of relief, the fold width should equate to the width of the area where the uplift rate is greater than the sedimentation rate. Using our parameterized deformation models for the Yakeng and Apilhai folds, we can plot the predicted emergence width as a function of the ratio between the shortening and sedimentation rates. The ratios consistent with the observed fold widths are 5.0 for Yakeng and 14.15 for Apilhai. Combining the predicted ratios with relevant magnetotelluric sedimentation rates yields first-order estimates of the mean shortening rates since fold emergence: 2.15 mm/yr at Yakeng and 11.2 mm/yr at Apilhai, much faster than the long-term averages. While the precision and reliability of this fold-width method will depend on our ability to understand the complexities of the post-emergence sedimentation regime, governed as it is by competition between two important geomorphic processes.



Anjihai fold

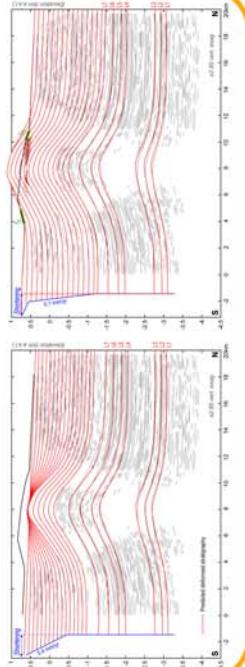
The surface fold is about 7 km wide, and exposes conglomerates of the Yuqiu lithology. Conglomerates, Nankou, and Dabachang (Nankeen) formations, unconformably overlain by Quaternary conglomerates and pebbles. On the flanks of the anticline, such Quaternary structures are well-preserved, forming talus slopes with slopes of $\sim 10\%$. We interpret these surfaces, rooted Tr1 and Tr2 as folded talus slopes which passively record deformation since their abandonment. Along the steep walls of the over-gap, the shallow structure of the fold is beautifully exposed, with dip angles up to 25° .

The seismic profile reveals a smooth, rather symmetric subsurface structure strongly suggestive of detachment folding, although the lower part of the section might be suggestive of small-scale ramping near the fold's core. Our fine-markers interpretation of the seismic data allows mapping 7 distinct markers across the fold (L1 to L7). For all seven markers, the structural relief areas are well-correlated with depth, constant, with <-1 km of finite shortening over a basal detachment located ~ 5 km below the surface. We conclude that the sediments below L7 are pre-tectonic strata.

Imposing a finite shortening of 1.55 km and a basal detachment depth of 4.5 km b.s.l., we can model the observed finite deformation of the neotectonic markers, using a 13-angle deformation model. The agreement between the predicted and observed dip angles is evidenced when all present-day seismic reflectors are un-correlated with depth, constant, with <-1 km of finite shortening over a basal detachment located ~ 5 km below the surface. We conclude that these reflectors therefore correspond to growth strata.

Simple scenarios of shortening (constant shortening & sedimentation rates) provide a good fit to the seismic data alone. However, all of these models predict that the fold should have no or negligible topographic relief, because the sedimentation rate exceeds the uplift rate. This is clear indication that the ratio of shortening to sedimentation rates has recently increased.

Two-stage scenarios involving a recent acceleration of shortening yield a better fit to the combined subsurface + surface data set. Additional data is still needed to quantify this acceleration.

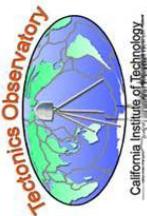


Topographic relief width is thus a function of the spatial distribution of uplift and the sedimentation rates. Shortly after the initiation of relief, the fold width should equate to the width of the area where the uplift rate is greater than the sedimentation rate. Using our parameterized deformation models for the Yakeng and Apilhai folds, we can plot the predicted emergence width as a function of the ratio between the shortening and sedimentation rates. The ratios consistent with the observed fold widths are 5.0 for Yakeng and 14.15 for Apilhai. Combining the predicted ratios with relevant magnetotelluric sedimentation rates yields first-order estimates of the mean shortening rates since fold emergence: 2.15 mm/yr at Yakeng and 11.2 mm/yr at Apilhai, much faster than the long-term averages. While the precision and reliability of this fold-width method will depend on our ability to understand the complexities of the post-emergence sedimentation regime, governed as it is by competition between two important geomorphic processes.

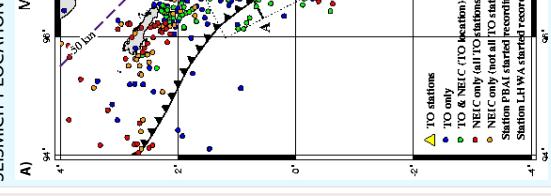


Seismicity in the Mentawai Region of Sumatra Using the Caltech Tectonic Observatory's Local Short-Period Seismic Network

Natalia I Deligne, Robert W Clayton, Erik Hauksson
Seismological Laboratory, California Institute of Technology

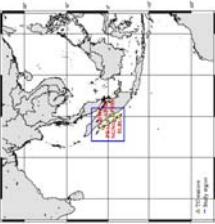


SEISMICITY LOCATION

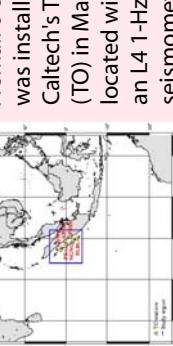


OVERVIEW

A small 6-station seismic array was installed in Sumatra by Caltech's Tectonic Observatory (TO) in May 2005. Each site, co-located with a GPS station, has an L4 1-Hz verticle seismometer recorded by a Nanometrics Taurus 24-bit logger. The data are continuously recorded on site and are retrieved every 3 to 5 months. Events are located using SEISAN. To date, 1094 earthquakes have been located for the period from May 19th - August 4th 2005. In general, for the entire study region (4N - 4S, 94E - 102E) only about 20% of the located earthquakes have also been located by the National Earthquake Information Center (NEIC), suggesting that background seismicity is far greater than previously estimated. In the region near Nyaung Nyang Island, the NEIC reports almost no events, but the TO stations pick up considerable activity. Most of the located earthquakes appear to be M4 - M5.



SAMPLE WAVEFORM DATA



DATA QUALITY

The two land-based stations (ABGS and PSKI) have the clearest signal. Proximity of an earthquake also strongly influences clarity of signal at a given station. Station PBAI has the clearest island-based signal. Station LHWA has a relatively good signal, and records numerous events, presumably to the north of the array, that are not recorded by any other stations. Station NGNG records the surf at a nearby beach, while station SLBU is exceedingly noisy.

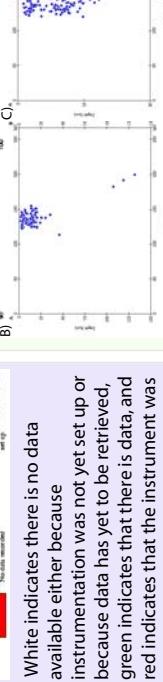
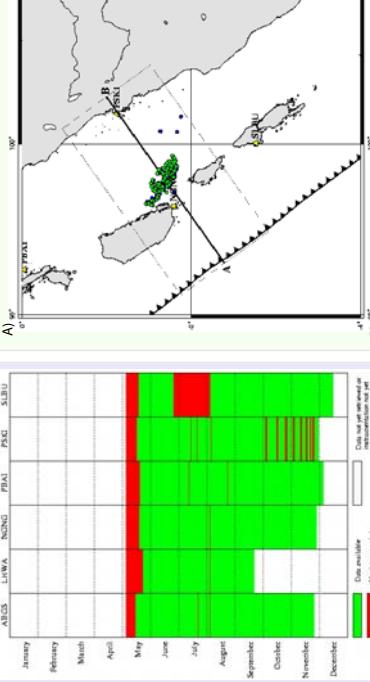
A) All located earthquakes in the study region (4N - 4S, 94E - 102E) for the period from May 19th - August 4th 2005, with the locations for crosssections for figure parts B, C, and D indicated. TO stations are shown as yellow triangles. Events were located from TO waveforms using SEISAN, and a minimum of four stations had a signal for an event to be located. Events in blue are only located by the TO, green events are located by both the TO and the NEIC, while red and orange events have only been located by the NEIC. For most such events, however, at least one or two TO stations picked up the signal. Slab contours are based on global seismicity (Gudmundsson and Sambridge, 1998).

B) Cross-section across the northern portion of the study region. For this and the other cross-sections, only TO-located earthquakes are shown (i.e. blue and green events on location map).

C) Cross-section across the central region. Note the linear feature occurring about 200 km from the trench, going to a depth of 50 km. Although focal mechanism have not yet be done, such a linear feature suggests strike slip motion. Earthquakes not along this feature seem to match the suggested slab location pretty well.

D) Cross-section across the southern portion of the study region. There is a lot of activity at shallow depth above the presumed slab location. This region was selected for the double difference relocation study (see below)

DOUBLE DIFFERENCE RELOCATION



FUTURE WORK

- 1) Update the current 1D velocity model, with the aim of eventually developing a 3D velocity model
- 2) Calibrate the amplitude from the stations to be able to determine earthquake magnitude.

A) Earthquakes selected for the double difference relocation study. Blue circles correspond to the original SEISAN-based TO location, and the green circles correspond to the re-located epicenter. TO stations are yellow triangles, and the cross-section profile is indicated.

B) SEISAN-based location cross-section. The three deeper earthquakes are the three earthquakes outside of the main cluster on the location map.

C) Re-located events. Note the strong linear feature; although we do not yet have focal mechanism, this feature suggests strike-slip motion. If this activity is indeed strike slip motion, this would imply that the Great Sumatran Fault does not take up all the strike-slip motion along the Sunda arc; there is a component taken off shore closer to the trench.

A) Example of event with high-quality signal. The NEIC reports this earthquake as having a magnitude of 5.8; the epicenter we obtained for this event is 1.579 N, 96.845 E with a depth of 19.3 km (the NEIC reports location / depth of 1.46 N, 97.15 E, 24 km). Note the regular spikes at station NGNG corresponding to the surf.

B) Example of typical data. This event was not located as less than 4 stations recorded a signal. Note the noise at station LHWA,

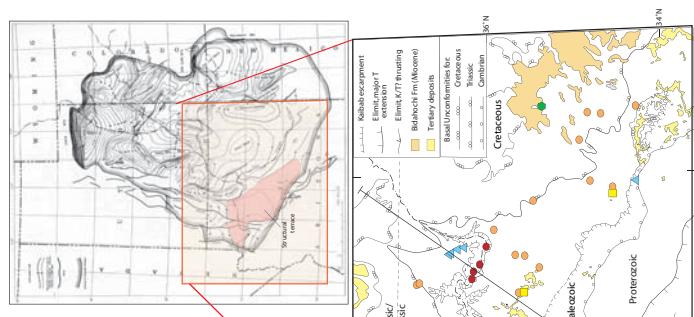
the extreme noise at station SLBU (this is what most of the data from station SLBU looks like), and the regular spikes at station NGNG corresponding to the surf.

White indicates there is no data available either because instrumentation was not yet set up or because data has yet to be recorded, green indicates that there is data, and red indicates that the instrument was not recording.

Unroofing of the southwestern Colorado Plateau from (U-Th)/He apatite thermochronometry

Rebecca M. Flowers (rflowers@gps.caltech.edu), Brian P. Wernicke, and Ken A. Farley
California Institute of Technology Division of Geological and Planetary Sciences

(U-Th)/He RESULTS

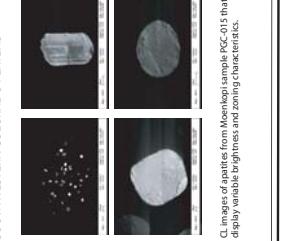
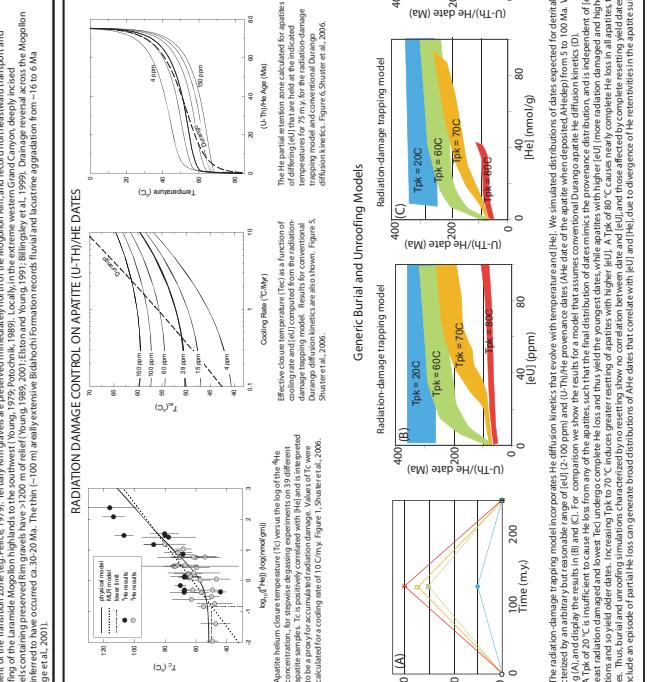


STRACT

A screenshot of the Microsoft Word ribbon interface. The 'Page Layout' tab is highlighted in blue, indicating it is the active tab. Other tabs visible include 'File', 'Home', 'Insert', 'Page Layout', 'References', 'Mailings', 'Review', and 'View'. The main content area shows a document with several sections of text and a large blue rectangular callout box overlaid on the page.

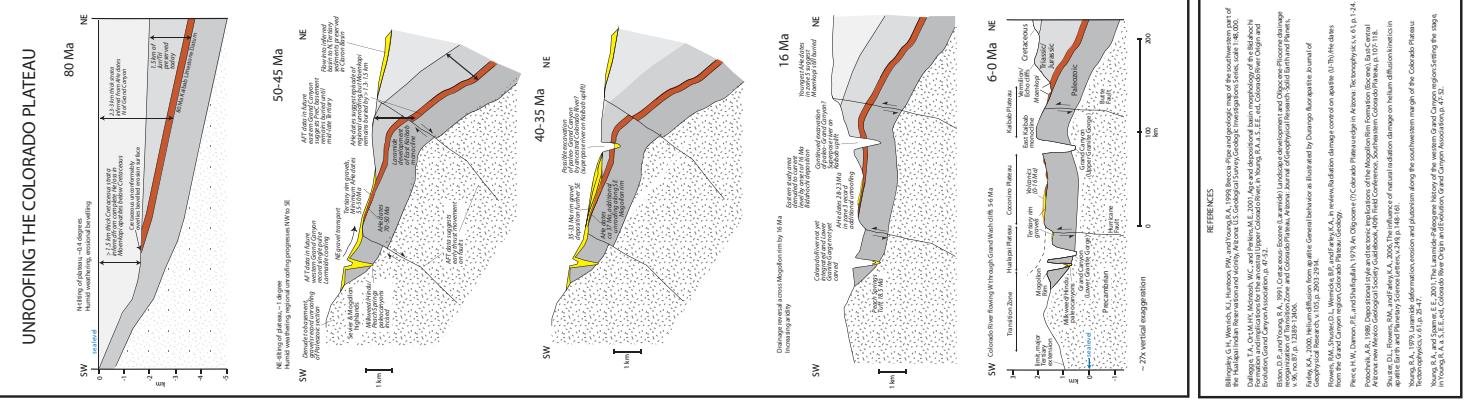
The southwestern portion of Grand Canyon, the Colorado Plateau is characterized by a broad, 1,000-m (3,280-ft) high, NE-dipping structural terrace that is the southern extension of the Tontoan Kasha-Katuwe plateau. The plateau surface is an elongation of the Tasse Mountain formation. To the northeast, progressively younger Mesozoic and Cenozoic formations are deposited as a series of cutouts known as the Great Box Canyons. To the west, the plateau is structurally delineated by major normal faults of the Basin and Range province, and to the southward it is bounded by the Grand Canyon where it cuts the topography high along the Colorado River. The Colorado Plateau is a relatively arid region, receiving an average annual precipitation of 500 mm (19.7 in). A north-south escarpment, the Mogollon Rim, extends 500 km (310 mi) across the plateau.

This pattern of relief (Young, 1989; Potocznik, 1989) locally in the extreme western Grand Northwest (Young, 1989; Potocznik, 1989). Locally, in the extreme western Grand

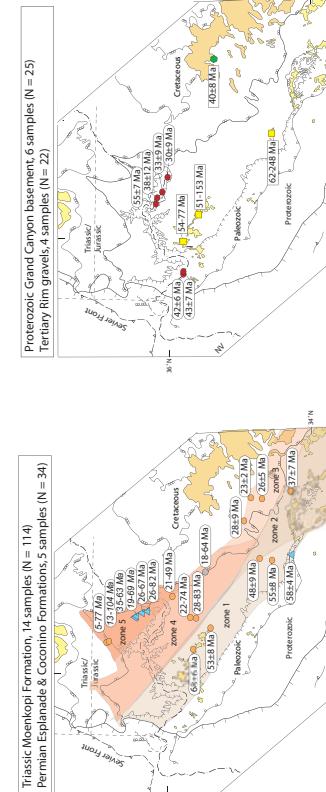


CL images of apatites from Moenkopi sample PGC-015 that display variable brightness and zoning characteristics.

UNROOFING THE COLORADO PLATEAU

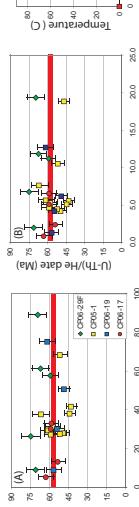


(U-Th)/He RESULTS

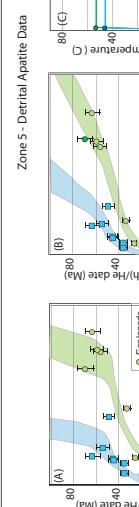


The clear patterns emerge from the Ane dataset. First, the spatiotemporal dates are younger to the northeast with increasing distance from the source. Second, the dates are older to the northwest, indicating that the dates are older to the northwest. Third, the dates along the Meghalaya range from 1.2 to 2.3 days later than those dates in patches to the northeast (Fig. 5j). For this reason, dates for the former are reported as a weighted mean, while the patchy dates are reported for the latter. Zone 3 samples are characterized by large ranges (Table 1) and can be explained by the effect of individual damage on the relatively low rainfall zone below.

Zone 1 - Detrital Apatite Data

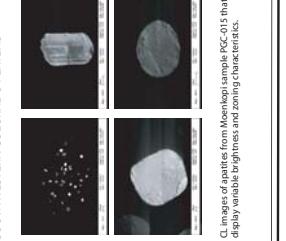
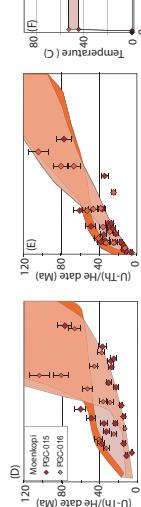


卷之三



A graph showing Force F (in N) on the y-axis versus time t (in seconds) on the x-axis. The force increases from 0 N at $t = 0$ s to 30 N at $t = 20$ s, and remains constant at 30 N until $t = 30$ s.

Time t (s)	Force F (N)
0	0
20	30
30	30



CL images of apatites from Moenkopi sample PGC-015 that display variable brightness and zoning characteristics.



GPS at CTO: Geodetic Arrays and Data Processing

Jeff Genrich, John Galetzka, Jean-Philippe Avouac, Kerry Sieh, Mark Simons, and Brian Wernicke

Introduction and Current Status

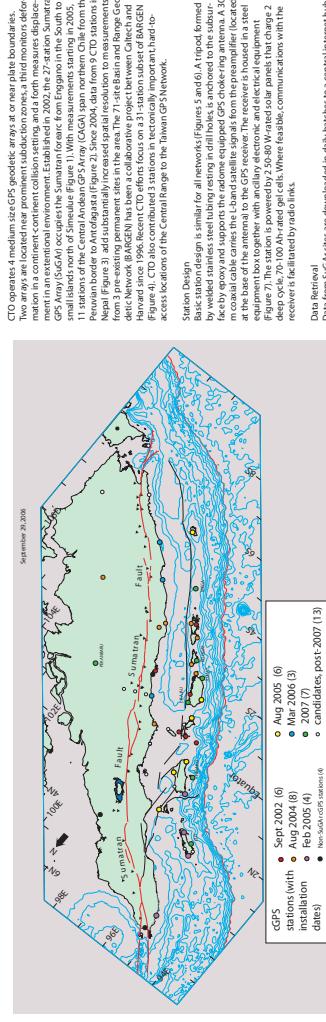


Figure 1. Map location of existing and future SugA sites in Sumatra.

Future Development

Site Augmentation

A major addition of sites is scheduled for SugA. A dense network in central Chile is key to the success of research activities in an extensional environment. Established in 2002, the 27-station SugA Array occupies the Santiago forearc from Engomar in the South to small islands north of Simeule in the North (Figure 1). With measurements starting in 2005, 15 stations of the Central Andes GPS Array (CAGA) span northern Chile from the Peruvian border to Arica (Avouac et al., 2004). Since 2004, data from 10 CAGA stations have been collected by the Chilean Geodetic Network (CONICET) (Figure 2). With measurements starting in 2005, 13 stations of the existing array are now active. The CAGA array has been expanded to the Atacama Desert in northern Chile and the western Bolivian Altiplano to the east. The array is currently under way in Tibet.

Real Time Data Streaming

If site measurements can be transmitted in near real time to central processing facilities, GPS geodetic networks located near plate boundaries will potentially become much more useful. The first real time data from the CAGA sites (JGRN, MCIA) have been tested successfully. Near real time data links for SugA have been explored earlier this year at the central facility of the satellite communications provider (ACG) in Baum. Using existing (low bandwidth) modems, data streams with a sampling interval of about 1 sec, were forwarded to a remote server in the U.S. The data were then processed in real time. Eliminating or reducing the range selected subsamples for data streams at 1 Hz or 2 Hz. Effort to expand real-time connectivity are underway for all CTO networks.

Data Processing

The displayed time series of station positions are based on a long term, independently of how many sources of errors. Omissions on their origin, thermal, and temporally correlated pseudorange and pseudo components with annual, diurnal contributions (Local multipathing, regional, and semi-annual frame induced errors produce additional, spatially correlated noise. Eliminating or reducing the range selected subsamples for data streams will be a valuable tool for real time data processing and analysis.

Station Design

Based on design similar for all networks (Figures 5 and 6), a tripod formed by welded stainless steel tubing resting in drift holes is anchored to the substrate by epoxy and supports the radome equipped GPS tracking antenna. A 30 m coaxial cable connects the antenna to the GPS. The tripod is connected to the base of the antenna by a 1 m long vertical segment. The antenna and equipment box together weigh about 10 kg. The equipment box contains a self contained power supply, a solar panel, a battery, and electrical equipment (Figure 7). The station is powered by 2.5-8.0 W wind sole panels that charge 2 deep cycle, 70/100 Ah rated gel cells. Where feasible, communications with the receiver are handled by dedicated links.

Data Retrieval

Data from SugA sites are downloaded in daily batches to a central internet hub in Bali, Indonesia. Through low bandwidth serial port modems of a regional communications satellite, several SugA sites have internet links by long range Ethernet radio bridges terminating at internet hubs of local universities (UCLA and UCI) and at the regional government administration plates. The regional network is fed in a chain along government, currently you download via satellite manual downloads.

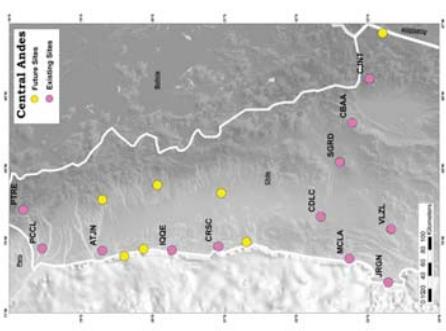


Figure 2. Map location of existing and future CAGA sites in northern Chile.

Data Processing

For each 24 hour time period, all available raw receiver image files are converted to Kxz Rnxz format. Using the GAMIT software package, each Rnxz file is processed with corresponding phase observations from nearby global reference stations to yield day station positions with respect to the network. The network is then refined using the 2000 reference stations, which are then used to calculate Fitter GPS with the help of solutions from several global IGS networks. North-, east- and up-components of the day station positions are displayed on network specific web pages (Figures 8 and 9).

Future

For the next 10 years, the SugA array will be expanded to cover the entire Andean mountain belt. The array will consist of approximately 1000 stations, which will be distributed in the Andes, the Andean Foothills, and the High Sierra. The array will be used to monitor the tectonic evolution of the Andes and the impact of the Andes on the rest of the world.

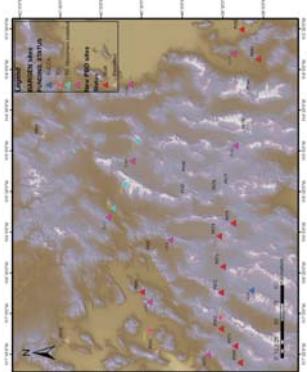


Figure 3. Map location of existing and future sites.



Figure 4. Map location of CTO subnetwork of BAGEN.



Figure 5. Interior of equipment box for station PCCL.

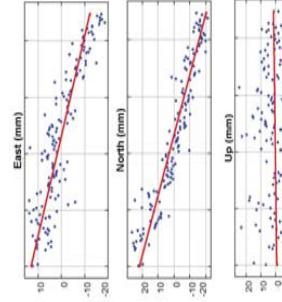


Figure 6. CAGA station PPRE.

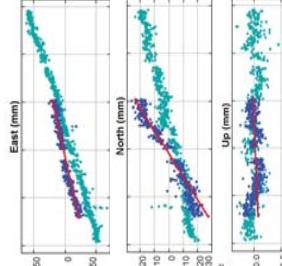
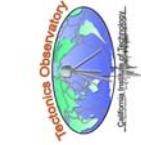


Figure 7. Interior of equipment box for station PCCL.

Figure 8. East-, north-, and up-components of daily positions for SugA site LEWK in ITR2000.

Figure 9. East-, north-, and up-components of daily positions for SugA site LEWK in ITR2000.

Figure 10. Map of GPS network. The map shows the Andes mountain range with various GPS station locations marked. A legend indicates: Nepal (blue triangles), Future TD Sites (yellow triangles), Existing TD Sites (pink triangles), and Other Existing Sites (black circles). A scale bar shows distances from 0 to 300 Kilometers.



Geodynamic models for the "uplift" and erosion of the Colorado Plateau



Frederic Herman and Mike Gurnis

Tectonics Observatory, Geological and Planetary Science
California Institute of Technology

Our aim is to test a range of geodynamic models that can explain the uplift of the Colorado Plateau, if uplift there was. Many mechanisms have been proposed and summarized in the literature [e.g. McGeechkin 1980]. The main mechanisms fall in three different categories: (a) Late Tertiary uplift related to early tertiary uplift related to Sevier-Laramide orogeny (80 to 40 Ma) [e.g. McQuarrie and Chase 2000]; (b) mid-Tertiary uplift related to removal of flat slab [Spencer 1996], release of negative dynamic topography [Mitrovica et al 1989, Gurnis 1992] or hydration of the lithosphere from volatiles derived from the Farallon slab [Humphries et al 2003]) and (c) Late-Tertiary uplift associated with regional extensional tectonism, either by large removal of instable lithosphere [Bird 1979] or heating from below. We propose to test these geodynamic models using the codes available through the Computational Infrastructure for Geodynamics framework.

We use a mantle convection model, so-called CitcomT [e.g. Billen et al 2003] to solving the equations of momentum, continuity and transport equations :

$$\begin{aligned} u_{ij} &= 0, \\ -P_{,j} + (\eta u_{i,j} + \eta u_{j,i})_j + R\alpha\dot{\varphi}\delta_{ij} &= 0, \\ (\dot{\varphi})_{,i} + u_i(\dot{\varphi})_j &= (\dot{\varphi})_{,ji} \end{aligned}$$

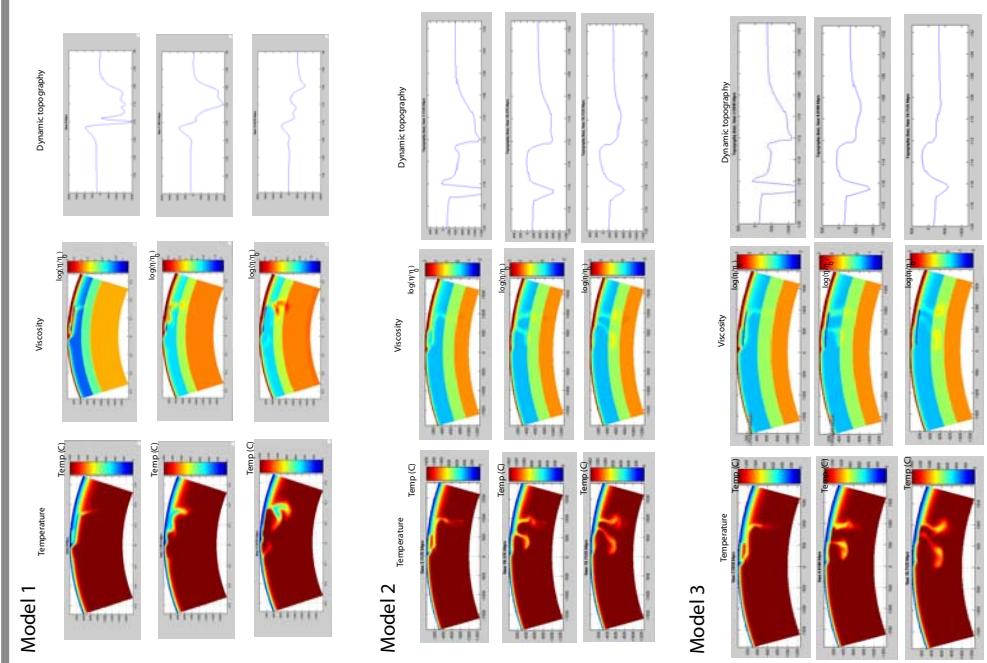
in 3D and in spherical coordinates.

It computes dynamic topographies using the boundary flux method [Zhong et al 1993].

The viscosity varies laterally and is temperature-pressure dependant. Lateral and radial viscosity variations can be included in the model, by setting the effective viscosity to the desire value in a given region.

We present three models which test the hypothesis of uplift of the Colorado Plateau due to removal of the flat subduction. In each model, we include a buoyant continental tectosphere that is shaved by the flat subduction. In each model, decoupling between the subducting plate and the overriding lithosphere is ensured by including a region of low viscosity between the two. Model 1 simulates an abnormally thick oceanic lithosphere below the plateau, Model 2 assumes a thinner oceanic plate and Model 3 is equivalent to model 2 except that the zone of decoupling is made larger.

We use a mantle convection model, so-called CitcomT [e.g. Billen et al 2003] to solving the equations of momentum, continuity and transport equations :



We here use a Surface Process Model, which simulates the evolution of hillslope uplift on the edges of the plateau. This, in turn, causes isostatic down erosion along these scarps and forming an internally drained basin. Ultimately, the basin will overflow and lead to rapid erosion along the Grand Canyon. The model starts with a basal level drop. Consequently, erosion is induced at the edges of the plateau. This, in turn, causes isostatic uplift on the edges of the plateau, slowing down erosion along these scarps and forming an internally drained basin. Ultimately, the basin will overflow and lead to rapid erosion along the Grand Canyon.

Computations were performed on the Pangaea facilities at the Geological and Planetary Sciences Division, California Institute of Technology

Thermal modelling of metamorphism and exhumation in Western Nepal and India

Frederic Herman*, Jean-Philippe Ayouac*, Amos Aikman†, Julien Celerier†, Mark Harrison @

* Tectonics Observatory, California Institute of Technology
+ Research School of Earth Sciences, Australian National University
@ Department of Earth and Space Sciences, University of California Los Angeles



The geological association of the Himalayan range is a juxtaposition of inverted metamorphic sequences in the footwall of the Main Central Thrust (MCT) with a belt of Miocene leucogranites emplaced above the fault (Figure 1). The MCT extends across the 250 km length of the Himalayan orogen and is the dominant structural feature of that mountain belt.

The inverted metamorphic sequences beneath the MCT ramp has been interpreted as a reactivation of the thrust following ~ 10 myr of inactivity [e.g. Harrison et al 1997]. Measured monazite ages from the lower Lesser have been used to interpret reactivation of the MCT at ca. 8 Ma and activation at ~ 6 Ma. More recently, Bollinger et al 2004 and 2006 showed that shortening across the Himalaya can be explained by accommodation by a single fault, the Main Himalayan Thrust (MHT), and that the growth of Himalayan wedge has resulted from underplating and development of a duplex. In this latter scenario, the MCT zone corresponds to the MHT exhumed at the surface.

In both instances, an increase of exhumation from about 8 Ma is required. We test here these models in Far West Nepal and India, where new thermochronological data have been collected (Ar³⁹/Ar⁴⁰ in muscovite, Raman Spectrometry and (U-Th)/He in zircon). We use a thermal-kinematic model which solves the heat transfer equation coupled with an inversion algorithm, the Neighbourhood Algorithm (Cambridge 1999).

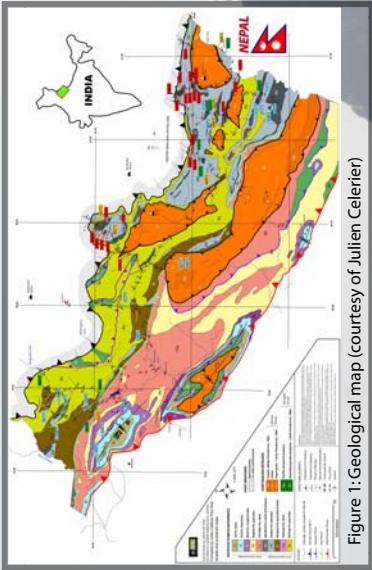


Figure 1: Geological map (courtesy of Julien Celerier)

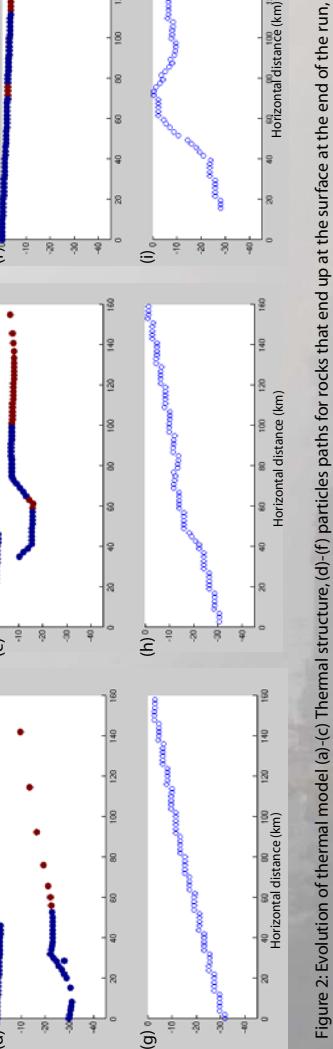
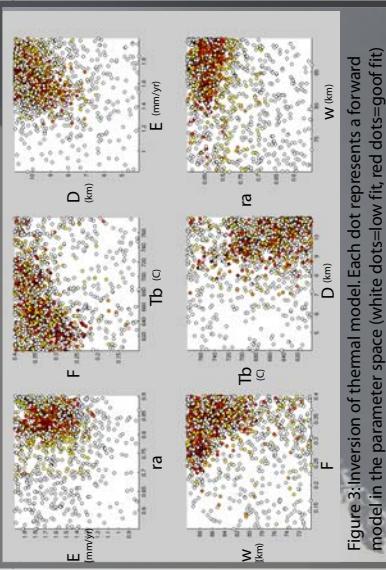


Figure 2: Evolution of thermal model (a)-(c) Thermal structure, (d)-(f) particles paths for rocks that end up at the surface at the end of the run, (g)-(i) evolution of MHT during model run.

References

- Harrison, T.M., et al 1997, A late Miocene-Pliocene origin for Central Himalayan inverted metamorphism, *Earth and Planetary Science Letters*, 146, p.E1-E8.
- Bollinger, L., et al 2004, Thermal structure and exhumation history of the Lesser Himalaya., *Tectonics*
- Cambridge 1999, Geophysical Inversion with a Neighbourhood Algorithm - I. Searching a parameter space *Geophys. J. Int.*, 138, 479-494
- Computations were performed on the Pangaea facilities at the Geological and Planetary Sciences, California Institute of Technology.

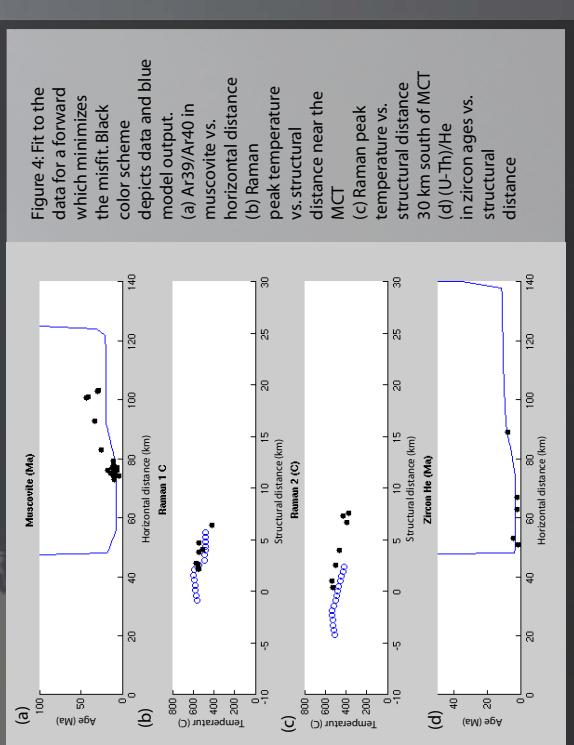
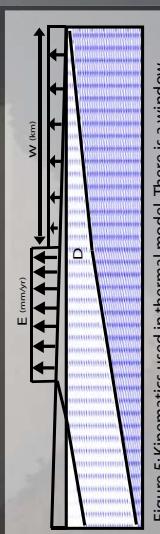


Figure 4: Fit to the data for a forward model. Each dot represents a forward model in the white dots=good fit, red dots=goof fit

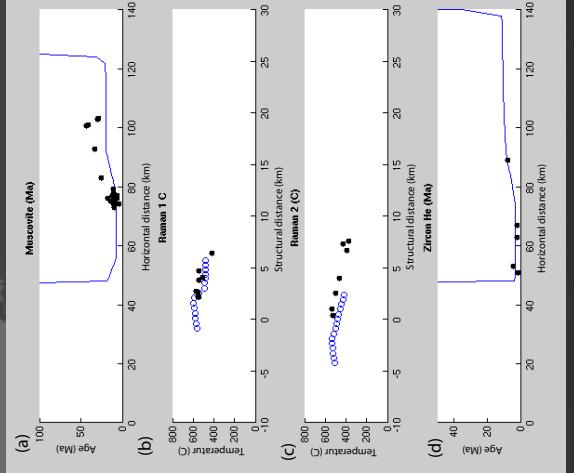


Figure 5: Kinematics used in thermal model. There is a window of accretion, the upper plate overthrust the lower plate at $(1-v)^{*}v$, with v being the total shortening rate.

Spectral element modeling of earthquake nucleation and spontaneous rupture on rate and state faults

Yoshihiro Kaneko¹, Nadia Lapusta^{1,2}, Jean Paul Ampuero³

1 Introduction
 Nucleation and spontaneous dynamic propagation of earthquakes on rate and state faults have been successfully modeled in the framework of boundary integral methods (BIM) (Rice and Ben-Zion, 1986; Lapusta et al., 2000). However, these studies have been mostly restricted to planar faults embedded into a uniform elastic space or half-space due to the nature of BIM. At the same time, observations show complicated crustal structures (such as asymering and fault damage zones) and non-planar fault geometries. It is important to include those factors into earthquake models, combining them with laboratory derived constitutive fault relations such as rate and state friction. In this work, we use 3-D spectral element method (SEM) to model earthquake nucleation and propagation of spontaneous rupture on a vertical strike-slip fault governed by rate and state friction. Our ultimate goal is to develop a SEM framework for simulating long-term deformation histories, in terms of sequences of earthquakes and combination of seismic and seismic sliding.

6 SCEC 3D code comparison

A study of Day et al., 2005, similar to our comparison,

- linear slip-weakening friction
- non-smooth initial conditions (both in time and space)
- the errors are represented as rupture arrival time (slip rates at 1 mm/sec)

7 Initial conditions for our test cases



To have similar effective lip weakening of friction at the comparison location, we choose $\alpha = 0.13$, $b = 0.18$, and $L = 0.133$ m. Hence, our initial conditions are similar to the initial conditions of the SCEC comparison problem.

8 2D SEM simulations using rate-state (aging law) friction

In the standard formulation with constant effective normal stress $\bar{\sigma}$, the shear strength τ can be expressed as:

$$\begin{aligned} \tau &= \sigma_1 \bar{\sigma} = \sigma [\mu_0 + b \ln(V/V_0) + b \ln(V_0/L)] \\ &\quad d\bar{\sigma}/dt = 1 - Vg/L \quad (\text{aging law}) \\ &\quad dV/dt = (Vg/L) \ln(V/V_0) \quad (\text{slip law}) \end{aligned}$$

3 Boundary integral method (BIM)

- Only the boundary (on the fault) is discretized
- Wave propagation is accounted for analytically through theoretically derived kernels
- Accurate and efficient but relatively limited in terms of geometry and bulk properties

4 Rate and state dependent friction

In the standard formulation with constant effective normal stress $\bar{\sigma}$, the shear strength τ can be expressed as:

$$\begin{aligned} \tau &= \sigma_1 \bar{\sigma} = \sigma [\mu_0 + b \ln(V/V_0) + b \ln(V_0/L)] \\ &\quad d\bar{\sigma}/dt = 1 - Vg/L \quad (\text{aging law}) \\ &\quad dV/dt = (Vg/L) \ln(V/V_0) \quad (\text{slip law}) \end{aligned}$$

5 Anti-plane test problem to compare SEM and BIM

Potentially unstable sliding when $\alpha > b$
 – Stable sliding when $\alpha < b$

6 2D SEM

Discretization: 45 km, periodic boundary, nucleation patch, fault boundary, slip boundary.

The rupture nucleates at the center of the fault, and SH wave propagates in the medium. The star indicates our comparison location. The right figure represents the SEM and BIM slip rates as a function of time at the comparison location for the best-resolution runs.

9 Measures of errors: rupture arrival time vs. cross-correlation

Two different measures of errors:

1. Rupture arrival time difference
 Rupture arrival time = time at a peak of slip rates (our current definition)
2. Cross-correlation time difference

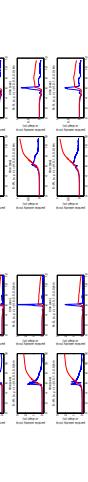
Our comparisons are similar to the SCEC code validation (Harris et al., 2004), and we consider the following questions:

- the accuracy of SEM solutions with respect to BIM
- appropriate measures of errors
- abrupt vs. smooth initial conditions
- the state variable updates, integration or direct use of evolution law
- simulations with aging law vs. slip law

14 Dynamic rupture simulation in homogeneous media (continued)

Snapshots of horizontal displacement (m/s) on the fault in the interval of one second. Homogeneous property ($V_s = 3464$ m/s, $V_p = 6000$ m/s) is imposed. The quantity $(\bar{\sigma} - \bar{\sigma}_0)/(\bar{\sigma}_0 - \bar{\sigma}_{\infty})$ in the weakening region, $(1, 11)$, in the re-strengthening region, $(1, 11)$, in the nucleation region "B1" is 120 MPa, and initial shear stress is 70 MPa except at "nucleation region" (B1) (6 MPa).

The rupture speed becomes super-shear near the free surface.



Displacement and velocity seismograms at 4-km horizontal distance (14 km away from the center of nucleation). Right and left insets correspond to the simulations above and below respectively. The top, middle, and bottom panels correspond to different off-fault locations (on the fault, 0.5 km away, and 1.0 km away from the fault).

15 Dynamic rupture simulation in a layered structure

Same as above except that the effective normal stress gradually increases with depth and constant (120 MPa) below 6-km depth. The transition from velocity-weakening to velocity-strengthening is smooth and occurs at 2-km depth.

16 Conclusions

- We have incorporated rate and state friction into 3D SEM dynamic rupture code (Ampuero, 2004). The snapshots of the slip rates on the fault are shown above. The initial conditions used for this simulation are similar to the smooth case in 2D, where the nucleation proceeds gradually. The rupture speed transitions from sub-shear to super-shear in the in-plane direction, consistently with daughter crack mechanism of Burridge-andrews for slip weakening friction (Andrews, 1976). Note that the transition is observed here for rate and state friction laws.
- SEM and BIM give virtually indistinguishable solutions to the test problem with the nucleation and spontaneous rupture propagation when the node spacing is small enough.

17 Future work

- Understand how much can be learned from the near-field seismic records in terms of the history of slip or slip rates on the fault with different weakening mechanisms.
- Develop SEM to include variable time steps to simulate long-term deformation history of a fault.

10 Convergence rates: rupture arrival time vs. cross-correlation

On the vertical axis, the rupture arrival time difference or the cross-correlation time difference between the highest-resolution run and lower-resolution runs is shown. In general, BIM and SEM give similar convergence rates in terms of the rupture arrival time difference. When the cross-correlation time difference is used, SEM shows smaller errors than those of BIM. For one of such cases, the slip rates for both BIM and SEM are shown on the bottom panels.

11 Geometry of 3D simulation



12 Nucleation and super-shear transition of spontaneous rupture



13 Dynamic rupture simulation in homogeneous media



We have incorporated rate and state friction into 3D SEM dynamic rupture code (Ampuero, 2004). The snapshots of the slip rates on the fault are shown above. The initial conditions used for this simulation are similar to the smooth case in 2D, where the nucleation proceeds gradually. The rupture speed transitions from sub-shear to super-shear in the in-plane direction, consistently with daughter crack mechanism of Burridge-andrews for slip weakening friction (Andrews, 1976). Note that the transition is observed here for rate and state friction laws.

14 Dynamic rupture simulation in homogeneous media (continued)

Snapshots of horizontal displacement (m/s) on the fault in the interval of one second. Homogeneous property ($V_s = 3464$ m/s, $V_p = 6000$ m/s) is imposed. The quantity $(\bar{\sigma} - \bar{\sigma}_0)/(\bar{\sigma}_0 - \bar{\sigma}_{\infty})$ in the weakening region, $(1, 11)$, in the re-strengthening region, $(1, 11)$, in the nucleation region "B1" is 120 MPa, and initial shear stress is 70 MPa except at "nucleation region" (B1) (6 MPa).



Displacement and velocity seismograms at 4-km horizontal distance (14 km away from the center of nucleation). Right and left insets correspond to the simulations above and below respectively. The top, middle, and bottom panels correspond to different off-fault locations (on the fault, 0.5 km away, and 1.0 km away from the fault).

15 Dynamic rupture simulation in a layered structure

Same as above except that the effective normal stress gradually increases with depth and constant (120 MPa) below 6-km depth. The transition from velocity-weakening to velocity-strengthening is smooth and occurs at 2-km depth.

16 Conclusions

- We have incorporated rate and state friction into 2D and 3D SEM dynamic rupture code to simulating a single earthquake.
- SEM and BIM give virtually indistinguishable solutions to the test problem with the nucleation and spontaneous rupture propagation when the node spacing is small enough.

17 Future work

- Understand how much can be learned from the near-field seismic records in terms of the history of slip or slip rates on the fault with different weakening mechanisms.
- Develop SEM to include variable time steps to simulate long-term deformation history of a fault.



Sense of shear and thermal evolution of the schist of Sierra de Salinas, California

Steven Kidder¹, Jason Saleeby¹, Frédéric Herman¹, Mihai Ducea²

¹California Institute of Technology ²University of Arizona

Abstract

Improved knowledge of kinematic and geochronologic relationships is important for understanding the late Cretaceous demise of the Salinia-Mojave continental arc, its effect on crustal composition, and the processes of sediment subduction, tectonic erosion, ridge collision and exhumation. The demise of the arc coincided with structural juxtaposition of forearc-type assemblages (the POR schists) and Eastern-zone plutonic rocks along the Salinas shear zone. We investigate this late Cretaceous episode using microstructural techniques, 2D modeling and Ar-Ar dating. Preliminary microstructural work indicates top-to-northeast ductile shear, similar to the sense of shear in lower-grade early Tertiary schists to the SE. Preliminary 2D modeling suggests that flat subduction may not be a feasible mechanism for schist emplacement as has previously been suggested.

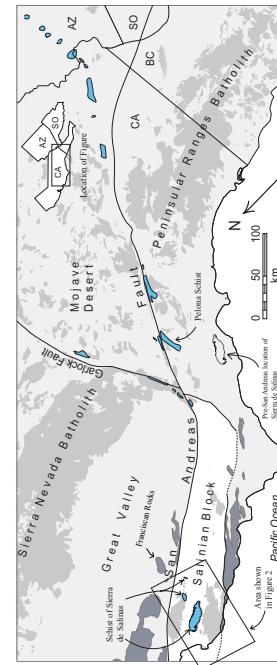


Figure 1. Map of parts of California, Arizona, Sonora, and Baja California showing some geologic features. The POR schists are shown in blue. Mesozoic granitic and related metamorphic rocks are shaded lightly. Fine-gridded areas are mainly Franciscan formation. Pre-San Andreas location of Sierra de Salinas from Powell (1993).

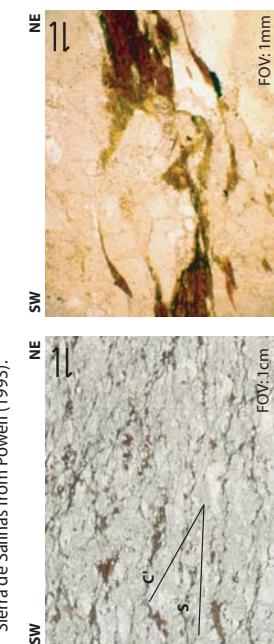


Figure 2. Micrographs of schist of Sierra de Salinas showing criteria used to determine sense of shear. The image at left shows typical C fabric. The image at right shows two biotite "fish."

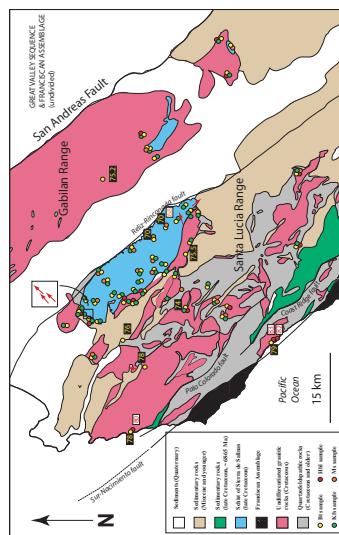


Figure 3. Geologic map of the Salinian central block (after Barbeau et al., 2005) showing sample locations for 40Ar/39Ar dating and sense-of-shear determinations. Boxed areas mark locations of samples previously dated by Barth et al. (2003) and Kistler and Champion (2001). Ar-work is in progress. See figure 1 for location. Red arrows indicate locations where a preliminary sense of shear is determined. Arrows indicate motion direction of upper plate.

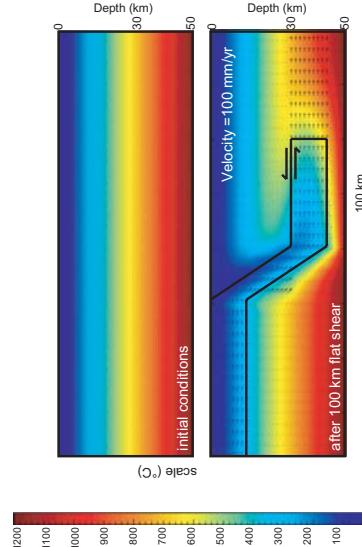


Figure 4. Thermal mechanical modeling can test the flat slab hypothesis (from East or West; e.g. Kidder et al., 2005; Grove et al., 2005). The key field observation is a convergence in upper and lower plate temperatures at initial upper plate conditions. Preliminary results shown above suggest insufficient heat is available with a flat slab scenario to reach observed temperatures in the schist.

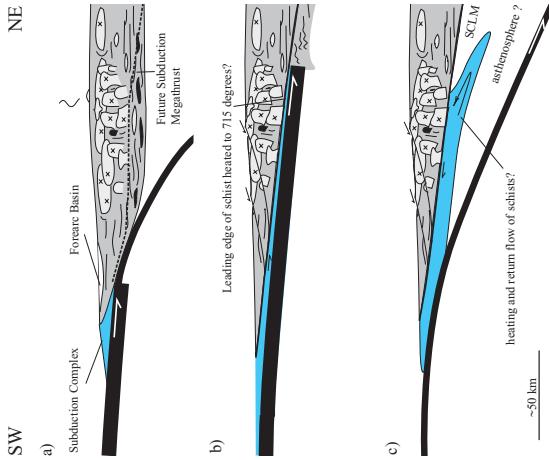


Figure 5. Cartoons depicting initial conditions and two possible scenarios for schist emplacement during the Late Cretaceous demise of the Salinia-Mojave continental arc. Schist is shown in blue. No vertical exaggeration. (a) Initial conditions of arc just prior to collision with overthickened oceanic crust (e.g. Saleeby, 2003). The dashed line in panel a depicts the future location of the megathrust. (b) Flat slab hypothesis for heating the schist (after Kidder & Dueca, 2006). As lower portions of the arc were sheared off, the leading edge of the schist was carried beneath the now-extinct volcanic arc heating the schist to over 700°. This hypothesis predicts top-to-southwest shear. (c) Alternative hypothesis for heating the schist (after Saleeby, 2003). Schist is heated at depth and returned to shallower levels by some exhumation process, possibly return flow. This process overprints structures that may have developed in an earlier deformation phase of top-to-southwest shear (e.g. panel b). This hypothesis is more consistent with results of 2D modeling and preliminary evidence of top-to-northeast sense of shear observed in the schist of Sierra de Salinas.

Receiver Function Analysis of the Middle American Subduction Zone in Central Mexico





YoungHee Kim, Robert W. Clayton
Seismological Laboratory, California Institute of Technology, Pasadena, CA 91125

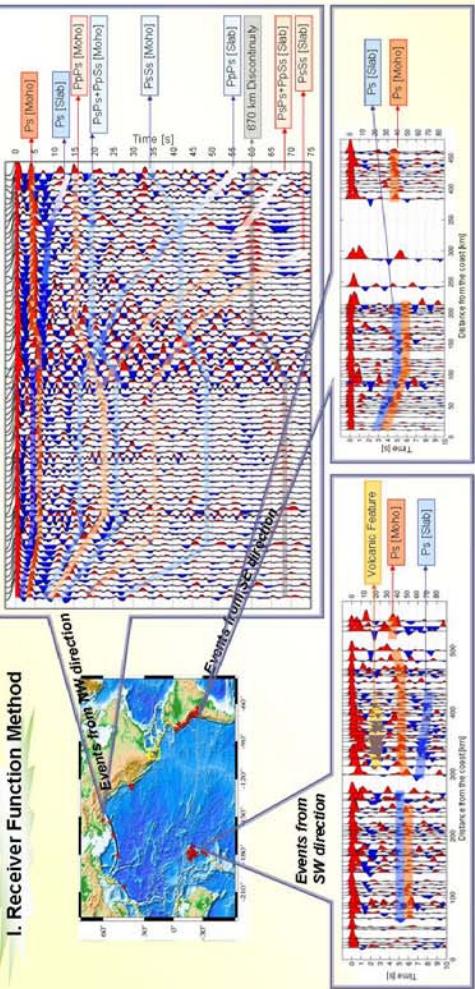
YoungHee Kim, Robert W. Clayton

ABSTRACT

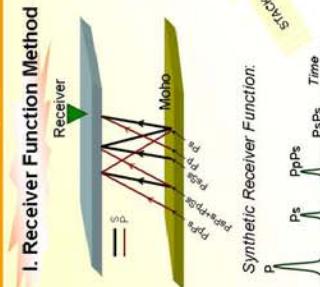
We have produced an image of the subducted Cocos plate beneath central Mexico with receiver functions utilizing data from the Meso America Subduction Experiment (MASE). The receiver function image shows that the subducting oceanic crust is shallowly dipping to the north at 15 degrees for 80 km from Acapulco and then horizontally underplates the continental crust for 200 km, all the way to the beginning of the Trans Mexican Volcanic Belt (TMVB). Both the continental and oceanic Moho are also clearly seen in the image. The cross section from Acapulco to a point 150 km to the north shows the erosion of the continental plate by the slab. There is no apparent evidence of crustal compressional features due to the underplating. Beneath the TMVB, the slab is dipping at about 20 degrees toward the end of the MASE array (near the Atlantic Coast). The continental Moho is about 40 km deep beneath the TMVB and shallows towards the north. An image is also produced by migration of the individual receiver functions. Multiples are also included in the image because in many instances they are stronger than the primary conversions. The image is tested by modeling the derived structure with a finite-difference wave propagation program.

RESULTS

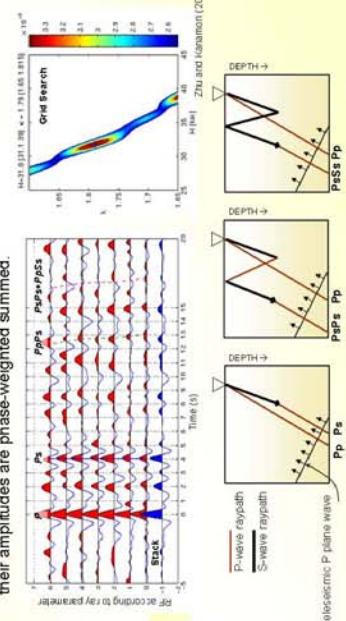
I. Receiver Function Method



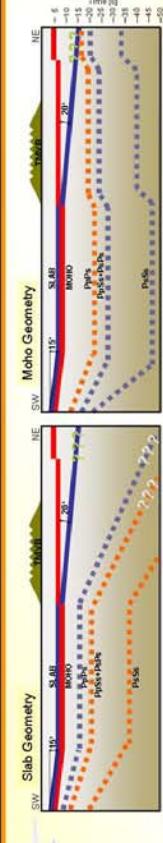
METHOD



H (moho depth) and κ (v_s/N) are estimated by maximizing the weighted summation function: $S(H, \kappa) = w_1(t_{Fe}) + w_2(t_{Fp}) - w_3(t_{Sp})$. The arrival times for Ps , $BnPs$, and $PsPc$ are calculated for given (H, κ) .

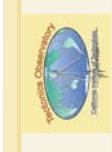


CONCLUSIONS & FUTURE WORK



Future Work

- Focus on obtaining the slab geometry at the northern end of the MASE line using more data from future events
 - To provide reference velocity from inversion
 - To develop migration algorithm to deliver an inversion of scattered wavefield for variations in the Earth's elastic properties
 - 2D & 3D elastic-wave finite-difference modeling
 - To image upper mantle discontinuities by RF migration



The Rupture Characteristics of the 1999 Izmit Earthquake Sequence Using Geodetic and Seismic Data

A. Ozgun Konca¹, Don Helmberger¹, Sébastien Leprince¹, Jean-Philippe Avouac¹

¹Tectonic Observatory, California Institute of Technology e-mail: orgun@gps.caltech.edu

Results for 17 Aug. 1999, Izmit Earthquake



Figure 1. Local network (Kocaeli) locations and Harvard CMT solutions of the events used in this study.



Figure 2. Comparison of variations in hundred-thousand arrival times.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

The conventional method of doing source inversions involves "backtracking" of the arrival times of S and SH waves. The variation of hundreds of arrival times from a reference one-dimensional model are improved by a pathfinding scheme.

2. Joint Inversion of Duzce Earthquake



Figure 3. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the unadjusted Arrival Times.

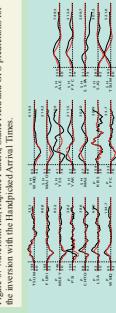


Figure 4. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the Calibrated Arrival Times.

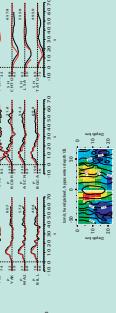


Figure 5. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the GPS data only.

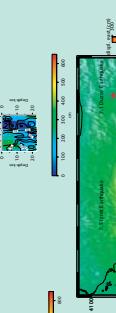


Figure 6. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the GPS and InSAR data only.



Figure 7. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the GPS and InSAR data and the GPS synthetics.

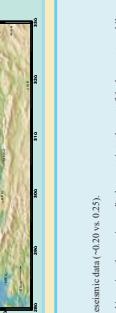


Figure 8. Waveform fits, Rupture Pattern, Static Field and GPS predictions for the inversion with the GPS and InSAR data and the GPS synthetics and the InSAR synthetics.



Figure 9. Slip model for joint tectonic, GPS and InSAR inversion of Duzce earthquake. Rupture time is between 1.5 s and 2.5 s.

Figures 3–9 show the results of the joint inversion of the tectonic, GPS and InSAR data. The results are consistent with the results shown in Figure 10. The results shown here do not include the strong-motion data yet available. The goal here is to see if data available. The results shown here do not include the strong-motion data yet available. The goal here is to see what difference it makes in teleseismic inversions. We have also

strengthened the Duzce source model with geodetic inversions (Dobrovolsky et al., 2002). The results shown here are quite consistent with the results proposed by Dobrovolsky et al. (2002) and GPS and strong-motion inversions. Another aim of this study is to understand whether supercavitation rupture speed proposed to explain strong-motion data is consistent with surface measurements. The geodetic data and surface offset measurements constrain the slip distribution significantly. Therefore, the time and input velocity as the two parameters to play with, in order to fit both teleseism and strong-motion data.

The goal of this comparison is to test whether teleseismic data together with near-field geodetic data can provide any significant constraint to rupture velocity or not.

(a) fast to the west

(b) fast to the east

(c) fast to the west

(d) fast to the east

(e) fast to the west

(f) fast to the east

(g) fast to the west

(h) fast to the east

(i) fast to the west

(j) fast to the east

(k) fast to the west

(l) fast to the east

(m) fast to the west

(n) fast to the east

(o) fast to the west

(p) fast to the east

(q) fast to the west

(r) fast to the east

(s) fast to the west

(t) fast to the east

(u) fast to the west

(v) fast to the east

(w) fast to the west

(x) fast to the east

(y) fast to the west

(z) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east

(ss) fast to the west

(tt) fast to the east

(uu) fast to the west

(vv) fast to the east

(ww) fast to the west

(xx) fast to the east

(yy) fast to the west

(zz) fast to the east

(aa) fast to the west

(bb) fast to the east

(cc) fast to the west

(dd) fast to the east

(ee) fast to the west

(ff) fast to the east

(gg) fast to the west

(hh) fast to the east

(ii) fast to the west

(jj) fast to the east

(kk) fast to the west

(ll) fast to the east

(mm) fast to the west

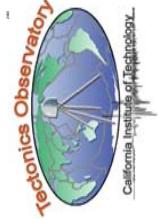
(nn) fast to the east

(oo) fast to the west

(pp) fast to the east

(qq) fast to the west

(rr) fast to the east



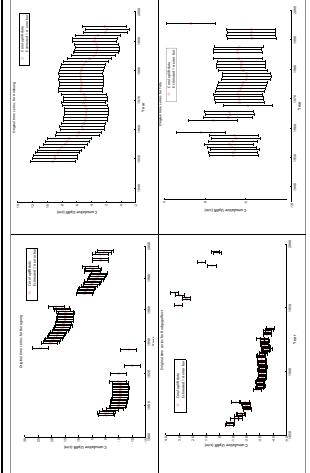
Time Dependent Slip: A Mathematical Study of Historic Displacement Near the Sunda Megathrust Subduction Zone

A. Kositsky, M. Chleb, J-P. Avouac, K. Sieh - California Institute of Technology, GPS division
ap@caltech.edu

ABSTRACT

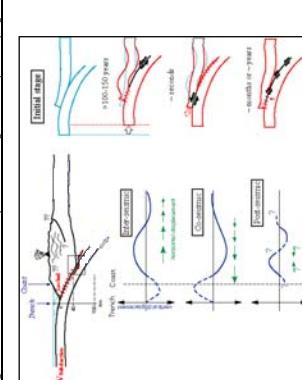
We invert displacement data using principal component analysis (PCA) and the Okada formulation for three-dimensional displacements using GPS data (Work in progress)

D. H. Natawidjaja - LIPI, Indonesia



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

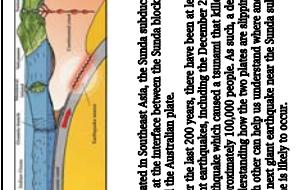
After taking into account factors such as a change in sea level height, the two-prong method provides a more accurate position of the crust to determine its relative to one part of the earth that has been taken during the same period. This allows us to calculate a reasonable detailed record of vertical deformation over the last 40 to 50 years.



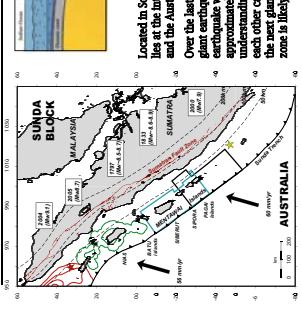
The earthquake cycle for two plates on the sections on cross primarily of three phases:
1. Initial seismic loading: As the subducting plate and the overriding plate converge, a section of the interface between the two plates breaks and stress builds over time.
2. Co-seismic Events: The stress eventually becomes too much for the plates to bear a sudden and violent slipping between the two plates, resulting in an earthquake. After the earthquake, the plates gradually return to a temporary state of equilibrium, often accompanied by gradually decreasing time-dependent slip.

3. Post-seismic Behavior: A return of tension in the crust and the resulting surface deformation are located in the same place on the right side of the diagram. This information implies that we are able to figure out what the displacement at the surface that has been mapped to solve for what is happening at depth.

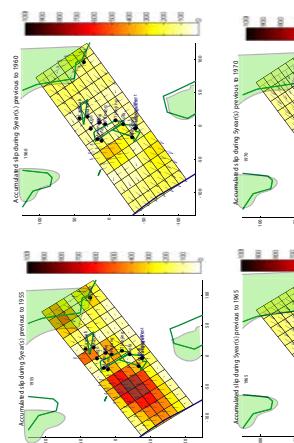
This diagram also shows the location of the Tondo Observatory, which is located at the interface between the Sunda block and the Australian plate.



Located in Southeast Asia, the Sunda subduction zone lies at the interface between the Sunda block and the Australian plate. Over the last 200 years, there have been at least five giant earthquakes, including the December 26, 2004 earthquake, which caused a tsunami that killed approximately 100,000 people. As such, a detailed understanding how the two plates are slipping past each other can help us understand where and when the next giant earthquake near the Sunda subduction zone is likely to occur.



Over the last 200 years, there have been at least five giant earthquakes, including the December 26, 2004

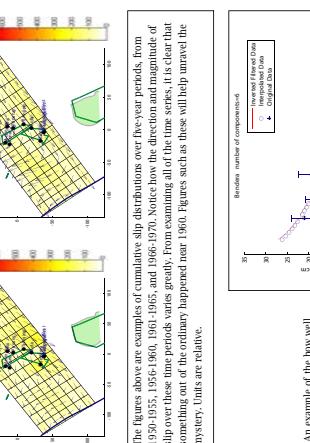


(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After taking into account factors such as a change in sea level height, the two-prong method provides a more accurate position of the crust to determine its relative to one part of the earth that has been taken during the same period. This allows us to calculate a reasonable detailed record of vertical deformation over the last 40 to 50 years.

After

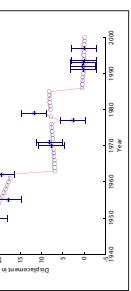
the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

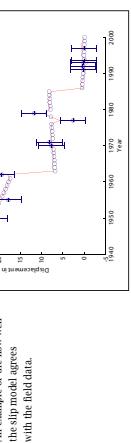
the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

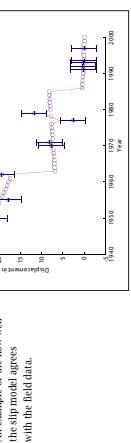
the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

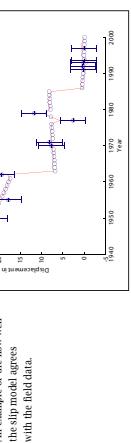
the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

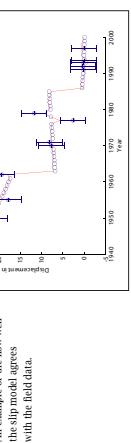
the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

the



(Left) The two-prong method used to measure surface motion due to plate motion. The two prongs are oriented vertically and horizontally, and record growth records, which is a proxy for vertical displacement. While GPS is more precise and provides data in three directions during the 1960 Sumatra subduction zone event, it is less accurate than the GPS station network and is not able to record quakes. (Right) Because crust occurs rarely at a slow growth rate, we can provide over a hundred years of displacement data [5].

After

the



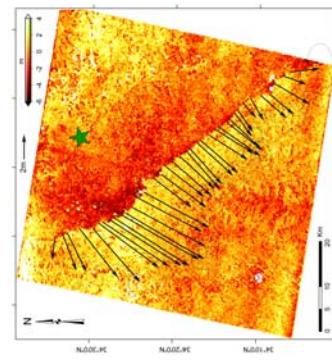
COSI-Corr v1.0 : Co-registration of Optically Sensed Images and Correlation

Toward an operational use of optical remotely sensed images for ground deformation measurements

S. Leprinic, F. Ayoub, J-P Avouac - California Institute of Technology, GPS department
leprinic@caltech.edu

In complement to seismological records, the knowledge of the ruptured fault geometry and of the co-seismic ground deformation are key data to investigate the mechanics of seismic rupture. This information can be retrieved from sub-pixel correlation of pre- and post-earthquake remotely sensed optical images. However, this technique suffers from a number of limitations, mostly due to uncertainties on the imaging systems and on the platform attitudes, leading to strong distortions and stereoscopic effects. Here, we propose an automated procedure that overcomes most of those limitations. In particular, we take advantage of the availability of accurate digital elevation models with global coverage (SRTM). This methodology will improve our ability to collect measurements of ground deformation, in particular in the case of large earthquakes occurring in areas with little or no local geophysical infrastructure. Measuring co-seismic deformations from remotely sensed optical images is attractive thanks to the operational status of a number of imaging programs (SPOT, ASTER, USGS-NAPP aerial programs, etc...) and to the broad availability of archived data.

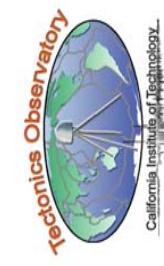
The 2005, Mw 7.6 Kashmir earthquake: Sub-pixel correlation of ASTER images



The color image represents the Northward ground displacements (positive to the north), determined from the correlation of ASTER images, 15-m ground resolution, taken on November 14, 2005 (AST_L1A003/20053267957) and on October 27, 2005 (AST_L1A003/20051572195). The incidence angle is 8.6° for both images. This correlation image was obtained with a sliding 32x32 pixels correlation window and 8-pixel step, leading to a ground resolution of 20 m. No measurement is assigned to white points, where the correlation is lost or where outliers (where the measured ground displacement was found to exceed 10 m) have been filtered out. Correlation is lost mainly due to land slides or variation of the snow cover. Oscillations with an amplitude of 3-mm with a period of 5 km have been removed from stacking. They were characteristic of altitude oscillations from ASTER.

Arrows represent the horizontal surface fault slip. They are defined from linear least square adjustment on each side of the fault and on each NS and EW images of stacked profiles running perpendicular to the rupture. Profiles are stacked over a width of 6 km and a length of 18 km. Ellipses show the 95% confidence intervals.

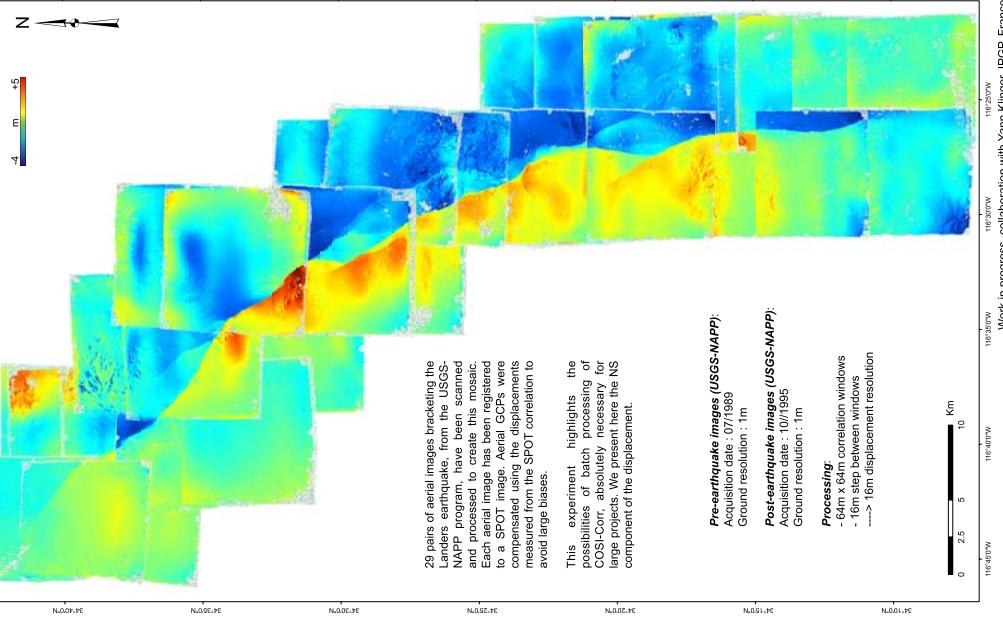
- References:**
- S. Leprinic, S. Barnot, F. Ayoub and J-P Avouac, "Automatic Precise Ortho-correlation and Co-registration for satellite Image Correlation, Application to Ground Deformation Measurements," IEEE Transactions on Geoscience and Remote Sensing, In Press 2006.
 - J. P. Avouac, F. Ayoub, S. Leprinic, O. Koenig and D. Helmberger, "The 2005, Mw 7.6 Kashmir surface rupture kinematics from subpixel correlation of ASTER images and seismic waveforms analysis," Earth and Planetary Science Letters, vol. 249, no. 3-4, pp. 514-528.
 - F. Ayoub, S. Leprinic and J-P. Avouac, "Aerial Photography for Seismic Ground Deformation Measurement," Submitted to ISPRS, 2006.
 - S. Leprinic, A. Samaniego and J-P. Avouac, "Edge Preserving Denoising of Optical Correlation Images, Applications to Ground Deformation Measurement," In preparation.



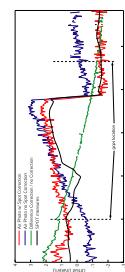
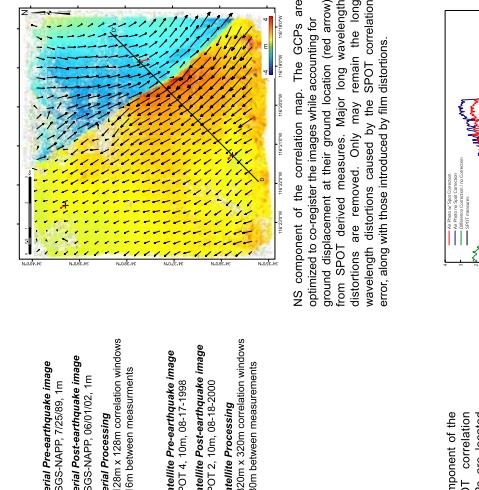
California Institute of Technology

In the general procedure consists of generating accurate ground control points (GCP) for each image. An accurate ortho-rectification model is then built, which allows accurate orthorectification and coregistration of the images. Correlation on the ortho-rectified images then delivers the horizontal ground displacements to analyse. As of November 2006, we are releasing a complete software package, COSI-Corr, developed in IDL and integrated under ENVI that allows for such processing in the following, we demonstrate its main capabilities, namely the processing of SPOT, ASTER and aerial images, the possibility of constraining aerial images with satellite derived measurements, and the convenient use of batch processing for intensive computations.

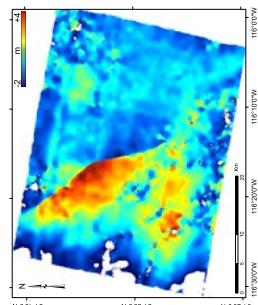
The 1992, Mw 7.3 Landers earthquake from Aerial images, batch processing. SPOT images taken as reference.



The 1999, Mw 7.1 Hector Mine earthquake: Using a priori information from Satellite images to better constrain Aerial images measurements



Profile DD: GCPs correction brings a 2nd order polynomial correction in each NS and EW directions. High frequency signal displacement at the fault is not affected however.



The 1999, Mw 7.1 Hector Mine earthquake: Cross-correlating SPOT and ASTER images

COSI-Corr allows the co-registration and correlation of several kinds of data. Here is an example where the co-seismic displacement of the Hector Mine earthquake is accurately retrieved from the cross-correlation of SPOT and ASTER images. Here, only the NS component of the displacement is shown. The correlation map has been denoised.

Processing:
- Each image is ortho-rectified at 15m
- 48m x 48m correlation windows
- 120m between measurements
- Waves artifacts to the pitch oscillations of ASTER have been removed by stacking

Satellite Pre-earthquake image
SPOT 4, 10m, 08/1998

Satellite Postearthquake image
ASTER, 15m, 04-24-2000

Pre-earthquake images (USGS-NAPP):
Acquisition date : 07/1989
Ground resolution : 1m

Post-earthquake images (USGS-NAPP):
Acquisition date : 10/1995
Ground resolution : 1m

Processing:
- 64m x 64m correlation windows
- 16m between windows
- 16m displacement resolution

Work in progress, collaboration with Yann Klingler, IPGP, France

Three-Dimensional Elastodynamic Simulations of Seismic and Aseismic Slip History of a Planar Strike-Slip Fault

Yi Liu (yil@caltech.edu) and Nadia Lapusta (lapusta@caltech.edu), California Institute of Technology

Abstract

Simulations of spontaneous slip accumulation in three-dimensional (3D) models enjoy a lot of interest because of their ability to clarify earthquake physics. We have been developing a D-methodology for simulating the entire seismic and aseismic slip history of a fault subjected to slow aseismic loading. The algorithm, extended from the 2D study by Lapusta et al. (2000), allows us to resolve all stages of spontaneous slip accumulation in a single computational procedure, including quasi-static nucleation process, dynamic rupture propagation, post-seismic deformation, and aseismic processes throughout the loading period. Simulating long-term aseismic deformation histories while accounting for dynamic effects of occasional earthquakes is quite challenging due to a variety of temporal and spatial scales.

We consider a vertical strike-slip fault embedded in an elastic half-space and governed by rate and steady-state friction. On the fault, there is a seismogenic region, 30 km long and 15 km deep, with steady-state friction. It is surrounded by steady-state strength-strengthening regions that stabilize shear stress (creep) under loading. We observe the following interesting phenomena:

- (1) The simulations produce realistic earthquakes and complicated patterns of interseismic slip. The earthquakes propagate with rupture speeds comparable to the shear wave speed of the surrounding bulk, and have average slip rates of order of 1 m/s. After each large earthquake, there is an accelerated post-seismic creep in the surrounding rate-strengthening regions. During interseismic periods, we observe very interesting patterns of aseismic slip with accelerating and decelerating patches and slow propagation of faster creeps along the surface. These patterns result in occasional small events.
- (2) The quasi-dynamic model which ignores wave-mediated stress changes and hence significantly simplifies the computation of dynamic responses, qualitatively captures most features of the fully dynamic computation but produces more sluggish aseismic behavior and events unable to produce some dynamic features such as the superseismic bursts.
- (3) An asperity (a small circular region 20% stronger than the surrounding fault) causes a supershear burst for the first earthquake in the simulation but not for subsequent events. This indicates that single-earthquake simulations, due to their strong dependence on initial conditions, may in some cases reach conclusions that would be sustained over a longer history of the fault.
- (4) All simulated large events have similar initial stages of their moment-rate function.

In future studies, we plan to (i) adopt more realistic friction laws, by combining rate and state friction with pore pressure evolution and flash heating effects during the dynamic rupture; (ii) incorporate the bimaterial configuration into our earthquake sequence simulations, to investigate its statistical influence on rupture propagation direction over many earthquake cycles; (iii) determine the model response for a wider range of frictional parameters, such as more realistic characteristic slip distance of state and friction; (iv) investigate the possibility of determining frictional parameters by comparing our simulations with observations; (v) study how complicated patterns of aseismic slip that we observe can explain recent observations of slow earthquakes and other interesting phenomena.

Case I: Fault with homogeneous seismic region

Model Geometry

We simulate earthquake sequences on a fault embedded in an infinite elastic half-space, subjected to slow tectonic loading ($f_0 = 35$ min). The fault properties are extended from 2D studies (i.e., Lapusta et al., 2000), where a steady-state velocity-weakening region of $a = 0.015$ and $b = -0.019$ is surrounded by steady-state velocity-strengthening regions of $a = -0.019$ and $b = 0.015$. Nucleation starts in the strip $15 \text{ km} < x < -10 \text{ km}$ as the initial shear stress there is set to be 10% higher than τ_{eq}/τ_0 , $L = 8 \text{ km}$, and the aseismic critical nucleation sizes for an air-plane problems are (Rice & Rummel, 1983; Kubili and Ampuero, 2005).

The calculation is implemented using spectral boundary integral method. Since the analytical internal kernels are available only for the whole infinite space, we make a mirror image of the simulated fault to approximately represent the effect of the free surfaces. We use variable time stepping. Throughout the computation, time steps change by more than 10 orders of magnitude, allowing us to do relatively fast through the quasi-static loading periods, and to consider carefully earthquake nucleation and dynamic rupture propagation periods.

Snapshots of slip rate distribution during 1st and 2nd events

The color scheme depicts the slip rate in our one-dimensional model. Each snapshot has two time marks. The first one is the simulated time t averaged seismic slip rate in years. The second one is the current time step Δt multiplied by 30.

Supershear Burst in the First Event

Snapshots of velocity distribution in the 1st event:



Average rupture speed during the first event is $c = 4.6 \text{ km/s}$ which is larger than $c_s = 4 \text{ km/s}$.

Dunham et al. (2003) studied this kind of supershear bursts in single-event simulations; rupture surrounds the heterogeneity and creates a supershear burst after breaking it.

No Supershear Burst in the Following Events

Snapshots of velocity distribution in the 2nd event:



Average rupture speed during the second event is $c = 4.6 \text{ km/s}$ which is larger than $c_s = 4 \text{ km/s}$.

Dunham et al. (2003) studied this kind of supershear bursts in single-event simulations; rupture surrounds the heterogeneity and creates a supershear burst after breaking it.

Model of a vertical strike-slip fault

Model of a vertical strike-slip fault

In simulations, we use rate-and-state friction law:

$$\tau = \sigma f = c \left[f_c + ab \frac{\dot{V}_e}{V_0} + b \ln \frac{f(t)}{f_0} \right]$$

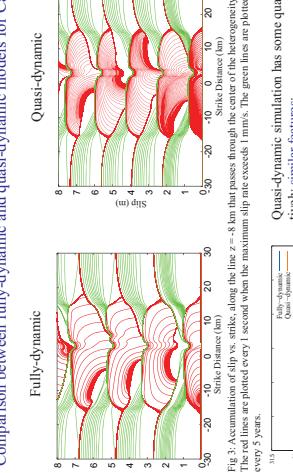
a, b are friction parameters, of the order of 0.01.

If $a < b$, the fault exhibits steady-state velocity weakening;

Basic friction $f_0 = 0.6$ at the sliding velocity $V_e = 10^6 \text{ m/s}$.

Supershear burst due to this heterogeneity is not repeatable in subsequent events due to redistribution of stress.

Comparison between fully-dynamic and quasi-dynamic models for Case II



Quasi-dynamic

Fully-dynamic

Quasi-dynamic simulation has some qualitative similarities to the fully-dynamic simulation:

slow nucleation process

fast rupture propagation

postseismic slip

interseismic slip

slower rupture propagation speed

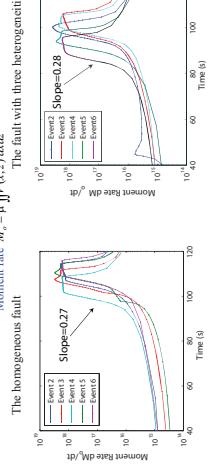
smaller slip rate and step events

smaller earthquake period T

Similar moment rate function for large events

The homogeneous fault

Moment rate $M_r = \mu \int f(x, z) dz$: The fault with three heterogeneities



Fully-dynamic simulation has important differences:

slower nucleation process

fast rupture propagation

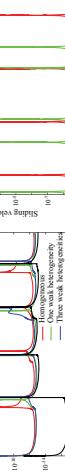
postseismic slip

interseismic slip

slower rupture propagation speed

smaller slip rate and step events

smaller earthquake period T



Different events have almost the same slope of growth S .

What is the physical meaning of the slope? Some discussion in Ampuero (2003).

Can it be inferred from seismic observation?

Influence of L on earthquake period T

Will happen for even smaller L ? Will T become insensitive to the existence of heterogeneities for small enough L ?

An Instantaneous Sub-Rayleigh-to-Supershear Transition Mechanism

Yi Liu¹ and Nadia Lapusta^{1,2}

¹The seismic ratio $S = (\tau^* - \tau^*) / (\tau^* - \tau^*)$. Without loss of generality, we assume $\tau^* = 0$ in simulations.

Rupture propagation is calculated with grid size Δx and time step Δt :

$$\Delta x = \frac{L'}{\sqrt{N}}, \quad \Delta t = \frac{\beta \Delta x}{\beta c_s} \quad (3)$$

3 Burridge-Andrews Supershear Transition Mechanism

Based on a 2D shear wave crack model Burridge (1973) found that there is a shear stress peak $\tau^* = \tau^* + S_{crit}(\frac{x}{L'})$, which propagates with the shear wave speed in front of the crack. Andrews (1976) observed that the shear stress τ^* at the shear wave peak gradually increases as the rupture propagates, and approaches the limiting value τ^* . If $\tau^* > \tau^*$, the seismic ratio $S < S_{crit} \approx 1.03$, τ^* overcomes the fault strength τ^* in the process of rupture propagation, and a daughter crack nucleates at the shear wave front, propagating with a supershear speed.

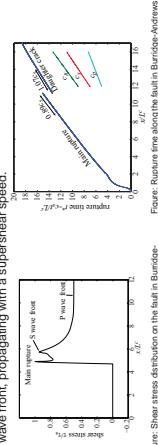


Figure: Shear stress distribution on the fault in Burridge-Andrews model.

4 Super shear transition in our model

In the following, we will show that various other approaches, described in the abstract, are also able to trigger the rupture to go supershear.

4.1 Advancing main rupture toward a preexisting subcritical crack

To smoothly initiate main rupture and a subcritical secondary crack, we impose the following loading stress $\tau^*(x, t')$:

$$\tau^*(x, t') = \tau^* + (\tau^* - \tau^*)[1 + (1 - e^{-1})e^{-\frac{|x|}{L'}}] + \frac{S_{crit}}{c_s^2}t'^2, \quad t' < t_{cut} \quad (4)$$

where $L_{cut} = 0.5(j_0d_0)(1 - \nu)^{0.5}$ is half of the critical nucleation length for in-plane crack obtained by Uenishi and Rice (2003). $D = 12L'$, and $\tau^* = \tau^*/\beta$. This form of loading stress induces two separate cracks at $x = 0$ and $x = 2L'$ at $t' = 0$. At $t' = t_{cut}$, the length of the crack around $x = 0$ reaches the critical length $2L_{cut}$, and it starts to propagate spontaneously, and the length of the crack around $x = D$ is only $1.2L_{cut}$, therefore it remains a subcritical crack. We stop increasing the loading stress at $t' = t_{cut}$, and set $t = t' - t_{cut}$.

We observe the following interesting features in our simulations, which we will present along with our preliminary analysis:

(1) Crack front can abruptly jump from the Rayleigh-wave speed to a supershear speed. We call this "direct" supershear transition. For example, consider a secondary crack nucleated by one of the ways described above under the advancing field of the main rupture. The secondary crack at sub-Rayleigh and it accelerates towards the Rayleigh wave speed. Once the Rayleigh wave speed is reached, the secondary crack jumps to a supershear speed instantaneously.

(2) The supershear transition mechanisms we have described work not only in 2D plane models, but also in 2D models under certain conditions.

(3) Once the transition takes place in our models, the supershear rupture propagation can be maintained under prestress levels that are much lower than ones predicted by the Burridge-Andrews mechanism. This shows that the level of prestress implied by the Burridge-Andrews mechanism is only needed to nucleate a crack at the site of the shear-wave peak, and not to drive the rupture to supershear speeds or to maintain that supershear propagation.

2 Simulated Model (2D)

A planar interface is embedded in an infinite, elastic and homogeneous space. The main rupture initiates from a length of $2L'$ given in (2). In some cases, a heterogeneity exists in front of the main rupture. Depending on simulated problems, the heterogeneity may be a preexisting subcritical crack, a patch with lower static friction strength, or a patch with higher prestress.

The fault strength Γ is assumed to be governed by a linear slip-weakening law:

$$\Gamma(\delta) = \begin{cases} \tau^* + (\tau^* - \tau^*) - (\delta - \delta_c), & \delta \leq \delta_c \\ \tau^* + d, & \delta > \delta_c \end{cases} \quad (1)$$

A singular shear crack with uniform prestress σ^* will propagate spontaneously if its half length exceeds a critical value L' (Andrews (1976)):

$$L' = \frac{2\mu(\lambda + \mu)}{\pi(\lambda + 2\mu)}(\tau^* - \tau^*)^{1/2} \quad (2)$$

the patch the main rupture propagates with sub-Rayleigh speeds. However, the patch (with the length $0.2L'$ in the simulation) induces the rupture to transition to supershear, and the rupture remains supershear afterward. This shows that the level of prestress implied by the Burridge-Andrews mechanism is not required to maintain real faults propagation. Hence, supershear propagation on real faults can occur under prestresses that are much lower than the values implied by the Burridge-Andrews mechanism.

4.3 A crack under over-stressing condition

Supershear transition can be induced by overstressing a crack. Consider the case with prestress τ^* over the region $x \in [-L', L']$ set to 0.8 times larger than the static friction strength τ^* . At the beginning of the simulation, the stress inside the patch drops from $1.8\tau^*$ to 0 instantaneously, and spontaneous rupture propagation starts. The rupture is initially sub-Rayleigh.

An interesting phenomenon is that both our abrupt supershear transition and the Burridge-Andrews daughter crack transition appear in the simulation. First, we observe our abrupt supershear transition without initiating a daughter crack (in the figure below). Later we observe a daughter crack that nucleates in front of the main rupture and propagates with supershear speeds (in the figure below).

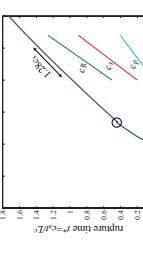
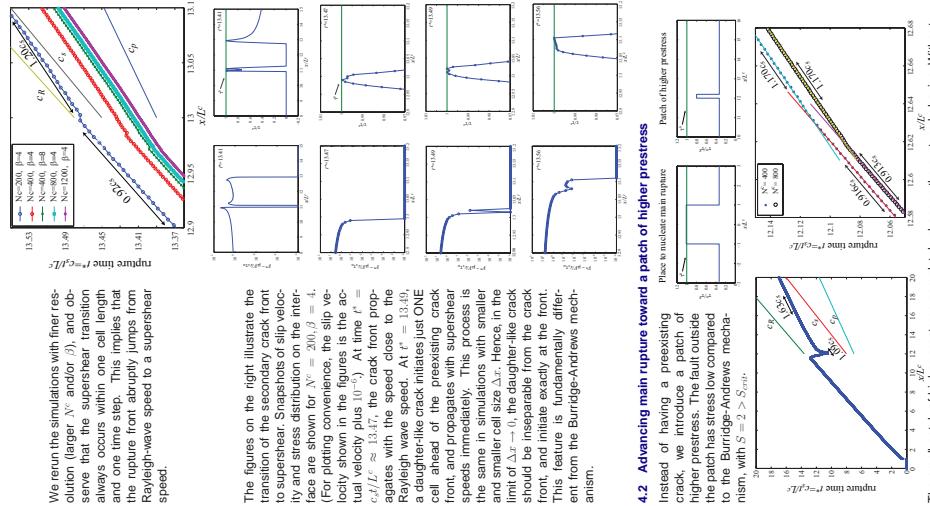


Figure: Rupture time along the fault for the case $N = 100, \beta = 1$.

5 Discussion

The described abrupt supershear transition mechanisms work in 3D fault models as well. We simulate the rupture propagation on a strike-slip fault interface, where a rectangular fault is surrounded by unbreakable barriers. If the model includes a patch of higher prestress or lower static friction strength, the rupture may transition to supershear speeds and the supershear can be maintained in spite of the prestress much lower than that predicted by the Burridge-Andrews mechanism. However, we notice that to trigger supershear transitions, the required patch size in the 3D strike-slip model should be much larger than that in the 2D in-plane model.

From simulations, it seems that a special loading stress τ^* is needed to cause cracks to transition to supershear. We hypothesize that the loading stress should move fast enough in the direction of the crack propagation. The most natural loading stress environment of this kind is the stress field in front of main rupture propagating with supershear Rayleigh speed advancing on a secondary crack or heterogeneity. However, there are other ways to create suitable loading stress environments. For example, we have tried to artificially impose a dynamic loading stress field $\tau^*(x, t) = f(t - c_s t)$ on a crack propagating with sub-Rayleigh speeds. We find that this also triggers supershear transition with features very similar to the preexisting crack case (4.1). Our current work is directed towards developing theoretical explanations for these phenomena.



The small patch of higher prestress completely changes the rupture behavior. Without

POLYGENETIC POST-LARAMIDE MANTLE LITHOSPHERE BENEATH THE MOJAVE DESERT: THE XENOLITH RECORD

Peter I. LUFFI, Jason B. SALEEBY, Minal N. DUCEA
Division of Geological & Planetary Sciences, California Institute of Technology,²Department of Geosciences, University of Arizona

*e-mail: pluffi@gps.ca/tech.edu

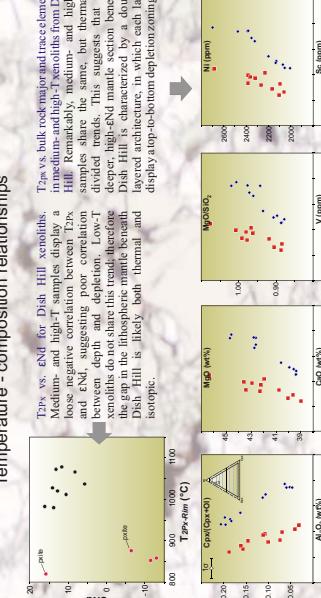


1. Introduction

1) Nd and Pb isotope composition of mantle-derived sills from the Mojave Desert reveal regional scale signatures of the deep lithosphere in the Miocene (Hart et al., 2000; GSA Bull.) enriched/Precambrian sub-continental lithospheric mantle (SCLM) of North America underlain the eastern part of the Desert, whereas underlain the western part of the Desert, enriched/Early Cenozoic mantle dominated its composition (Fig. 1). The boundary between these domains coincides with the SR-16 line of Kuster and Peterson (1973, GSA Bull.). In the eastern part of the Mojave Desert model is sustained by trace element and radiogenic isotope enrichment peridotite xenoliths emplaced in the Cima volcanic field in Plio-Pleistocene-Holocene basins (Dish Hill, 1997; GSA Bull.). Paleogeographic models (Milner et al., 1996) predict a bounding fault zone between the eastern and western domains, and therefore the eastern domain is bounded by the San Andreas fault which can help to refine the boundary between the inferred main domains, and 2) relationships between the inferred main domains, and 2) relationship between the inferred main domains, and 2) relationship between the inferred main domains, and 2)

Temperature (°C)	Sn90Ag10 (wt%)	Sn90Cu10 (wt%)
100	90.0	90.0
200	89.5	89.0
300	88.5	87.5
400	87.0	85.5
500	85.5	83.5
600	84.0	81.5
700	82.5	79.5
800	81.0	78.0
900	79.5	76.5
1000	78.0	75.0
1100	76.5	73.5

2. Dish Hill xenoliths



Information about the Majestic Document

E. S. Mazzoni

- Mantle xenoliths
- Juxtaposition of American lithosphere and Farallon slab beneath the Late Pleistocene constrains geodynamic evolution of the United States.
- Compositional differences suggest imbrication in the Late Pleistocene.
- Adiabatic bathe fractionating melt environment may trend observed in

MENTS
ACKNOWLEDGMENTS
ASimow
observator

THE "SIMEULE SADDLE"; AN ASEISMIC PATCH OF THE SUnda MEGATHRUST?

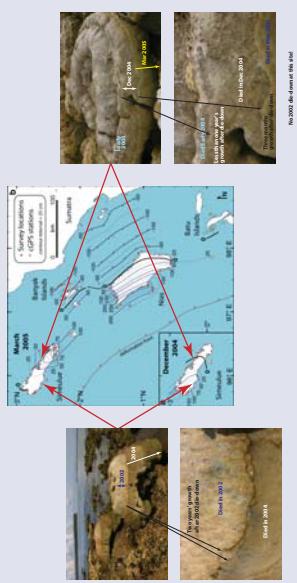
ron Meltzner, Kerry Sieh, Rich Briggs, John Galetzka, Willy Amidon (Tectonics Observatory, California Institute of Technology), Lambang Suwardi, Danny Hilman Natawidjaja, Dudi Pravudi (Indonesian Institute of Sciences - LIPI Geotechnology), and Imam Suprihanto

INTRODUCTION: THE QUESTIONS

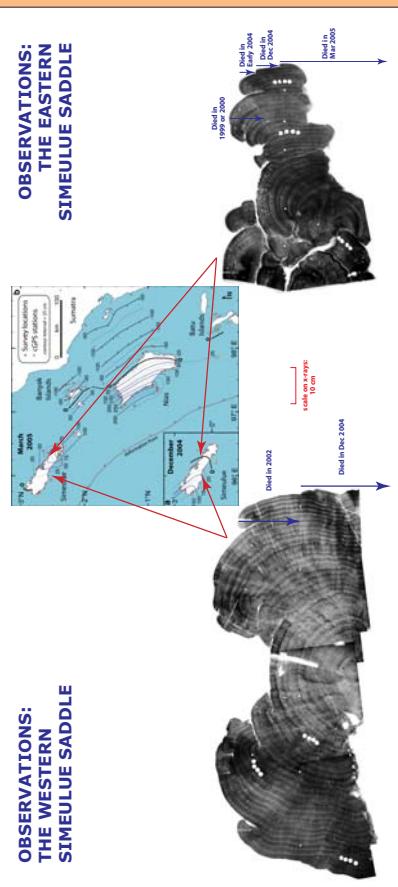
The December 2004 and March 2005 Sunda megathrust earthquakes nucleated northwest and southeast of Simbulua Island, respectively, and each ruptured sequentially into the ~100-km-long island. Uplift at the northwest end of Simbulua during the 2004 earthquake, and uplift at the southeast end up to 2005 were 1.5 m or more. Uplift associated with each earthquake diminished toward the center of the island, and did not dip on the underlying megathrust, according to our inversion. Cumulative uplift was as little as 0.5 m along the west coast of central Simbulua. Hence, although there is an uplift deficit on central Simbulua, it is not as large as the ~1.5 m of subsidence produced during an M_w 7.3 earthquake in 2002 produced up to ~200 m west of the sites of lowest uplift in 2004–2005, even including that uplift, the saddle persists. No events similar to 2002 can be found in the historical record for at least the previous 54 years.

The occurrence of major earthquakes during the 2002–2004 and 2005–2006 field seasons and of X-ray images of slabbed Porites coral microatolls, coupled with a paucity of large historical earthquakes in the region, suggests that the central Banda Sea region may experience occasional seismic slip events, or that moderate earthquakes ($M > 7$) may account for significant fracturing of the uplift during a

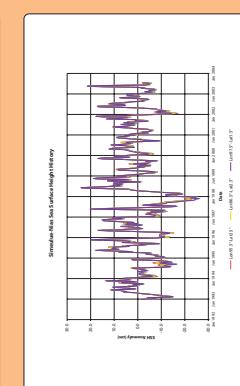
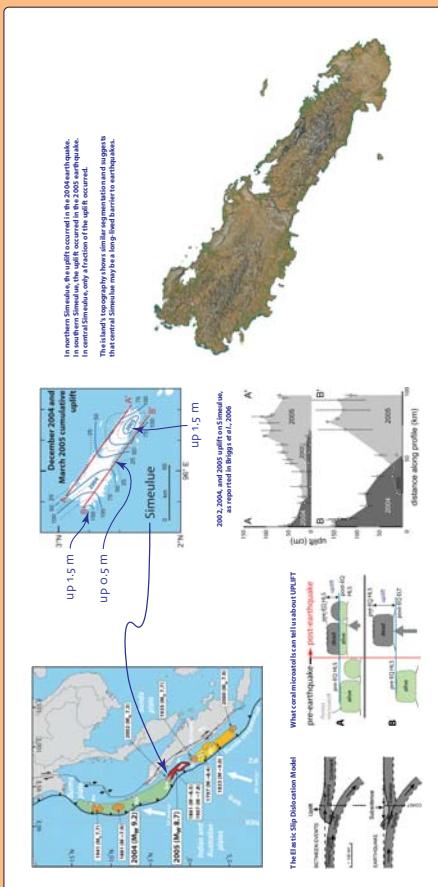
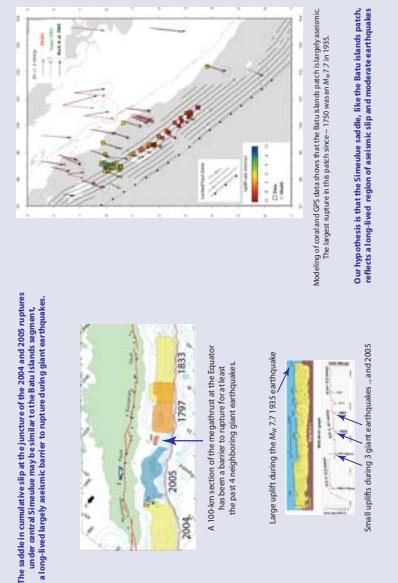
OBSERVATIONS: THE WESTERN SIMEULUE SADDLE



OBSERVATIONS: THE EASTERN SIMEULUE SADDLE



**OBSERVATIONS:
THE EASTERN
SIMEULUE SADDLE**



Phosphorus Zoning in Olivine: An Additional Constraint on Magmatic Processes in Arc Volcanism

Zachary T. Morgan, Edward M. Stolper, Michael B. Baker, Daniel Vielzeuf, and Fidel Costa

Introduction/Context

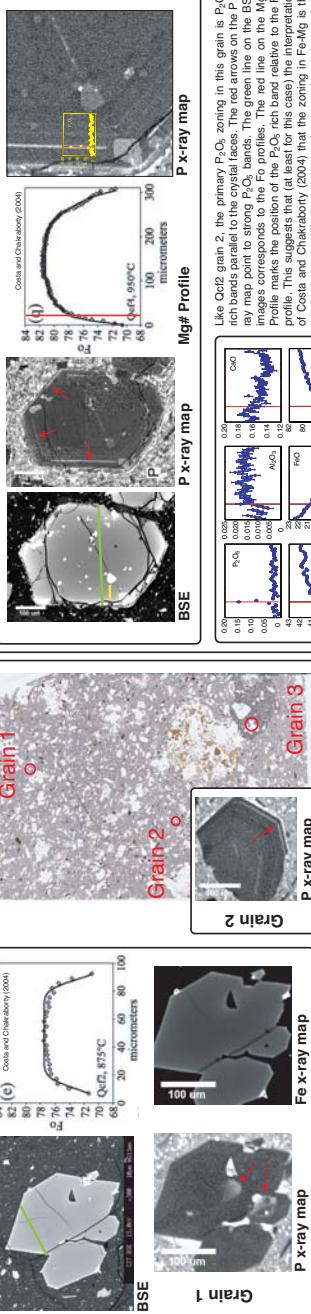
Arc lavas contain abundant evidence of disequilibrium processes, in the form of concentration profiles, reaction rims, and dissolution features. Many of these disequilibrium features can be used to determine the time scales of the magmatic events resulting in their formation. Previous studies have used zoning in plagioclase and/or olivine, reaction rims on olivine, and amphibole break down rims to infer the time from the disequilibrium event to eruption.

Previous work by Fidel Costa used zoning in olivine or plagioclase to infer the time scales of magmatic processes. Costa and Chakraborty (2004) looked at Fe-Mg zoning in olivine to calculate the time scales of olivine incorporation into dacite and andesite magmas. They focus on four lavas from Volcan San Pedro in the Southern Volcanic Zone of Chile (a and b). The lavas are three dacites (Qcf1, Qcf2, and Qcf3), and one andesite (Qcf4). None of these magmas has a low-pressure olivine phase volume so its presence is assumed to be the result of magma mixing with a more mafic magma. (i.e., basaltic andesite, Costa and Chakraborty, 2004) or assimilation of mafic to ultramafic lithologies (Costa and Dungan, 2005). The Fo (Mg#) profiles of olivine reported in Costa and Chakraborty (2004) are reproduced for each of the studied lavas. As illustrated in the Fo (Mg#) zoning can be complicated but then modelled on the outermost zone developed during the last stage prior to eruption. This is year 1 for Qcf1 to 6 years for Qcf2, and 10 years for Qcf3 and Qcf4.

Recent studies of zoning in olivine suggest that just focusing on the major elements such as a FeO and MgO profile is insufficient. Certain trace elements such as the slowly diffusing $P_{2}O_5$, Al_2O_3 , and Cr_2O_3 may record important information about the olivine growth history that is not retained by the faster diffusing species. Milman-Barris et al. (2006) described complex zoning of $P_{2}O_5$, Al_2O_3 , and Cr_2O_3 in natural olivines from Hawaiian lavas, komatiites, and a Martian meteorite. In addition, their linear cooling rate experiments on a Hawall bulk composition produced olivine with strong zoning in $P_{2}O_5$, Al_2O_3 , and Cr_2O_3 . They attributed the $P_{2}O_5$, Al_2O_3 , and Cr_2O_3 zoning to rapid growth of the olivine.

The purpose of this study is to examine olivine from arc lavas for zoning patterns, and to explore the effects this zoning will have on the mode and time scales estimated by Costa and Chakraborty (2004).

Qcf2: Dacite, $SiO_2 = 64.5\%$, $T = 875\text{ C}$



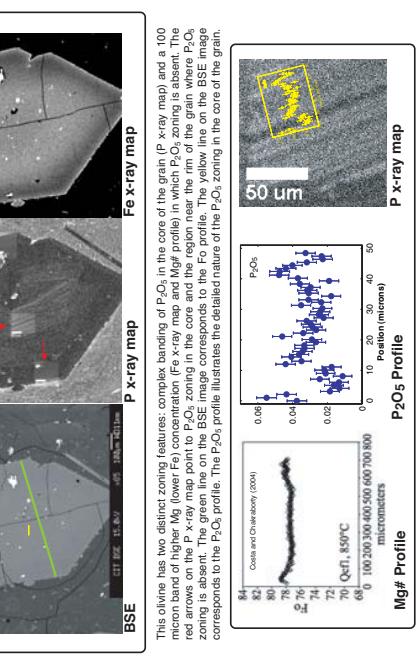
Grain 1: Olivine zoning in Qcf2 dacite

Grain 1: This olivine has two distinct zoning features: a higher $P_{2}O_5$ concentration (outer 10 microns, Fe-X ray map) where $P_{2}O_5$ zoning is absent and $P_{2}O_5$ rich zones in the interior of the grain. The red arrows on the P-X ray Map point to the $P_{2}O_5$ zoning in the core of the grain. The green line on the BSE image corresponds to the Fo profile.

Grain 2: The $P_{2}O_5$ zoning in this olivine grain occurs as bands parallel to the crystal faces.

Grain 3: This grain occurs in a more mafic inclusion in the dacitic magma. The $P_{2}O_5$ zoning in this olivine is a combination of patchy $P_{2}O_5$ -rich zones in the rim and a band parallel to the crystal faces. Similar to the Qcf1, there is a 100 microns thick $P_{2}O_5$ zoning in the grain, with a higher Fo (Mg#) value in the core than in the outermost $P_{2}O_5$ zoning. The red arrow on the BSE image points to features in the $P_{2}O_5$ zoning region. The green line on the BSE image corresponds to the Fo profile.

Qcf1: Dacite, $SiO_2 = 65.8\%$, $T = 850\text{ C}$



Grain 1: Olivine zoning in Qcf1 dacite

This olivine has two distinct zoning features: complex banding of $P_{2}O_5$ in the core of the grain (P-X ray map) and a 100 microns band of higher Mg# lower Fo point to $P_{2}O_5$ zoning in the core and the region near the rim of the grain where $P_{2}O_5$ zoning is absent. The red arrows on the P-X ray map point to the core and the region near the rim on the BSE image corresponds to the Fo profile. The yellow line on the BSE image corresponds to the $P_{2}O_5$ profile. The $P_{2}O_5$ profile illustrates the detailed nature of the $P_{2}O_5$ zoning in the core of the grain.

The phenanitized P-X ray map illustrates the $P_{2}O_5$ banding, on which I've superimposed a quantitative profile taken across a portion of the crystal. The quantitative profile illustrates the $P_{2}O_5$ concentrations that are qualitatively described by the P-X ray map.

Additional evidence for our hypothesis of a later stage rim on the olivine comes from spinel inclusions hosted within the olivine grain. The Fe, Ti, Al and P-X ray maps illustrate that spinel inclusions on the outer edge of the grain have high Fe, Ti and P and low Al relative to spinel inclusions in the core of the lava.

Summary/Future Work:

As with previously examined olivine from other tectonic environments (Milman-Barris et al. 2006), olivines from arc lavas contain complex zoning in $P_{2}O_5$. The types of $P_{2}O_5$ zoning in olivine we have observed are bands parallel to the crystal faces (Qcf2 grains 1 and 2, Qcf3 grain 2, and Qcf4), bands in the interior of the grain (Qcf2 grain 1, and Qcf3). Ultimately, the complex zoning will better constrain the growth histories of the olivine, and lead to a greater understanding of the processes involved in nucleation and emplacement events in arc lavas. However, the shorter term the $P_{2}O_5$ zoning coupled with Fe-Mg zoning can be used to refine models and estimates for olivine residence times in lavas, as is illustrated with Qcf4. The next step of this study is to collect P-X ray maps for a population of different olivine in the same section to characterize the $P_{2}O_5$ zoning for each lava. Are the types of $P_{2}O_5$ zoning in olivine distinct for a given lava or is there a continuum of $P_{2}O_5$ zoning common to all of the lavas?

In addition, we will modify the model used by Costa and Chakraborty (2004) to estimate the olivine residence times to include a rim of new olivine similar to Qcf4. What is the effect on time estimates when a rim of new olivine is added to a pre-existing olivine? Are the changes to the olivine residence times consistent for all olivine in a given lava?

References:
Costa, F., and Chakraborty, S. (2004). Olivine zoning in arc lavas: implications for melt evolution and melt migration. *Journal of Petrology*, 45, 517-530.
Costa, F., Dungan, M. (2005). Short timescales of magmatic evolution in arc lavas. *Geology*, 33, 837-840.
Milman-Barris, M.S., Milman, A.E., Baker, M., Stolper, E.M. (2006). Zoning of phosphate in magnetite inclusions. *American Mineralogist*, 91, 1161-1170.

H e a v e n l y M o u n t a i n s , D o w n - t o - E a r t h J o b

BY ELISABETH NADIN

Tien Shan is Chinese for “heavenly mountains”—the name alone reflects how hard it is to get there. In June 2006 a group of 21 Caltech students, postdocs, and professors made the trip to this remote region of northwestern China as participants in a two-week research field trip, sponsored by the Institute’s Division of Geological and Planetary Sciences. After landing in Beijing, we flew to Urumqi, capital of the Xinjiang Province—where the range is located—and caught up with the 15 Chinese students and four professors who had traveled 48 hours by train from the east-coast city of Nanjing to join us. A 10-hour drive to the southeast Tien Shan foothills brought us to this arid central Asian landscape of red, yellow, and brown rocks cut by tectonic rivers.



Top to bottom:
Jean-Philippe Avouac discusses the Quilitak anticline.
An Uyghur merchant passes tilted rock beds.
A spectacular scene of the stark wilderness with Paleozoic rocks in the background.
At right, Jean-Philippe Avouac explains the interactions between the tilted rock layers and overlying terraces as Nanjing University geology professor Shangli Wang looks on.



The banner at the top of the page shows the Tien Shan’s Quilitak anticline, a large fold that formed over a blind thrust. Fault-bend folds formed the one of the field investigations. Multibeam in the foreground give a sense of how enormous it is.

Below that, a closeup of one of the limbs of the fold, showing sheared thrusts offsets (tops to left) consistent with the compressive field.

Above: A map of eastern Asia—red stars show our short stops, and the blue star is on the Tien Shan.

Left: Satellite image centered on the Tien Shan.



Middle left: Coal mining in the Tien Shan.

Middle right: Nanjing University student Yao Xue and Caltech undergrad Diana Hsiao check the quality of seismic images coming from geophones into a field computer.

Above: Growth strata!

Left: A nice meal hosted by the Tien oil company.



Top: Our happy group at the end of a field day.

Above right: Uyghur men and their delicious fruit.

Left: Caltech geophysicist Rob Clayton, with the help of two students, triggers seismic waves with the jetty gun. Below: Built along the Silk Road in the third century and abandoned in the tenth century when Islam became a dominant force in the region, these hillside Buddhist caves known as the Kizil grottoes still stand today, although many paintings from their ceilings and walls were plundered for their gold gilt or removed by archaeologists.



Top left: Min Chen collects samples as Rob Clayton looks on.

Top right: Our happy home for two weeks.

Middle left: Coal mining in the Tien Shan.

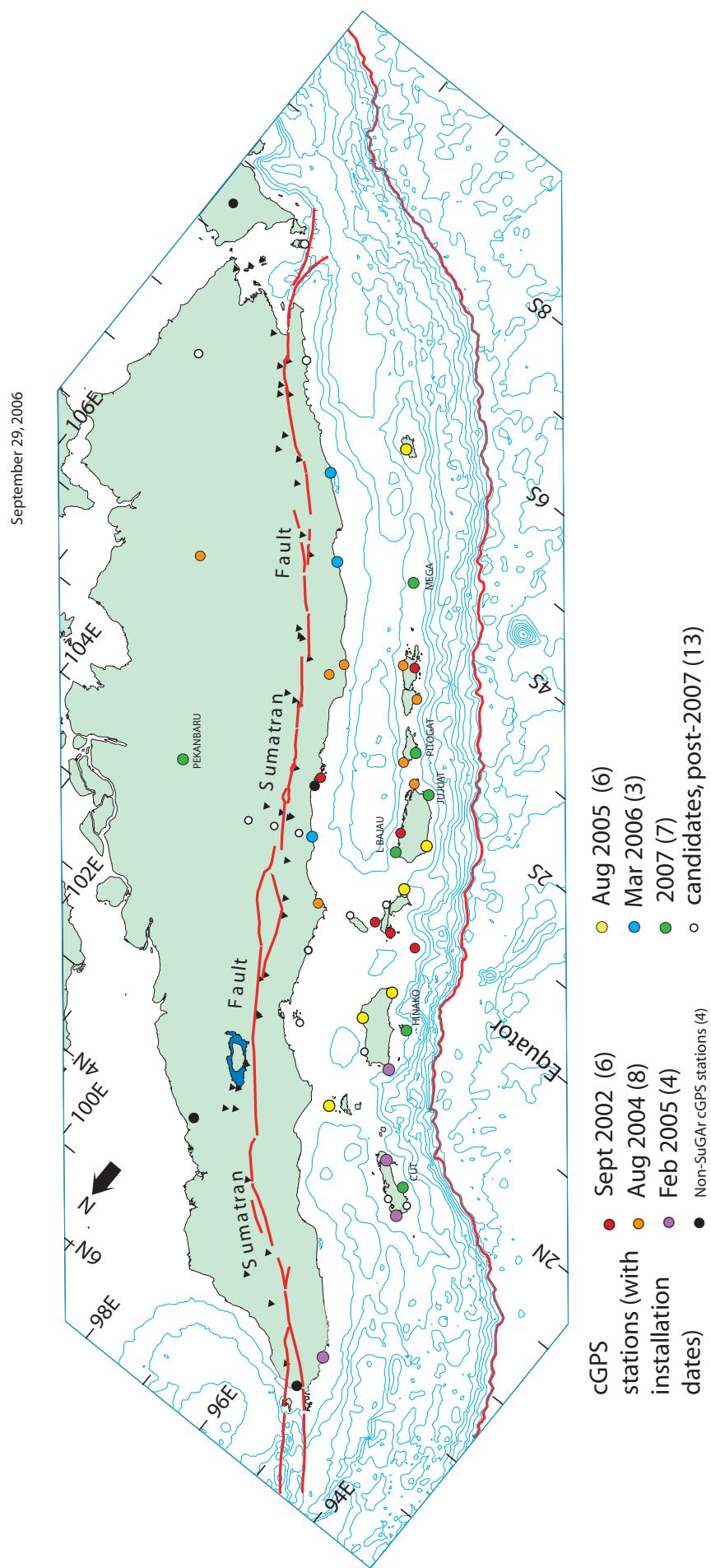
Middle right: Nanjing University student Yao Xue and Caltech undergrad Diana Hsiao check the quality of seismic images coming from geophones into a field computer.

Above: Growth strata!



Left: A nice meal hosted by the Tien oil company.





Assimilation of Plate Tectonic Reconstructions into Geodynamic Flow Models



Mark Turner, Mike Gurnis, Lydia Taylor, Vlad Manea, Sonja Kisin
Caltech, Tectonics Observatory

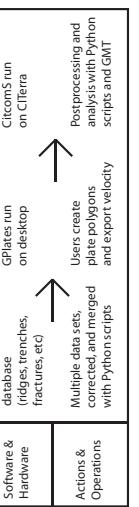
Introduction

One of our goals is to attempt to place the tectonic features of individual boundaries within a global context and understand the way they are moving differently. There is a large number of new subduction zones identified in the new version of CitoNS that are not included in the older version. In order to address these, their migration has been developed, using an entirely new generation of tools that are computationally advanced while being consistent with the actual structure and kinematics of plate boundaries. Thus far we have made considerable progress in this direction.

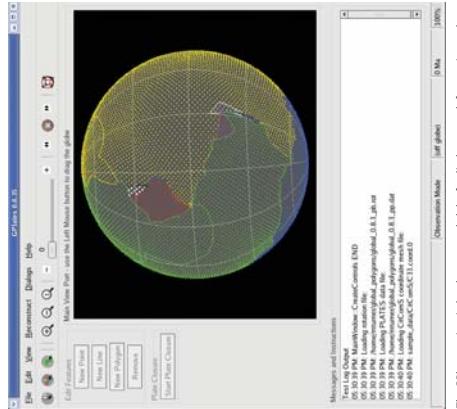
One goal is to assimilate plate tectonic reconstructions into global and regional geodynamic models. Within the University of Sydney and the Geological Survey of Norway, the CitoNS has been a key partner in the development of GPlates, a plate tectonic reconstruction software package. We are using GPlates as the preprocessor from and to global models of mantle convection using the CitoNS finite element code.

Using GPlates, we have developed a method for representing the evolving geometry of tectonic plates. A single plate is represented by all the margins around the plate reconstructed according to the Euler pole of the margin, and an algorithm for computing the intersections of all the margins. This is an essential prerequisite for merging dynamic models with paleo-reconstructions. We have used this software to build a global set of plates over the last 80 million years, to merge these data into a single global circulation model of the mantle. This work has given us a new tool that allows us to explore the dynamics of changes in plate motions and shapes over the next several years. We are now routinely running this software on the CfTera supercomputer.

Typical Workflow Assimilating Data into Models



GPlates: Software for Tectonic Reconstructions



An idealized tectonic plate showing component margin line data features. Each line feature has its own euler pole and rotates according to the rules of plate tectonics.

$$(1) \quad \nabla = \nabla$$

$$(2) \quad \nabla = \partial/\partial\theta - \nabla_\perp$$

$$(3) \quad \frac{\partial}{\partial t} = \dot{\theta}\nabla + \nabla^2\theta + H$$

where u is velocity, σ is the stress tensor, P is dynamic pressure, T is temperature, K is thermal diffusivity, H internal heat and g is gravitational acceleration.

These equations are solved with CitoNS-py (Tan et al., 2006). With the finite element, the model domain is a spherical shell representing the entire mantle and lithosphere. CitoNS-py uses a decomposition scheme such that the spherical shell is first decomposed into 2 caps so that it may view the elements are approximately equal over the entire surface of the sphere. There each cap is further divided such that the edges of the cap are equally divided.

Orthographic projection of processors from a full CitcomS mesh in which there are 15 caps in map view for each cap. The CitcomS cap is shown as distinct colors while the processor domains within the caps are indicated by the intensity of the color. This example was produced for a run with 2 processors in radius such that the total number of processors was $2 \times 160 = 384$. This is the largest case we have solved on the CfTera machine so far and most of the cases have been solved with 96 processors with 128x128 elements in map view for each processor.

References:

Tan, E., Choi, P., Hourcade, M., Gurnis, M., Alvarez-Gaume, Coupling multiple models of mantle convection within a computational framework, Geophysics, Geochimica et Cosmochimica Acta, 70(6), 1029-1039, doi:10.1016/j.gca.2005.11.014, 2006.

CitcomS: Mantle Convection Models

The plate tectonic reconstructions are assimilated into a model of mantle convection solved with the finite element package CitcomS-py developed at Caltech. CitcomS solves for conservation of mass, momentum and energy.

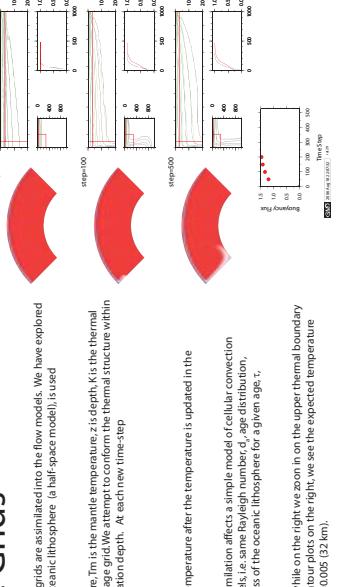
(1) $\nabla = \nabla$

(2) $\nabla = \partial/\partial\theta - \nabla_\perp$

(3) $\frac{\partial}{\partial t} = \dot{\theta}\nabla + \nabla^2\theta + H$

where u is velocity, σ is the stress tensor, P is dynamic pressure, T is temperature, K is thermal diffusivity, H internal heat and g is gravitational acceleration.

These equations are solved with CitoNS-py (Tan et al., 2006). With the finite element, the model domain is a spherical shell representing the entire mantle and lithosphere. CitoNS-py uses a decomposition scheme such that the spherical shell is first decomposed into 2 caps so that it may view the elements are approximately equal over the entire surface of the sphere. There each cap is further divided such that the edges of the cap are equally divided.



References:

Tan, E., Choi, P., Hourcade, M., Gurnis, M., Alvarez-Gaume, Coupling multiple models of mantle convection within a computational framework, Geophysics, Geochimica et Cosmochimica Acta, 70(6), 1029-1039, doi:10.1016/j.gca.2005.11.014, 2006.

case27 (Age=0.0 Ma, Steps=500)

case27 (Age=0.0 Ma, Steps=600)

case27 (Age=0.0 Ma, Steps=700)

case27 (Age=0.0 Ma, Steps=800)

case27 (Age=0.0 Ma, Steps=900)

case27 (Age=0.0 Ma, Steps=1000)

case27 (Age=0.0 Ma, Steps=1100)

case27 (Age=0.0 Ma, Steps=1200)

case27 (Age=0.0 Ma, Steps=1300)

case27 (Age=0.0 Ma, Steps=1400)

case27 (Age=0.0 Ma, Steps=1500)

case27 (Age=0.0 Ma, Steps=1600)

case27 (Age=0.0 Ma, Steps=1700)

case27 (Age=0.0 Ma, Steps=1800)

case27 (Age=0.0 Ma, Steps=1900)

case27 (Age=0.0 Ma, Steps=2000)

case27 (Age=0.0 Ma, Steps=2100)

case27 (Age=0.0 Ma, Steps=2200)

case27 (Age=0.0 Ma, Steps=2300)

case27 (Age=0.0 Ma, Steps=2400)

case27 (Age=0.0 Ma, Steps=2500)

case27 (Age=0.0 Ma, Steps=2600)

case27 (Age=0.0 Ma, Steps=2700)

case27 (Age=0.0 Ma, Steps=2800)

case27 (Age=0.0 Ma, Steps=2900)

case27 (Age=0.0 Ma, Steps=3000)

case27 (Age=0.0 Ma, Steps=3100)

case27 (Age=0.0 Ma, Steps=3200)

case27 (Age=0.0 Ma, Steps=3300)

case27 (Age=0.0 Ma, Steps=3400)

case27 (Age=0.0 Ma, Steps=3500)

case27 (Age=0.0 Ma, Steps=3600)

case27 (Age=0.0 Ma, Steps=3700)

case27 (Age=0.0 Ma, Steps=3800)

case27 (Age=0.0 Ma, Steps=3900)

case27 (Age=0.0 Ma, Steps=4000)

case27 (Age=0.0 Ma, Steps=4100)

case27 (Age=0.0 Ma, Steps=4200)

case27 (Age=0.0 Ma, Steps=4300)

case27 (Age=0.0 Ma, Steps=4400)

case27 (Age=0.0 Ma, Steps=4500)

case27 (Age=0.0 Ma, Steps=4600)

case27 (Age=0.0 Ma, Steps=4700)

case27 (Age=0.0 Ma, Steps=4800)

case27 (Age=0.0 Ma, Steps=4900)

case27 (Age=0.0 Ma, Steps=5000)

case27 (Age=0.0 Ma, Steps=5100)

case27 (Age=0.0 Ma, Steps=5200)

case27 (Age=0.0 Ma, Steps=5300)

case27 (Age=0.0 Ma, Steps=5400)

case27 (Age=0.0 Ma, Steps=5500)

case27 (Age=0.0 Ma, Steps=5600)

case27 (Age=0.0 Ma, Steps=5700)

case27 (Age=0.0 Ma, Steps=5800)

case27 (Age=0.0 Ma, Steps=5900)

case27 (Age=0.0 Ma, Steps=6000)

case27 (Age=0.0 Ma, Steps=6100)

case27 (Age=0.0 Ma, Steps=6200)

case27 (Age=0.0 Ma, Steps=6300)

case27 (Age=0.0 Ma, Steps=6400)

case27 (Age=0.0 Ma, Steps=6500)

case27 (Age=0.0 Ma, Steps=6600)

case27 (Age=0.0 Ma, Steps=6700)

case27 (Age=0.0 Ma, Steps=6800)

case27 (Age=0.0 Ma, Steps=6900)

case27 (Age=0.0 Ma, Steps=7000)

case27 (Age=0.0 Ma, Steps=7100)

case27 (Age=0.0 Ma, Steps=7200)

case27 (Age=0.0 Ma, Steps=7300)

case27 (Age=0.0 Ma, Steps=7400)

case27 (Age=0.0 Ma, Steps=7500)

case27 (Age=0.0 Ma, Steps=7600)

case27 (Age=0.0 Ma, Steps=7700)

case27 (Age=0.0 Ma, Steps=7800)

case27 (Age=0.0 Ma, Steps=7900)

case27 (Age=0.0 Ma, Steps=8000)

case27 (Age=0.0 Ma, Steps=8100)

case27 (Age=0.0 Ma, Steps=8200)

case27 (Age=0.0 Ma, Steps=8300)

case27 (Age=0.0 Ma, Steps=8400)

case27 (Age=0.0 Ma, Steps=8500)

case27 (Age=0.0 Ma, Steps=8600)

case27 (Age=0.0 Ma, Steps=8700)

case27 (Age=0.0 Ma, Steps=8800)

case27 (Age=0.0 Ma, Steps=8900)

case27 (Age=0.0 Ma, Steps=9000)

case27 (Age=0.0 Ma, Steps=9100)

case27 (Age=0.0 Ma, Steps=9200)

case27 (Age=0.0 Ma, Steps=9300)

case27 (Age=0.0 Ma, Steps=9400)

case27 (Age=0.0 Ma, Steps=9500)

case27 (Age=0.0 Ma, Steps=9600)

case27 (Age=0.0 Ma, Steps=9700)

case27 (Age=0.0 Ma, Steps=9800)

case27 (Age=0.0 Ma, Steps=9900)

case27 (Age=0.0 Ma, Steps=10000)

case27 (Age=0.0 Ma, Steps=10100)

case27 (Age=0.0 Ma, Steps=10200)

case27 (Age=0.0 Ma, Steps=10300)

case27 (Age=0.0 Ma, Steps=10400)

case27 (Age=0.0 Ma, Steps=10500)

case27 (Age=0.0 Ma, Steps=10600)

case27 (Age=0.0 Ma, Steps=10700)

case27 (Age=0.0 Ma, Steps=10800)

case27 (Age=0.0 Ma, Steps=10900)

case27 (Age=0.0 Ma, Steps=11000)

case27 (Age=0.0 Ma, Steps=11100)

case27 (Age=0.0 Ma, Steps=11200)

case27 (Age=0.0 Ma, Steps=11300)

case27 (Age=0.0 Ma, Steps=11400)

case27 (Age=0.0 Ma, Steps=11500)

case27 (Age=0.0 Ma, Steps=11600)

case27 (Age=0.0 Ma, Steps=11700)

case27 (Age=0.0 Ma, Steps=11800)

case27 (Age=0.0 Ma, Steps=11900)

case27 (Age=0.0 Ma, Steps=12000)

case27 (Age=0.0 Ma, Steps=12100)

case27 (Age=0.0 Ma, Steps=12200)

case27 (Age=0.0 Ma, Steps=12300)

case27 (Age=0.0 Ma, Steps=12400)

case27 (Age=0.0 Ma, Steps=12500)

case27 (Age=0.0 Ma, Steps=12600)

case27 (Age=0.0 Ma, Steps=12700)

case27 (Age=0.0 Ma, Steps=12800)

case27 (Age=0.0 Ma, Steps=12900)

case27 (Age=0.0 Ma, Steps=13000)

case27 (Age=0.0 Ma, Steps=13100)

case27 (Age=0.0 Ma, Steps=13200)

case27 (Age=0.0 Ma, Steps=13300)

case27 (Age=0.0 Ma, Steps=13400)

case27 (Age=0.0 Ma, Steps=13500)

case27 (Age=0.0 Ma, Steps=13600)

case27 (Age=0.0 Ma, Steps=13700)

case27 (Age=0.0 Ma, Steps=13800)

case27 (Age=0.0 Ma, Steps=13900)

case27 (Age=0.0 Ma, Steps=14000)

case27 (Age=0.0 Ma, Steps=14100)

case27 (Age=0.0 Ma, Steps=14200)

case27 (Age=0.0 Ma, Steps=14300)

case27 (Age=0.0 Ma, Steps=14400)

case27 (Age=0.0 Ma, Steps=14500)

case27 (Age=0.0 Ma, Steps=14600)

case27 (Age=0.0 Ma, Steps=14700)

case27 (Age=0.0 Ma, Steps=14800)

case27 (Age=0.0 Ma, Steps=14900)

case27 (Age=0.0 Ma, Steps=15000)

case27 (Age=0.0 Ma, Steps=15100)

case27 (Age=0.0 Ma, Steps=15200)

case27 (Age=0.0 Ma, Steps=15300)

case27 (Age=0.0 Ma, Steps=15400)

case27 (Age=0.0 Ma, Steps=15500)

case27 (Age=0.0 Ma, Steps=15600)

case27 (Age=0.0 Ma, Steps=15700)

case27 (Age=0.0 Ma, Steps=15800)

case27 (Age=0.0 Ma, Steps=15900)

case27 (Age=0.0 Ma, Steps=16000)

case27 (Age=0.0 Ma, Steps=16100)

case27 (Age=0.0 Ma, Steps=16200)

case27 (Age=0.0 Ma, Steps=16300)

case27 (Age=0.0 Ma, Steps=16400)

case27 (Age=0.0 Ma, Steps=16500)

case27 (Age=0.0 Ma, Steps=16600)

case27 (Age=0.0 Ma, Steps=16700)

case27 (Age=0.0 Ma, Steps=16800)

case27 (Age=0.0 Ma, Steps=16900)

case27 (Age=0.0 Ma, Steps=17000)

case27 (Age=0.0 Ma, Steps=17100)

case27 (Age=0.0 Ma, Steps=17200)

case27 (Age=0.0 Ma, Steps=17300)

case27 (Age=0.0 Ma, Steps=17400)

case27 (Age=0.0 Ma, Steps=17500)

case27 (Age=0.0 Ma, Steps=17600)

case27 (Age=0.0 Ma, Steps=17700)

case27 (Age=0.0 Ma, Steps=17800)

case27 (Age=0.0 Ma, Steps=17900)

case27 (Age=0.0 Ma, Steps=1800

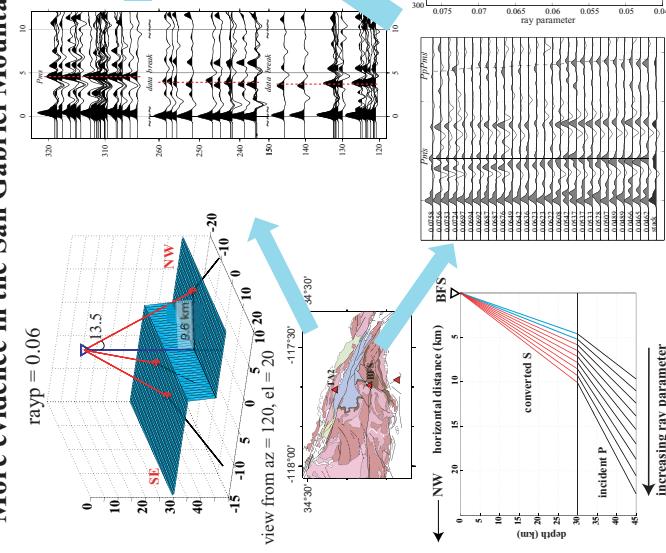
Evidence for large Moho offset in Southern California from Receiver Functions

Zhimei Yan & Robert W. Clayton

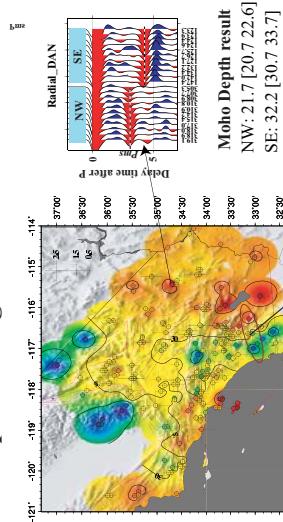
Abstract

Large offsets on the Moho are imaged among the back-azimuthal grouped Receiver Function (RF) beneath several stations in Southern California. Some of them correlate very well with the surficial geologically mapped faults, such as station TA2 on the San Andreas Fault (SAF), some of them occur in places where no surficial major faults exist, such as station DAN in Fennel Valley, Mojave Desert. Combined with synthetic RF waveform modeling, a notch structure is inferred on the eastern San Gabriel Mountains, where Moho shallows from 38 km north of the SAF, 34 km south of the San Gabriel Fault to ~29 km in between beneath the Mt. Baldy block.

More evidence in the San Gabriel Mountains

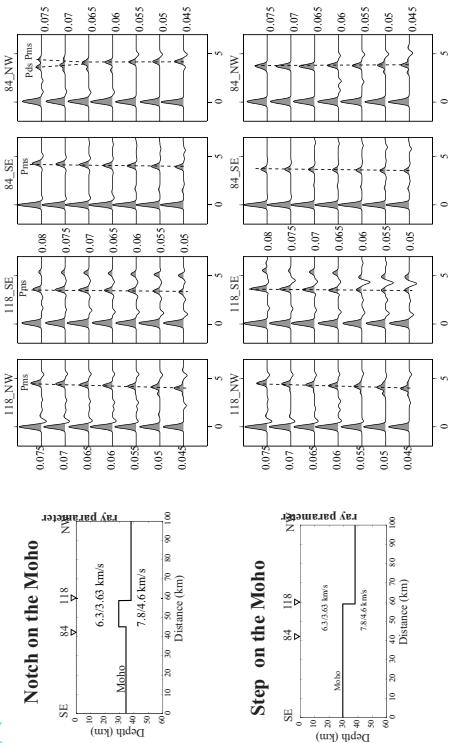


Examples of large offset on the Moho

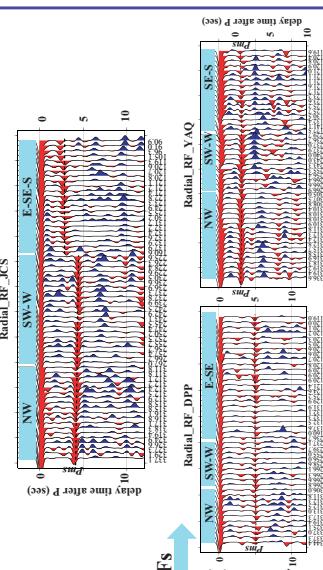
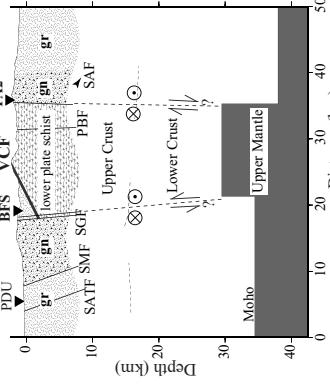


At station DAN in the Mojave Desert, a large Moho offset of ~10 km is inferred. At station JCS, which locates on the Elsinore Fault, large Pms of 4.5 sec is observed for the RFs from West, while small Pms arrival of 3 sec is observed for events from SE. The large Moho offset here is confirmed by the two nearby stations DPP and YAQ.

Synthetic Waveform modeling



Summary



Synthetic waveform modeling for the anomalous features in the receiver functions for two stations (TA2 and BFS) in the eastern San Gabriel Mountains indicates that a notch structure exists on the Moho, where the Moho shallows from 38 km north of the SAF, 34 km south of SGF to ~29 km beneath the Mt. Baldy block. The shallow Moho correlates very well with the surficial exposure of the lower plate Pelona Schist or places where the schist resides at very shallow depth. This is either related to the strike-slip movements along the two major faults or caused by differential uplifting.



Simulated Nonlinear Response of High-rise Buildings for the 2003 Tokachi-oki Earthquake Mw8.3

California Institute of Technology Civil Engineering Jing Yang jingy@caltech.edu
Thomas Heaton t@caltech.edu John Hall johnhall@caltech.edu

1. Introduction

Seismic waves from large subduction earthquakes are rich in long period waves that may be especially large in regions with local site amplification. The long-term global goal of our research is to investigate how well-designed modern high-rise buildings may perform in giant subduction earthquakes (e.g. Cascadia). Towards this goal, we are studying the Tokachi-oki 2003 earthquake (Mw8.3) which is the largest well recorded earthquake till now and was recorded by 276 strong motion stations located in Hokkaido Island. We use records from these stations to simulate the fully nonlinear seismic responses of 6- and 20-story steel moment-frame buildings designed according to both the U.S. 1994 UBC and also Japanese building code published in 1987. We consider buildings with both perfect welds and also with brittle welds whose fracture characteristics are similar to those observed in the 1994 Northridge earthquake.

From this research, we find that although Japanese code buildings are stronger, they are also stiffer which tends to increase the global forces experienced by Japanese buildings by more than 20% compared with U.S. code buildings. The net effect is that when considering collapse potential the Japanese buildings can sustain motions about 68% larger than the U.S. buildings. Moreover, our simulations indicate the building would have been strongly excited throughout the coastal region, with the potential for collapses in some locations.

3. Computational Models

We use Frame 2-D, a finite element method based on a fiber-element model that includes both material nonlinearities as well as geometric nonlinearities.

Steel Moment-Frame Building Models (symmetric):
• Heights: 6-story and 20-story buildings (plus one basement)
• Design Codes: 1994 UBC at seismic zone 3 and Japanese building code (1987)

• Welds conditions: brittle (prior to fracture) and perfect (weld fracture)

Figure 4: Pushover curves for 6-story and 20-story buildings. This analysis measures the actual strength of buildings. We can find that Japanese buildings are stronger than U.S. buildings and the presence of brittle welds significantly decreases the strength of a building.

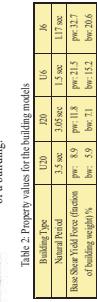


Figure 3: Definitions of the response parameters.

5. Collapse Factor

We introduce a new parameter named the "collapse factor" to describe the collapse possibility associated with simulated buildings. This safety factor is defined to be the scalar multiplier of the recorded ground motion that is required to cause collapse of the simulated building.

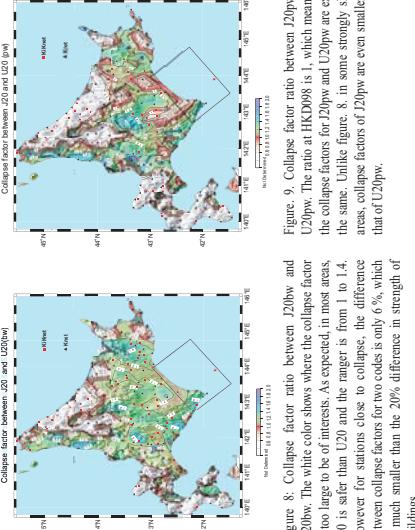


Figure 5: Collapse factor between J20 and U20bw

2. The 2003 Tokachi-oki Earthquake Mw8.3

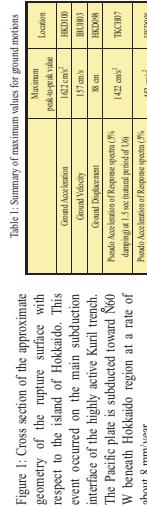
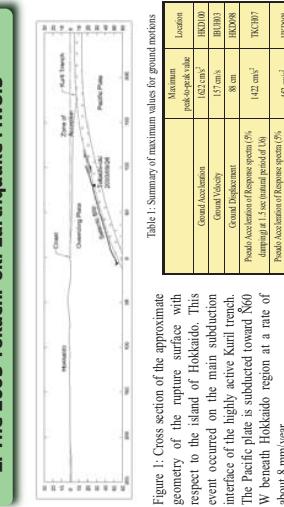


Figure 7: Maximum inter-story drift (in percent) for the U20 (20-story steel frame designed to IBC94) with brittle welds (bw). The maximum value is 5.38% and occurred at HKD098. Notice the buildings located in the northeast and the southwest regions in Hokkaido will suffer strong shaking although they are almost 200 km away from the epicenter.

4. Simulated Nonlinear Responses

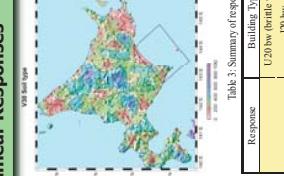


Figure 8: Collapse factor ratio between J20bw and U20bw

Figure 7: Collapse factor for U20bw. Although no buildings shown collapse in these simulations, increasing the amplitude of station HKD098 by only 6 % caused simulated collapse for U20bw. This is well within the uncertainty of this type of calculation.

2. The 2003 Tokachi-oki Earthquake Mw8.3

6. Future Work

We plan to simulate strong ground motions of the 2004 Sumatra earthquake by considering the strong motion recordings from the 2003 Tokachi-oki earthquake as empirical Green functions. Simulation the broad-band seismic body waves in the frequency band of interest is important for providing constraints on the strong motion simulation.

7. Conclusions

- The long-period ground motions recorded in the 2003 Tokachi-oki earthquake would have caused large inter-story drifts in 20-story flexible steel moment-resisting frame buildings designed according to both current U.S. and Japanese building codes.

- Although Japanese buildings are 20%-30% percent stronger than U.S. buildings, their capacity to resist collapse does not proportionally increase. Japanese buildings with brittle welds can sustain motions only 6% larger than corresponding U.S. buildings for station with significant collapse potential. And in some areas, Japanese buildings with perfect welds can sustain motions even smaller than U.S. buildings.

- Local soil geology plays an important role in the performance of high-rise buildings. Some basin areas which locate more than 200 km away from the epicenter amplify the long period motions large enough so that one could expect irreparable damage for 20-story buildings.

- The fracture of welds in the connections of beams and columns would dramatically reduce the strength of the buildings, as the fraction of welds in the connections of beams and columns would dramatically reduce the strength of the buildings, as

ACKNOWLEDGEMENT:
We thank the Tectonics Observatory funded by the Moore Foundation. Data from NIED